Universidade do Minho

## Escola de Ciências

Carla Maria Alves Ferreira

The Unsymmetric Tridiagonal Eigenvalue Problem

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# The Unsymmetric Tridiagonal Eigenvalue Problem 

Tese de Doutoramento em Ciências
Área de Conhecimento - Matemática

> Trabalho efectuado sob a orientação de
> Emeritus Professor Beresford Neill Parlett Professor Doutor Rui Manuel Silva Ralha

É AUTORIZADA APENAS A CONSULTA DESTA TESE PARA EFEITOS DE INVESTIGAÇÃO, MEDIANTE DECLARAÇÃO ESCRITA DO INTERESSADO, QUE A TAL SE COMPROMETE.

## Acknowledgments

The author would like to express her extreme gratitude to Professor Beresford Parlett for sharing his precious time and for making the work of this thesis so rewarding. His expertise, vision and clarity of thought greatly improved this work, as well as its presentation. Our discussions, over many afternoons, have proven enormously valuable in the development of this work and in shaping new ideas for future investigation.

The author is also grateful to Professor Rui Ralha for introducing her to the field of numerical linear algebra and for all the opportunities of research he has offered her, especially the initial encouragement for a long-term visit to the University of California at Berkeley. His suggestions and comments were always important.

The author is also thankful to many friends in the Mathematics Department of the University of Minho, particularly the ones with whom she had the chance to work in the last two academic years.

The author would like to thank to Professor Lisa Santos, Chair of the Mathematics Department, for having been extremely helpful with all the formalities that have made possible the completion of this thesis.

The author could never forget to express many thanks to her mother for her unconditional love and emotional support in every single stage of her journey. Also many thanks to her brothers for their constant advice and understanding.

The author also thanks all her friends who have made life in Berkeley an enriching experience.

Finally, the author is appreciative to the Center of Mathematics of the University of Minho for the financial support for this research.

# The Unsymmetric Tridiagonal Eigenvalue Problem 


#### Abstract

The development of satisfactory methods for reducing an unsymmetric matrix to tridiagonal form has been greatly hampered by the fact that there is not an accepted good algorithm for exploiting this form. Nevertheless, recently, promising elimination techniques for achieving a stable reduction to this form have been developed. But the standard QR algorithm destroys it immediately. Our work aims to fill this gap in the armoury of software tools for the matrix eigenvalue problem and so encourage the refinement of methods to reduce a matrix to tridiagonal form.

The progressive quotient difference algorithm with shifts (qds) was presented by Rutishauser as early as 1954. It is equivalent to the shifted LR algorithm written in a special notation for tridiagonals. The much more recent differential qds (dqds) is a sophisticated variant of qds. The first contribution of this thesis is a new algorithm, 3dqds, that consists of three dqds steps performed implicitly and such that real arithmetic is maintained in the presence of complex eigenvalues. One advantage of our algorithm over the Hessenberg QR algorithm is that it preserves the tridiagonal form and thus reduces both storage and time. We present some accuracy results comparing a Matlab implementation of 3dqds with the function eig of that software. These preliminary results suggest the robustness of 3dqds algorithm.

In contrast to the symmetric case, unsymmetric matrices can have a mixture of eigenvalues, some robust in the face of perturbations while others extremely sensitive. We present several condition numbers, some new, that take advantage of tridiagonal form. Ideally an algorithm should report these numbers along with each computed eigenvalue.

On the theoretical side, we present a rigorous proof of a surprising result. It is well known that the greater the ratio of adjacent eigenvalues, the faster LR converges. Nevertheless, in exact arithmetic, LR still converges even when all the eigenvalues are equal and the Jordan form is one big block.


# Cálculo de Valores Próprios de Matrizes Tridiagonais Não Simétricas 

## Resumo

O desenvolvimento de métodos satisfatórios para reduzir uma matriz não simétrica à forma tridiagonal tem sido fortemente travado pelo facto de que não existe um bom algoritmo aceite para explorar esta forma. Contudo, recentemente, promissoras técnicas de eliminação para realizar esta redução de maneira estável foram desenvolvidas. Mas o algoritmo QR standard destrói a forma tridiagonal imediatamente. O nosso trabalho pretende colmatar esta lacuna no conjunto das ferramentas computacionais para o problema de cálculo de valores próprios e assim encorajar o aperfeiçoamento de métodos para redução de uma matriz à forma tridiagonal.

O algoritmo qds (progressive quotient difference with shifts) foi introduzido por Rutishauser e remonta a 1954. É equivalente à versão shifted do algoritmo LR escrita numa notação especial para matrizes tridiagonais. O muito mais recente algoritmo dqds (differential qds) é uma versão sofisticada do qds. Uma contribuição desta tese é um novo algoritmo, 3dqds, que consiste em três passos do dqds realizados implicitamente e tal que a aritmética real é mantida na presença de valores próprios complexos. Uma vantagem do nosso algoritmo em relação ao Hessenberg QR é que preserva a forma tridiagonal e assim reduz a necessidade de espaço em memória e o tempo de execução. Apresentamos alguns resultados numéricos comparando uma implementaçao do 3dqds em Matlab com a função eig daquele software. Estes resultados preliminares sugerem a robustez do algoritmo 3dqds.

Em contraste com o caso simétrico, matrizes não simétricas podem ter um misto de valores próprios, alguns resistentes em face de perturbações enquanto outros extremamente sensíveis. Apresentamos vários números de condição, alguns novos, que tiram partido da forma tridiagonal. Idealmente, um algoritmo deve fazer acompanhar com estes números cada valor próprio calculado.

Do ponto de vista teórico, apresentamos também uma demonstração rigorosa de um resultado surpreendente. É bem conhecido que quanto maior for a razão entre valores próprios adjacentes, mais rapidamente o algoritmo LR converge. No entanto, em aritmética exacta, o algoritmo LR também converge mesmo quando todos os valores próprios são iguais e a forma de Jordan é um único bloco.

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## Introduction

A great deal of effort in science and engineering goes into eigenvalue computations. The symmetric case is well studied and there are good methods available. The unsymmetric case is intrinsically harder. The Matlab ${ }^{1}$ system [33] lets the user compute eigenvalues with one line of code, but this system, although wonderful for developing new numerical methods is too inefficient for day to day work in design and manufacturing. In addition, it cannot deal with the really large matrices, with order greater than 10000, that occur more and more often in applications.

So, we begin with the assumption that good algorithms are needed for real unsymmetric square matrices. Such algorithms have to be iterative in nature. When all eigenvalues are wanted then the preferred methods employ a sequence of similarity transformations which preserve the eigenvalues and gradually change the matrix to upper triangular form.

The reader might object at this point and say that the obvious strategy is to find the characteristic polynomial in a finite number of steps and then find the zeros of the polynomial. This approach died in the 1950's when it was appreciated that the coefficients of the characteristic polynomial are a too compact representation of the matrix eigenvalues; the eigenvalues are extremely sensitive to any uncertainty in the coefficients.

The next most compact pratical representation of a matrix is a tridiagonal form (all entries $(i, j)$ are zero unless $|i-j| \leq 1)$. All matrices can be reduced to such a form in a finite number of steps but the reduction is much easier in the symmetric case than in the unsymmetric one.

[^0]Now any iterative method will be much more efficient if the matrix is tridiagonal and if this form is preserved at each step. For example, Matlab does not preserve this form and that is the reason why it is not suitable for large matrices.

Now comes the question: is the tridiagonal form, in the unsymmetric case, also too sensitive, just as the companion matrix that gives the characteristic polynomial? If it is too sensitive, then, in general, there are two avenues of escape. Either we can try to determine classes of tridiagonals that do determine their eigenvalues to adequate accuracy or we can accept that the Hessenberg form is as far as we should go in reducing a full matrix to a more compact form.

No one really knows the answer to this question. Some work has been done in defining suitable measures of sensitivity, called condition numbers, but no one has studied them carefully. The formulae for the sensitivity of polynomial zeros as functions of the coefficients showed immediately that the condition numbers were going to be huge as soon as the degree goes into the hundreds. For tridiagonals the situation is not so clear. We explore this question in more detail in our study.

The differential quotient difference algorithm with shifts (dqds) was introduced by Fernando and Parlett in 1994 [16] to compute singular values of bidiagonal matrices to high relative accuracy but it may also be used to compute eigenvalues of tridiagonal matrices. In this thesis, based on previous work of Z . Wu [64], we propose a new algorithm for finding all the eigenvalues of a real unsymmetric tridiagonal matrix. This new algorithm, triple dqds (3dqds), incorporates three dqds steps implicitly. The motivation for 3dqds is to keep real arithmetic in the presence of complex shifts, efficiency and some elegance. One advantage of 3dqds over the standard Hessenberg QR (used by Matlab) is this property: it preserves the tridiagonal form. The preliminary numerical results show that 3dqds is a vital tool in the context of eigenvalue problems.

It is possible that perturbations to the tridiagonal entries have just as much effect as perturbing entries far from the diagonal. That is where the more refined condition numbers we are also presenting come in, particularly the ones based on the derivative of an eigenvalue
with respect to various matrix entries.
The reader will find that there is a great deal of technical detail in this thesis. The algorithms that we will present are quite complicated and it was necessary to get all the details correct in order for the programs to work properly.

Next we are going to sketch an outline of this thesis.
In Chapter 1 we introduce notation, give some background and explain the importance of the unsymmetric tridiagonal eigenvalue problem.

Chapter 2 describes different representations for tridiagonal matrices beginning with the LU factorization. Then, new relative condition numbers for measuring the sensitivity of eigenvalues of tridiagonals are presented, to see if tridiagonal form plays a big role in a perturbation study.

In Chapter 3 we describe the LR algorithm, giving emphasis to the implicit double shifted version on an Hessenberg matrix. Then we explain the qds algorithm and show the relation to shifted LR on a tridiagonal matrix with superdiagonal entries of 1's.

Chapter 4 deals with convergence results for basic LR algorithm. The central part of this chapter is the proof of a new convergence result of basic LR algorithm on a real unreduced tridiagonal matrix with a one-point spectrum. The Jordan form is one big Jordan block.

Chapter 5 shows all the details of the derivation of a first version of the triple dqds - an algorithm that performs implicitly three steps of simple dqds keeping real arithmetic in the presence of complex shifts.

Chapter 6 explores the connection between dqds and the Gram-Schmidt orthogonalization process to reveal new results about triple dqds. These results lead to the more elegant and more efficient final version 3dqds.

Chapter 7 gives a preliminary numerical comparison between 3dqds and existing Matlab's function eig. Although the set of test matrices is small, the numerical results permit us to conclude that 3dqds is a competitive algorithm.

Finally, in Chapter 8 we present a summary of our work and briefly set up some plans for future work.

## Chapter 1

## Setting the scene

This chapter describes notation, introduces definitions and some basic results, discusses canonical forms and perturbation theory in the context of eigenvalue problems. Then devotes attention to the unsymmetric tridiagonal eigenvalue problem describing the connection to the general unsymmetric eigenvalue problem and to Bessel matrices.

### 1.1 The eigenvalue problem

The eigenvalues of matrices or linear operators play a part in a very large number of applications, both theoretical and practical. In Chatelin [5, Chapter 3] we can find examples from diverse disciplines that show the extent of applications: they range from mathematics to chemistry, and to the dynamics of structures, even on economics. The determination of matrix eigenvalues is generally called the eigenvalue problem and it is a central topic in numerical linear algebra. From Golub [22] we quote what follows.

The eigenvalue problem for square matrices $A$, that is the determination of nontrivial solutions of $A x=\lambda x$, is inherently nonlinear and this leads to many computational problems. Computation of the eigenvalues $\lambda$ via the characteristic equation

$$
\operatorname{det}(A-\lambda I)=0
$$

is, except for very special cases, not an option since the coefficients of the characteristic equation cannot be computed from determinant evaluations in a numerical stable way. And even if the characteristic equation could be determined accurately, then the computation of its roots, in finite precision, may be highly unstable since small perturbations in the coefficients may lead to large perturbations of the roots.

The numerical computation of the associated eigenvectors and generalized eigenvectors is even more delicate, in particular when eigenvectors of $A$ make small angles with each other. In the limiting case, when the matrix is defective, $A$ can be reduced to the Jordan canonical form, but arbitrary small perturbations in $A$ may yield a nondefective matrix. This leads to many challenging numerical questions, which give rise to the central problem: how can we compute eigenvalues and eigenvectors in an efficient manner and how accurate are they? [...]

A method that is of great significance and serves as the basis for many algorithms is the Power iteration. It is still in use, but most frequentely as (implicit) part of more efficient techniques, e.g., krylov methods, inverse iteration, QR-method.

What becomes clear is that all these methods are of an iterative nature, and this is necessarily the case, since if there were a method of computing the eigenvalues of an $n$th order matrix in a finite number of computations, depending only on $n$, then this would be in contradiction with the fundamental theorem of Abel-Ruffini (and also a well-known result in Galois theory) that no such algorithm exists for the computation of the roots of a general polynomial of degree greater than 4. Hence, an algorithm for a matrix with a general structure (that is, not a diagonal matrix or a triangular matrix or alike) is necessarily iterative and the problem is to identify iterative algorithms which have a fast rate of convergence and lead to accurate results.

### 1.2 Notation, definitions and basic results

We describe the notation used hereafter and briefly set up well-known definitions and basic facts needed in this thesis.

The vector space of all real $m \times n$ matrices is denoted by $\mathbb{R}^{m \times n}$ and the vector space of real column $n$-vectors by $\mathbb{R}^{n}$. Similarly, $\mathbb{C}^{m \times n}$ denotes the vector space of $m \times n$ matrices with complex entries and $\mathbb{C}^{n}$ the vector space of complex column $n$-vectors.

Generally, we will have

| capital letters | $A, H, \Delta$ | for matrices |
| ---: | :---: | :--- |
| (double) subscripted lower case letters | $a_{i j}, h_{i j}, \delta_{i j}$ | for matrix elements |
| boldfaced lower case letters | $\boldsymbol{x}, \boldsymbol{c}, \boldsymbol{h}$ | for column vectors |
| subscripted lower case letters | $x_{k}, c_{k}, h_{k}$ | for vector elements |
| lower case Greek letters | $\alpha, \beta, \gamma, \theta$ | for scalars |

We may also denote a matrix $A=\left(a_{i j}\right) \in \mathbb{C}^{m \times n}$ by its columns, that is, we may choose to say $A=\left[\begin{array}{llll}\boldsymbol{a}_{1} & \boldsymbol{a}_{2} & \cdots & \boldsymbol{a}_{n}\end{array}\right]$. For instance, the $n \times n$ identity matrix will be denoted by $I_{n}$ and we have $I_{n}=\left[\begin{array}{llll}\boldsymbol{e}_{1} & \boldsymbol{e}_{2} & \cdots & \boldsymbol{e}_{n}\end{array}\right]$.

We denote the transpose of $A$ by $A^{T}$ and, if $A$ is complex, the conjugate transpose by $A^{*}$, not $A^{H}$. A square matrix $A$ is symmetric when $A=A^{T}$, otherwise is said to be unsymmetric or nonsymmetric. If $A=A^{*}$ then $A$ is an hermitian matrix.

A square matrix $A$ is upper triangular (lower triangular) if $a_{i j}=0$ when $i>j(i<j)$. And $A$ is said to be diagonal if $a_{i j}=0, i \neq j$. A diagonal matrix will be written $\operatorname{diag}\left(a_{11}, a_{22}, \ldots, a_{n n}\right)$ and, for $\boldsymbol{d}=\left(d_{i}\right) \in \mathbb{C}^{n}$, we define $\operatorname{diag}(\boldsymbol{d}):=\operatorname{diag}\left(d_{1}, \ldots, d_{n}\right)$.

Other definitions and notation will be introduced when needed.

Precise proofs of the results that will be presented in the following sections can be found in [51, Chapter 6], [34, Chapter 3], [21, Chapter 7], [11, Chapter 4] or [45, Chapter 12].

### 1.2.1 Eigenvalues

The notions of eigenvalue and eigenvector do not depend on length, angle or inner product.

Definition 1.2.1 Let $A \in \mathbb{C}^{n \times n}$ and $\boldsymbol{x} \in \mathbb{C}^{n}$. Then $\boldsymbol{x}$ is a (right) eigenvector of $A$ for the eigenvalue $\lambda$ if $\boldsymbol{x}$ satisfies

$$
\begin{equation*}
A \boldsymbol{x}=\lambda \boldsymbol{x}, \quad \boldsymbol{x} \neq \mathbf{0} . \tag{1.1}
\end{equation*}
$$

Each nonzero multiple of $\boldsymbol{x}$ is also an eigenvector for $\lambda$ and $\lambda-\mu$ is an eigenvalue of $A-\mu I$ with eigenvector $\boldsymbol{x}$, for every scalar $\mu$.

Fact 1.2.1 Eigenvectors corresponding to distinct eigenvalues are linearly independent.

By the theory of linear equations, (1.1) has a non-zero solution $\boldsymbol{x}$ if and only if the characteristic polynomial of $A, \chi$, defined by

$$
\chi(\lambda) \equiv \chi_{A}(\lambda) \equiv \operatorname{det}(A-\lambda I),
$$

verifies $\chi(\lambda)=0$.
The roots of $\chi$ are then the eigenvalues ${ }^{1}$ of $A$ and are also called the characteristic values, characteristic roots or latent roots of $A$.

Polynomial $\chi$ is of degree $n$ and its leading term is $(-1)^{n} \lambda^{n}$. So $A$ can at most have $n$ eigenvalues, some of which may be repeated.

Let $\lambda_{1}, \lambda_{2}, \ldots, \lambda_{m}$ be the distinct eigenvalues of $A$. Then $\chi$ can be represented as

$$
\chi(\lambda)=\left(\lambda_{1}-\lambda\right)^{n_{1}} \cdots\left(\lambda_{m}-\lambda\right)^{n_{n}}
$$

where the $n_{i}$ are positive integers with $\sum_{i=1}^{m} n_{i}=n$. The number $n_{i}$ is the algebraic multiplicity of $\lambda_{i}, i=1, \ldots, m$. If $n_{i}=1$, then $\lambda_{i}$ is called simple.

Since the characteristic polynomial of a real matrix has real coefficients, the complex eigenvalues of a real matrix occur in conjugate pairs.

[^1]The set of the eigenvalues $\lambda_{i}, i=1, \ldots, m$, in the complex plane, constitutes the spectrum of $A$ and it will be denoted by $\operatorname{spec}(A)$. We call

$$
\rho(A) \equiv \max \{|\lambda|: \lambda \in \operatorname{spec}(A)\}
$$

the spectral radius of $A$.
Fact 1.2.2 $\operatorname{spec}\left(A^{*}\right)=\{\bar{\lambda}: \lambda \in \operatorname{spec}(A)\}$.

This result states that if $\lambda$ is an eigenvalue of $A$, there is a nonzero vector $\boldsymbol{y}$ such that $A^{*} \boldsymbol{y}=\boldsymbol{y} \bar{\lambda}$, or, equivalently, $\boldsymbol{y}^{*} A=\lambda \boldsymbol{y}^{*}$. We say that $\boldsymbol{y}^{*}$ is a row eigenvector of $A$. Commonly, $\boldsymbol{y}$ is called a left eigenvector of $A$. This means that $\overline{\boldsymbol{y}}$ is a right eigenvector of $A^{T}$.

Fact 1.2.3 For right and left eigenvectors $\boldsymbol{x}$ and $\boldsymbol{y}$ associated with a simple eigenvalue $\lambda$, $\boldsymbol{y}^{*} \boldsymbol{x} \neq 0$.

Fact 1.2.4 If $\lambda$ and $\mu$ are two distinct eigenvalues with right and left eigenvectors $\boldsymbol{x}_{\lambda}$ and $\boldsymbol{y}_{\mu}$, respectively, then $\boldsymbol{y}_{\mu}^{*} \boldsymbol{x}_{\lambda}=0$.

For an eigenvalue $\lambda$ with right and left eigenvectors $\boldsymbol{x}$ and $\boldsymbol{y}$, respectively, $\{\lambda, \boldsymbol{x}\}$ is called an eigenpair and $\left\{\lambda, \boldsymbol{x}, \boldsymbol{y}^{*}\right\}$ an eigentriple.

### 1.2.2 Canonical forms

Most of the computational methods involve reducing a matrix into simpler or even canonical forms, from which it is easy to compute its eigenvalues and eigenvectors. These transformations are called similarity transformations. The two most common canonical forms are the Jordan form and the Schur form. For historical reasons we will also refer to the Frobenius canonical form. The Schur canonical form is the one that is more useful for pratical use.

Fact 1.2.5 Let $B=S A S^{-1}$. Then $A$ and $B$ have the same eigenvalues and $x$ is a right eigenvector of $A$ if and only if $S x$ is a right eigenvector of $B$.

The mapping $A \mapsto S A S^{-1}$ is a similarity transformation of $A$. If $S$ is unitary (orthogonal in the real case), that is, $S^{-1}=S^{*}\left(S^{-1}=S^{T}\right)$, we say that the transformation is a unitary (orthogonal) similarity transformation. Matrices $A$ and $B=S A S^{-1}$ are called similar matrices.

To motivate the Jordan and Schur forms, let us just recall that for a diagonal or a triangular matrix the eigenvalues are easy to compute: they are simply its diagonal entries.

Below we will see that a matrix in Jordan or Schur form is triangular.
Another particularly important form is the quasi-triangular form, which is a special case of the block triangular form. A square matrix $A$ is block upper triangular if it can be partitioned in the form

$$
A=\left[\begin{array}{cccc}
A_{11} & A_{12} & \ldots & A_{1 l}  \tag{1.2}\\
0 & A_{22} & \ldots & A_{2 l} \\
\vdots & \vdots & \ddots & \vdots \\
0 & 0 & \ldots & A_{l l}
\end{array}\right],
$$

where each diagonal block $A_{i i}$ is square. If each diagonal block is of order at most two, then $A$ is said to be in quasi-triangular form. Because $\operatorname{det}(A-\lambda I)=\prod_{i=1}^{l} \operatorname{det}\left(A_{i i}-\lambda I\right)$ we have

$$
\operatorname{spec}(A)=\bigcup_{i=1}^{l} \operatorname{spec}\left(A_{i i}\right)
$$

Next theorem introduces Jordan canonical form and it is followed by some related facts that we will need in our convergence proofs in Chapter 4.

Theorem 1.2.1 Jordan canonical form. Given $A \in \mathbb{C}^{n \times n}$, there exists an invertible matrix $X$ such that $X^{-1} A X=J$, where $J$ is in Jordan canonical form. This means that $J$ is block diagonal, with $J=\operatorname{diag}\left(J_{n_{1}}\left(\lambda_{1}\right), J_{n_{2}}\left(\lambda_{2}\right), \ldots, J_{n_{k}}\left(\lambda_{k}\right)\right)$ and

$$
J_{n_{i}}\left(\lambda_{i}\right)=\left[\begin{array}{cccc}
\lambda_{i} & 1 & & 0 \\
& \ddots & \ddots & \\
& & \ddots & 1 \\
0 & & & \lambda_{i}
\end{array}\right]_{n_{i} \times n_{i}}
$$

$J$ is unique, up to permutations of its diagonal blocks.

There is a good deal of terminology associated with Jordan canonical form. First, each block $J_{n_{i}}\left(\lambda_{i}\right)$ is called a Jordan block with eigenvalue $\lambda_{i}$. The determinants

$$
\operatorname{det}\left(\lambda I_{n_{i}}-J_{n_{i}}\right)=\left(\lambda-\lambda_{i}\right)^{n_{i}}
$$

are called the elementary divisors of $A$. If $n_{i}=1$ the elementary divisor is called linear.
If all the elementary divisors are linear, so that $J$ is diagonal, $A$ is said to be diagonalizable or nondefective; otherwise $A$ is called defective.

If we denote matrix $X$, that reduces $A$ to Jordan form, as

$$
X=\left[\begin{array}{llllllllll}
x_{1}^{(1)} & \ldots & x_{n_{1}}^{(1)} & x_{1}^{(2)} & \ldots & x_{n_{2}}^{(2)} & \ldots & \boldsymbol{x}_{1}^{(k)} & \ldots & \boldsymbol{x}_{n_{k}}^{(k)}
\end{array}\right]
$$

then, for each Jordan block $J_{n_{i}}\left(\lambda_{i}\right)$, the corresponding vectors $\boldsymbol{x}_{1}^{(i)}, \ldots, \boldsymbol{x}_{n_{i}}^{(i)}, i=1, \ldots, k$, satisfy

$$
\begin{aligned}
A \boldsymbol{x}_{1}^{(i)} & =\lambda_{i} \boldsymbol{x}_{1}^{(i)} \\
A \boldsymbol{x}_{j+1}^{(i)} & =\lambda_{i} \boldsymbol{x}_{j+1}^{(i)}+\boldsymbol{x}_{j}^{(i)}, \quad j=1, \ldots, n_{i}-1 .
\end{aligned}
$$

Vectors $x_{j+1}^{(i)}, j=1, \ldots, n_{i}-1$ are called generalized eigenvectors or principal vectors of grade $j+1$, verifying

$$
\left(A-\lambda_{i} I\right)^{j+1} \boldsymbol{x}_{j+1}^{(i)}=0, \quad\left(A-\lambda_{i} I\right)^{j} \boldsymbol{x}_{j+1}^{(i)}=\boldsymbol{x}_{1}^{(i)} \neq \mathbf{0}
$$

Fact 1.2.6 If all the eigenvalues of a matrix $A$ are distinct, then $A$ is diagonalizable.
A defective matrix does not have $n$ linearly independent eigenvectors.
Fact 1.2.7 A Jordan block has one right eigenvector $\boldsymbol{e}_{1}=\left[\begin{array}{llll}1 & 0 & \cdots & 0\end{array}\right]^{T}$ and one left eigenvector $\boldsymbol{e}_{n}=\left[\begin{array}{llll}0 & 0 & \cdots & 1\end{array}\right]^{T}$.

The number of linearly independent eigenvectors associated with an eigenvalue $\lambda$ is called the geometric multiplicity of $\lambda$ and does not exceed its algebraic multiplicity. It corresponds to the number of Jordan blocks associated with $\lambda$.

A matrix is defective if and only if the geometric multiplicity of at least one of its eigenvalues is less than the algebraic multiplicity - in this case, commonly, we call these eigenvalues briefly as multiple eigenvalues.

A matrix $A$ is nonderogatory if every eigenvalue has geometric multiplicity 1 . This means that different Jordan blocks $J_{n_{i}}$ correspond to distinct $\lambda_{i}$.

The $(i, j)$ minor of a square matrix $A$ is defined as the determinant of the submatrix obtained by removing the $i^{\text {th }}$ row and the $j^{\text {th }}$ column of $A$. The $k^{\text {th }}$ leading principal minor is the determinant of the first $k$ rows and $k$ columns of $A$.

The Companion matrix of a monic polynomial

$$
p(\lambda) \equiv \lambda^{n}+a_{n-1} \lambda^{n-1}+\ldots+a_{1} \lambda+a_{0}
$$

is the matrix

$$
C_{p}:=\left[\begin{array}{ccccc}
0 & & & & -a_{0} \\
1 & 0 & & & -a_{1} \\
& \ddots & \ddots & & \vdots \\
& & \ddots & 0 & -a_{n-2} \\
& & & 1 & -a_{n-1}
\end{array}\right] .
$$

Since $p(\lambda)=(-1)^{n} \operatorname{det}\left(C_{p}-\lambda I\right)=\operatorname{det}\left(\lambda I-C_{p}\right)$, the zeros of the polynomial $p$ are the eigenvalues of $C_{p}$. It is easy to see that $C_{p}$ is a nonderogatory matrix: for all $\lambda$, the $(1, n)$ minor of $C_{p}-\lambda I$ is always 1 ; consequently $\operatorname{rank}\left(C_{p}-\lambda I\right) \geq n-1$ and the dimension of the null space of $C_{p}-\lambda I$ is either 0 (when $\lambda$ is not an eigenvalue) or 1 (when $\lambda$ is an eigenvalue); this means that there is only one linearly independent eigenvector to each eigenvalue.

Fact 1.2.8 $A$ matrix $A$ is similar to the companion matrix of its characteristic polynomial if and only if $A$ is nonderogatory.

A $m \times n$ matrix $R=\left(r_{i j}(x)\right)$ whose elements $r_{i j}$ are rational functions, that is, ratios of polynomials, is a rational matrix.

Theorem 1.2.2 Frobenius, or rational, canonical form Given $A \in \mathbb{C}^{n \times n}$, there exists an invertible matrix $P$ such that $P^{-1} A P=F$, where $F$ is in Frobenius canonical form. This means that $F$ is the direct sum of the companion matrices of the elementary divisors of $\lambda I_{n}-A$.

If $A \in \mathbb{R}^{n \times n}$, the Frobenius canonical form of $A$ depends on which field ( $\mathbb{R}$ or $\mathbb{C}$ ) over which it is considered.

This form does not say much about eigenvalues but it is, by construction, the matrix with fewest nonzero entries that can be achieved by rational operations on the data, but it is too condensed. The Jordan canonical form tells us all we want to know about eigenvalues but can not be computed stably in general. So, it is used in theory but is very hard to compute because it is very unstable in the face of uncertainty. Thus, these two forms are not used in eigenvalue computations and most of the algorithms will aim to compute the Schur form instead.

Theorem 1.2.3 Schur canonical form. Given $A \in \mathbb{C}{ }^{n \times n}$, there exists a unitary matrix $Q$ and an upper triangular matrix $T$ such that $Q^{*} A Q=T$.

A matrix $A$ is said to be normal if $A A^{*}=A^{*} A$.

Fact 1.2.9 $A$ matrix $A$ is normal if and only if there exists a unitary matrix $Q$ such that $Q^{*} A Q=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right)$.

As said before, a real matrix $A$ can have complex eigenvalues and, therefore, there is not always a real triangular matrix with the same eigenvalues as $A$. So, we must either use complex numbers or sacrifice the triangular canonical form. Because it will be cheaper to compute, we prefer a canonical form that uses only real numbers and we will settle for a quasi-triangular form.

Theorem 1.2.4 Real Schur canonical form. If $A \in \mathbb{R}^{n \times n}$, there exists a real orthogonal matrix $V$ such that $V^{T} A V=T$ is real quasi-upper triangular. Its eigenvalues are the eigenvalues of its diagonal blocks. The $1 \times 1$ blocks corresponds to real eigenvalues and the $2 \times 2$ blocks to complex conjugate pairs of eigenvalues.

### 1.2.3 Perturbation theory

In what follows we will discuss how changes in the entries of a matrix $A$ affect the spectrum. If $\lambda$ is a simple eigenvalue of $A$, for a given matrix $\delta A$, we can identify an eigenvalue $\lambda+\delta \lambda$ of the perturbed $A+\delta A$ corresponding to $\lambda$ : the closest one to $\lambda$.

So, we want to relate the size of the matrix perturbation to the size of the eigenvalue change. The norms most widely used in matrix computations are the 1 -norm, the Frobenius norm and the $\infty$-norm. Unlike the 2 -norm (an example of a $p$-norm and also of great importance) the former matrix norms are easy to calculate. And for a consistent matrix norm $\|$.$\| we have, for any matrix A$ of order $n$,

$$
\rho(A) \leq\|A\| .
$$

For a real matrix $A$ and a matrix norm $\|\cdot\|$, the quantity

$$
\operatorname{cond}(A) \equiv\|A\|\left\|A^{-1}\right\|
$$

defines the condition number of the matrix $A$. Note that $\operatorname{cond}(A)=\operatorname{cond}(\tau A)$ for any scalar $\tau \neq 0$, that is, cond is scaling invariant by multiplication.

Next theorem relates eigenvalue condition numbers to the condition number of the matrix of all eigenvectors.

Theorem 1.2.5 Bauer-Fike (classical version). If $\mu$ is an eigenvalue of $A+\delta A \in \mathbb{C}^{n \times n}$ and $X^{-1} A X=D=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right)$, then

$$
\min _{\lambda \in \operatorname{spec}(A)}|\lambda-\mu| \leq \operatorname{cond}_{p}(X)\|\delta A\|_{p}
$$

where $\|\cdot\|_{p}$ denotes any of the $p$-norms.

Extreme eigenvalue sensitivity for a matrix $A$ cannot occur if $A$ is normal. On the other hand, nonnormality does not necessarily imply eigenvalue sensitivity. Indeed, a nonnormal matrix can have a mixture of well-conditioned and ill-conditioned eigenvalues. For this reason, it is beneficial to refine the perturbation theory so that it is applicable to individual eigenvalues and not the spectrum as a whole.

Theorem 1.2.6 Let $\lambda$ be a simple eigenvalue of $A$ with right eigenvalue $\boldsymbol{x}$ and left eigenvalue $\boldsymbol{y}$, normalized so that $\|\boldsymbol{x}\|_{2}=\|\boldsymbol{y}\|_{2}=1$. Let $\lambda+\delta \lambda$ be the corresponding eigenvalue of $A+\delta A$. Then

$$
|\delta \lambda| \leq \frac{\|\delta A\|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}+\mathcal{O}(\|\delta A\|)^{2}=\frac{\|\delta A\|}{\cos (\theta(\boldsymbol{y}, \boldsymbol{x}))}+\mathcal{O}(\|\delta A\|)^{2}
$$

where $\theta(\boldsymbol{y}, \boldsymbol{x})$ is the acute angle between $y$ and $x$. In other words, $\kappa_{\lambda} \equiv \frac{1}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}$ is the condition number of the eigenvalue $\lambda$. It is known as Wilkinson's condition number.

If the eigenvectors are not normalized the condition number $\kappa_{\lambda}$ is given by

$$
\kappa_{\lambda}=\frac{\|\boldsymbol{x}\|\|\boldsymbol{y}\|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|} .
$$

Each eigenvalue has its own condition number. Some eigenvalues can be well-conditioned and some can be ill-conditioned. A big condition number means sensitivity; ill-conditioned eigenvalues are hard to compute accurately because they are not well-defined.

The right and left eigenvectors of a Jordan Block are orthogonal, and so Wilkinson's condition number of its eigenvalue $\lambda$ will be $\kappa_{\lambda}=\infty$. At the other extreme are the normal matrices which have condition number $\kappa_{\lambda}=1$. So, the symmetric eigenvalue problem is always a well-conditioned problem.

Theorem 1.2.6 is useful only for sufficiently small $\|\delta A\|$. Next theorem increases the condition number by a factor on $n$ but is true for any modification $\delta A$.

Theorem 1.2.7 Bauer-Fike. Let A have all simple eigenvalues, i.e., be diagonalizable. Call these eigenvalues $\lambda_{i}, i=1, \ldots, n$, and let $\boldsymbol{x}_{i}$ and $\boldsymbol{y}_{i}$ be the right and left eigenvectors for $\lambda_{i}$, respectively, normalized so that $\left\|\boldsymbol{x}_{i}\right\|_{2}=\left\|\boldsymbol{y}_{i}\right\|_{2}=1$. Then the eigenvalues of $A+\delta A$ lie in disks $B_{i}$, where $B_{i}$ has center $\lambda_{i}$ and radius $\frac{n}{\left|\boldsymbol{y}_{i}^{*} \boldsymbol{x}_{i}\right|}\|\delta A\|_{2}$.

More generally, we can say that

$$
\forall \tilde{\lambda} \in \operatorname{spec}(A+\delta A): \min _{i}\left|\tilde{\lambda}-\lambda_{i}\right| \leq n \max _{i} \frac{1}{\left|\boldsymbol{y}_{i}^{*} \boldsymbol{x}_{i}\right|}\|\delta A\|_{2} .
$$

The proof of this theorem uses a useful inclusion result for eigenvalues, the Gerschgorin theorem, which we write below.

Theorem 1.2.8 Gerschgorin. For a matrix $A=\left(a_{i j}\right) \in \mathbb{C}^{n \times n}$ it holds that

$$
\operatorname{spec}(A) \subset \bigcup_{i=1}^{n} \mathcal{D}_{i} \quad \text { with } \quad \mathcal{D}_{i}:=\left\{\lambda \in \mathbb{C}:\left|\lambda-a_{i i}\right| \leq \sum_{j \neq i}\left|a_{i j}\right|\right\}, i=1, \ldots, n
$$

Disks $\mathcal{D}_{i}$ are called Gerschgorin disks of the matrix $A$.

The proof of the theorem shows not only that each eigenvalue of $A$ must lie in a Gerschgorin disk, but also that if the $i$ th component of an eigenvector is maximal, then the corresponding eigenvalue must lie in the $i$ th disk.

Theorem 1.2.9 If $k$ Gerschgorin disks of the matrix $A$ are disjoint from the other disks, then exactly $k$ eigenvalues of $A$ lie in the union of the $k$ disks.

Since a multiple eigenvalue has infinite Wilkinson's condition number, being close to a matrix which has multiple eigenvalues implies ill-conditioning. The bigger is the condition number of an eigenvalue, the closer is the matrix to one with a multiple eigenvalue.

Theorem 1.2.10 Let $\lambda$ be a simple eigenvalue of $A$, with unit right and left eigenvectors $\boldsymbol{x}$ and $\boldsymbol{y}$ and condition number $c=\frac{1}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}$. Then there is a $\delta A$ such that $A+\delta A$ has a multiple eigenvalue at $\lambda$ and

$$
\frac{\|\delta A\|_{2}}{\|A\|_{2}} \leq \frac{1}{\sqrt{c^{2}-1}}
$$

When $c \gg 1$, that is the eigenvalue is ill-conditioned, then the upper bound on the distance is $\frac{1}{\sqrt{c^{2}-1}} \approx \frac{1}{c}$, the reciprocal of the condition number.

Finally, we relate the condition numbers of the eigenvalues to the smallest possible condition number cond $(S)$ of any $S$ that diagonalizes $A$.

Theorem 1.2.11 Let A be diagonalizable with eigenvalues $\lambda_{i}$ and right and left eigenvectors $\boldsymbol{x}_{i}$ and $\boldsymbol{y}_{i}$, respectively, normalized so that $\left\|\boldsymbol{x}_{i}\right\|_{2}=\left\|\boldsymbol{y}_{i}\right\|=1$. Let us suppose that $S$ satisfies $S^{-1} A S=\Delta=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right)$. Then

$$
\max _{i} \frac{1}{\left|\boldsymbol{y}_{i}^{*} \boldsymbol{x}_{i}\right|} \leq\|S\|_{2}\left\|S^{-1}\right\|_{2} .
$$

If we choose $S=\left[\begin{array}{llll}x_{1} & x_{2} & \cdots & x_{n}\end{array}\right]$ then

$$
\|S\|_{2}\left\|S^{-1}\right\|_{2} \leq n \max _{i} \frac{1}{\left|\boldsymbol{y}_{i}^{*} \boldsymbol{x}_{i}\right|},
$$

that is the condition number of $S$ is within a factor of $n$ of its smallest value.

### 1.2.4 Relative errors and model of arithmetic

Let $\hat{\sigma}$ be an approximation to a real number $\sigma$. The most useful measures of the accuracy of $\hat{\sigma}$ is the absolute error

$$
E_{\mathrm{abs}}(\hat{\sigma})=|\sigma-\hat{\sigma}|,
$$

and its relative error

$$
E_{\mathrm{rel}}(\hat{\sigma})=\frac{|\sigma-\hat{\sigma}|}{|\sigma|}, \quad \sigma \neq 0
$$

Writing

$$
\hat{\sigma}=\sigma(1+\rho),
$$

an equivalent definition of relative error will be $E_{\text {rel }}(\hat{\sigma})=|\rho|$. Some authors omit the absolute values from these definitions.

In scientific computation, where answers to problems can vary enormously in magnitude, it is usually the relative error that is of interest, because it is scale independent: scaling $\sigma \rightarrow \alpha \sigma$ and $\hat{\sigma} \rightarrow \alpha \hat{\sigma}$ leaves $E_{\text {rel }}(\hat{\sigma})$ unchanged.

Relative error is connected with the notion of correct significant digits. But while the number of correct significant digits provides a useful way in which to think about the accuracy of an approximation, the relative error is a more precise measure.

## Floating point arithmetic

The maximum relative representation error in a floating point arithmetic with $p$ digits and base $\beta$ is $\frac{1}{2} \beta^{1-p}$. When the true value of a computation $a \odot b$ (where $\odot$ is one of the four binary operations,,$+- *$ and $/$ ) cannot be represented exactly as a floating point number, it must be approximated by a nearby floating point number. We denote this approximation by $f l(a \odot b)$. The difference $(a \odot b)-f l(a \odot b)$ is called the roundoff error. If $a \odot b$ is within the exponent range (otherwise we get overflow or underflow), then we assume the model of arithmetic

$$
\begin{equation*}
f l(a \odot b)=(a \odot b)(1+\delta), \quad|\delta| \leq \varepsilon \tag{1.3}
\end{equation*}
$$

The quantity $\varepsilon$ is called variously unit roundoff, machine precision or macheps. If we round as accurately as possible, $\varepsilon$ is equal to the maximum relative representation error $\frac{1}{2} \beta^{1-p}$.

The terms accuracy and precision are often confused or used interchangeable, but it is worth making a distinction between them. Accuracy refers to the absolute or relative error of an approximate quantity. Precision is the accuracy with which the basic arithmetic operations $+,-, *, /$ are performed and for floating point arithmetic is measured by the unit roundoff (1.3).

The IEEE standard for binary arithmetic is now common. It includes two kinds of floating point numbers: single precision (32 bits long) and double precision ( 64 bits long) that allow 24 and $53 p$ bits of precision, respectively.

Normally the "big oh" notation $\mathcal{O}$ implies a limiting process. However, in this thesis $\mathcal{O}$ will be a synonym for "of the order of magnitude of" and the usage will be clear from the context. For example, $\mathcal{O}(\varepsilon)$ will mean "of the order of magnitude of machine precision".

### 1.3 The unsymmetric tridiagonal eigenvalue problem

Tridiagonal matrices have received a great deal of attention since the 1950's. In the symmetric case, every matrix can be stably reduced to a similar tridiagonal matrix by a finite number of elementary orthogonal similarity transformations. In the unsymmetric case, reduction to tridiagonal form is also possible but we have to use non-orthogonal similarities transformations and the reduction may not be stable.

But why should we be interested in tridiagonal matrices? There is a need for eigenvalue methods capable of exploiting and respecting the elegant structure of an unsymmetric tridiagonal matrix so that the development of methods for tridiagonal reduction could be encouraged.

Unsymmetric tridiagonal matrices arise naturally as a result of the execution of the two-sided Lanczos algorithm (see, for instance, Parlett [39]). When applied to a matrix, the two-sided Lanczos algorithm builds a $k \times k$ unsymmetric tridiagonal matrix at the end of the Lanczos step $k$. This is a candidate method for the reduction of a nonsymmetric matrix to tridiagonal form but in practice it is confined to large sparse matrices. Other methods to perform this task have been proposed in Geist [20], Dongarra [13] and, more recently, in Sidje [50].

Tridiagonals also appear as primary data. For example, they are related to orthogonal polynomials because there is a remarkable three-term recurrence relation among these polynomials. In particular, there are the special Bessel polynomials, which also satisfy a three-term recurrence relation. Tridiagonals also arise in other numerical methods such as exponential interpolation [1].

The eigenvalues and eigenvectors of a real nonsymmetric matrix $A$ traditionally have been computed by first reducing $A$ to Hessenberg form $H$ and then computing the eigendecomposition of $H$ by the QR method. But the approach of beginning by reducing the initial matrix to nonsymmetric tridiagonal form is attractive because finding eigenvalues of a tridiagonal matrix is much faster than for a Hessenberg matrix. So, the development of an algorithm for the nonsymmetric tridiagonal eigenvalue problem is therefore a major
topic of research.
In contrast to the symmetric eigenproblem, that is always well-conditioned, the unsymmetric eigenproblem can be effectively very ill-posed. However, there is also a need for a careful study of the sensitivity of eigenvalues of tridiagonals to perturbations.

### 1.3.1 Orthogonal polynomials

From [38, Chapter 7]. Important in applied mathematics are real functions $\phi, \psi, \ldots$ of one real variable and, in particular, the set $\mathcal{P}_{n}$ of polynomials of degree not exceeding $n$. We shall not consider the general integral inner products

$$
(\phi, \psi) \equiv \int_{a}^{b} \omega(x) \phi(x) \psi(x) d x
$$

but go straight to the discrete case

$$
\begin{equation*}
(\phi, \psi) \equiv \sum_{i=1}^{n} \omega_{i} \phi\left(\xi_{i}\right) \psi\left(\xi_{i}\right) . \tag{1.4}
\end{equation*}
$$

To each set of $n$ distinct real numbers $\left\{\xi_{1}, \ldots, \xi_{n}\right\}$ and possible weights $\left\{\omega_{1}, \ldots, \omega_{n}: \omega_{i}>0\right\}$ there corresponds one, and only one, inner product function as defined by (1.4).

Polynomials are rather special functions and for each inner product there is a unique family of monic orthogonal polynomials $\left\{\phi_{0}, \phi_{1}, \ldots, \phi_{n-1}\right\}$; that is, $\phi_{j}$ has degree $j$, leading coefficient 1 and $\left(\phi_{j}, \phi_{k}\right)=0$ for $j \neq k$. This family is the distinguished basis of the inner product space $\mathcal{P}_{n-1}$ enriched with the given inner product.

Tridiagonal matrices come into the picture because there is a remarkable three-term recurrence relation among the $\phi^{\prime} s$; for $j=1,2, \ldots, n-1$ and $\beta_{0}=0$,

$$
\phi_{j}(\xi)=\left(\xi-\alpha_{j}\right) \phi_{j-1}(\xi)-\beta_{j-1}^{2} \phi_{j-2}(\xi) .
$$

Once such a relationship has been guessed, it is straightforward to verify what the $\alpha$ 's and $\beta$ 's must be

$$
\begin{aligned}
\alpha_{j+1} & =\left(\eta \phi_{j}, \phi_{j}\right) /\left(\phi_{j}, \phi_{j}\right), \quad j=0, \ldots, n-1 \\
\beta_{j}^{2} & =\left(\eta \phi_{j-1}, \phi_{j}\right) /\left(\phi_{j-1}, \phi_{j-1}\right), \quad j=1, \ldots, n-1,
\end{aligned}
$$

where $\eta$ denotes the identity function $\eta(\xi) \equiv \xi$. These numbers may be put into an unreduced symmetric tridiagonal matrix $T$ in the obvious way,

$$
T=\left[\begin{array}{ccccc}
\alpha_{1} & \beta_{1} & & & \\
\beta_{1} & \alpha_{2} & \beta_{2} & & \\
& \ddots & \ddots & \ddots & \\
& & \beta_{n-2} & \alpha_{n-1} & \beta_{n-1} \\
& & & \beta_{n-1} & \alpha_{n}
\end{array}\right]
$$

and then, for $j=1, \ldots, n$, we have

$$
\phi_{j}(\xi)=\operatorname{det}\left(\xi I_{j}-T_{j}\right)
$$

where $I_{j}$ denotes the identity matrix or order $j$ and $T_{j}$ denotes the $j^{\text {th }}$ principal submatrix of $T$, that is, $T_{j}$ is the submatrix obtained with the the first $j$ rows and $j$ columns of $T$. Thus the $\xi^{\prime} s$ and the $\omega^{\prime} s$ determine the unique $T$.

The question we pose now is how to determine the $\xi$ 's and the $\omega^{\prime} s$ from a given $T$. In other words, which inner products make the $\phi_{j}, j=0,1, \ldots, n$, mutually orthogonal?

Theorem 1.3.1 Let $T=S \Lambda S^{*}$ be the spectral decomposition of an unreduced symmetric $T$ with $S=\left[\begin{array}{llll}s_{1} & s_{2} & \cdots & s_{n}\end{array}\right]$ and $\Lambda=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right)$. Then the associated inner product of the form (1.4) is given by

$$
\xi_{i}=\lambda_{i}, \quad \omega_{i}=\gamma s_{i 1}^{2}, \quad i=1, \ldots, n,
$$

for any positive $\gamma$. And $\gamma=\sum_{1}^{n} \omega_{i}$.

It is customary to take $\gamma=1$.
The question "how much of this result extends to the unsymmetric case?" has been studied under the tittle "formal orthogonal polynomials", but that is outside the scope of this thesis. See [26].

### 1.3.2 The ordinary and generalized Bessel polynomials

The definition of the generalized Bessel polynomials (GBP) is given in terms of the associated family of differential equations

$$
\begin{equation*}
z^{2} u^{\prime \prime}+(a z+b) u^{\prime}-n(n+a-1) u=0, \quad n \in \mathbb{N}, a, b \in \mathbb{C} . \tag{1.5}
\end{equation*}
$$

It is known [44] that nontrivial polynomial solutions of equation (1.5) exist for every $n \in \mathbb{N}$ and $a, b \in \mathbb{C}$. If $b \neq 0$, there exists a unique polynomial solution $u_{n}^{(a, b)}$ that satisfies the condition $u_{n}^{(a, b)}(0)=1$. The polynomials $u_{n}^{(a, b)}$ defined in this way are precisely the generalized Bessel polynomials. They can be represented in the form

$$
\begin{equation*}
u_{n}^{(a, b)}(z)=\sum_{k=0}^{n} c_{n, k}^{(a, b)} z^{k}, \tag{1.6}
\end{equation*}
$$

where

$$
\begin{equation*}
c_{n, k}^{(a, b)}:=\binom{n}{k} \frac{(n+k+a-2)^{[k]}}{b^{k}}, \quad k=0,1, \ldots, n, \tag{1.7}
\end{equation*}
$$

with

$$
(x)^{[0]}:=1, \quad(x)^{[k]}:=x(x-1) \cdots(x-k+1), \quad k \in \mathbb{N} .
$$

Moreover, it can be shown that they satisfy the following recurrence relations,

$$
\left.\begin{array}{rl}
(n+a-1)(2 n+a-2) u_{n+1}^{(a, b)}(z)= & {\left[(2 n+a)(2 n+a-2) \frac{z}{b}+a-2\right] .} \\
& .(2 n+a-1) u_{n}^{(a, b)}(z)+n(2 n+a) u_{n-1}^{(a, b)}(z)
\end{array}\right)
$$

An important result is that for every $n \in \mathbb{N}, a \in \mathbb{C}$ and $b \in \mathbb{C} \backslash\{0\}$, the GBP $u_{n}^{(a, b)}$ has only simple zeros.

Formulas (1.7) and (1.8) show that the constant $b$ is a scaling factor and almost all authors assume $b=2$; and the case $a \in \mathbb{R}$ is the most investigated one in literature. Taking $a=b=2$ in (1.5) and (1.6) leads to the particular case of the ordinary Bessel polynomials.

Let $z_{n i}^{(a, b)}$ be the zeros of the GBP's $u_{n}^{(a, b)}$ and consider the problem of computing these zeros.

Matrices whose eigenvalues are the zeros $z_{n i}^{(a, b)}$ can be derived either from the coefficients $c_{n, k}^{(a, b)}$ of the polynomials $u_{n}^{(a, b)}$ in (1.6), or from the three-term recurrence relation (1.8). The former procedure leads to the so-called Companion matrices which are Hessenberg matrices. The latter procedure is the classical one, usually adopted to compute the zeros of orthogonal polynomials. The matrix one gets this way is tridiagonal and if denoted by $B_{n}^{(a, b)}$, we have

$$
B_{n}^{(a, b)}=\left[\begin{array}{ccccc}
\alpha_{1}^{(a, b)} & \gamma_{1}^{(a, b)} & & & \\
\beta_{1}^{(a, b)} & \alpha_{2}^{(a, b)} & \gamma_{2}^{(a, b)} & & \\
& \ddots & \ddots & \ddots & \\
& & \beta_{n-2}^{(a, b)} & \alpha_{n-1}^{(a, b)} & \gamma_{n-1}^{(a, b)} \\
& & & \beta_{n-1}^{(a, b)} & \alpha_{n}^{(a, b)}
\end{array}\right]
$$

where

$$
\alpha_{1}^{(a, b)}:=-\frac{a}{b}, \quad \gamma_{1}^{(a, b)}:=-\alpha_{1}^{(a, b)}, \quad \beta_{1}^{(a, b)}:=\frac{\alpha_{1}^{(a, b)}}{a+1}
$$

and

$$
\begin{aligned}
\alpha_{j}^{(a, b)}:=-b \frac{a-2}{(2 j+a-2)(2 j+a-4)}, \quad j=2, \ldots, n, \\
\gamma_{j}^{(a, b)}:=b \frac{j+a-2}{(2 j+a-2)(2 j+a-3)}, \\
\beta_{j}^{(a, b)}:=-b \frac{j}{(2 j+a-1)(2 j+a-2)}, \quad j=2, \ldots, n-1 .
\end{aligned}
$$

Several methods can be adopted to calculate the eigenvalues of $B_{n}^{(a, b)}$.

## Chapter 2

## Representations and measures of sensitivity

This chapter devotes attention to the unsymmetric tridiagonal eigenvalue problem beginning with various representations for tridiagonal matrices. Our first contribution is to provide new measures of eigenvalue sensitivity that exploit the tridiagonal form.

### 2.1 LU factorization

From Higham [27, Chapter 9]. Much insight into Gaussian Elimination (GE) is obtained by expressing it in matrix notation. The strategy of GE is to reduce the initial matrix $A^{1}:=A$ to an upper triangular matrix using elementary row operations. Let $A^{k}=\left(a_{i j}^{k}\right)$, $i, j=1, \ldots, n$, denote the matrix obtained in the $k^{\text {th }}$ stage of GE. The purpose of the $k$ th stage of the elimination is to annihilate the elements below the diagonal in the $k$ th column of $A^{k}$. This is accomplished by the operations

$$
a_{i j}^{k+1}=a_{i j}^{k}-m_{i k} a_{k j}^{k}, \quad i=k+1, \ldots, n, j=k+1, \ldots, n
$$

where $m_{i k}=a_{i k}^{k} / a_{k k}^{k}, i=k+1, \ldots, n$, are the so-called multipliers. At the end of the ( $n-1$ )st stage $A^{n-1}$ is upper triangular.

We can write

$$
A^{k+1}=M_{k} A^{k}:=\left[\begin{array}{ccccc}
I_{k-1} & & & & \\
& 1 & & & \\
& -m_{k+1, k} & 1 & & \\
& -m_{k+2, k} & & \ddots & \\
\vdots & & & \ddots & \\
& -m_{n, k} & & & \\
& &
\end{array}\right] A^{k} .
$$

The matrix $M_{k}$ can be expressed compactly as $M_{k}=I-\boldsymbol{m}_{k} \boldsymbol{e}_{k}^{T}$, where $\boldsymbol{e}_{k}$ is the $k$ th unit vector and $\boldsymbol{m}_{k}=\left[\begin{array}{llllll}0 & \ldots & 0 & m_{k+1, k} & \ldots & m_{n k}\end{array}\right]^{T}$. We have $\boldsymbol{e}_{i}^{T} \boldsymbol{m}_{k}=0$ for $i \leq k$. To invert $M_{k}$, just flip the signs of the multipliers: $M_{k}^{-1}=I+\boldsymbol{m}_{k} \boldsymbol{e}_{k}$. Overall,

$$
M_{n-1} M_{n-2} \ldots M_{1} A=A^{n}=: U
$$

and so

$$
\begin{aligned}
A & =M_{1}^{-1} M_{2}^{-1} \ldots M_{n-1}^{-1} U \\
& =\left(I+\boldsymbol{m}_{1} \boldsymbol{e}_{1}^{T}\right)\left(I+\boldsymbol{m}_{2} \boldsymbol{e}_{2}^{T}\right) \ldots\left(I+\boldsymbol{m}_{n-1} \boldsymbol{e}_{n-1}^{T}\right) U \\
& =\left(I+\sum_{i=1}^{n-1} \boldsymbol{m}_{i} \boldsymbol{e}_{i}^{T}\right) U \\
& =\left[\begin{array}{ccccc}
1 & & & \\
m_{2,1} & 1 & & \\
m_{3,1} & m_{3,2} & \ddots & & \\
\vdots & \vdots & & \ddots & \\
m_{n, 1} & m_{n, 2} & \ldots & m_{n, n-1} & 1
\end{array}\right] U=: L U .
\end{aligned}
$$

The conclusion is that GE computes an LU factorization of $A: A=L U$, where $L$ is unit lower triangular and $U$ is upper triangular.

There are two problems with the method described. First, there is a break-down with division by zero if $a_{k k}^{k}=0$. Second, if we are working in finite precision and some multiplier
$m_{i k}$ is large, there is a possible loss of significance: in the subtraction $a_{i j}^{k}-m_{i k} a_{k j}^{k}$, low-order digits of $a_{i j}^{k+1}$ could be lost. These observations motivate the strategy of partial pivoting.

Let $A_{k}$ denote the $k^{\text {th }}$ leading principal submatrix of $A$, that is $A_{k}$ is the submatrix of the first $k$ rows and first $k$ columns of $A$.

Theorem 2.1.1 There exists a unique $L U$ factorization of $A \in \mathbb{R}^{n \times n}$ if and only if $A_{k}$ is nonsingular for $k=1, \ldots, n-1$. If $A_{k}$ is singular for some $1 \leq k \leq n-1$ the factorization may exist, but if so it is not unique.

### 2.1.1 LU and LDU factorizations of a tridiagonal

A $n \times n$ matrix $H=\left(h_{i j}\right)$ is said to be upper Hessenberg (lower Hessenberg) if $h_{i j}=0$ when $i>j+1(i<j-1)$. It is unreduced, or irreducible, if $h_{i, i-1} \neq 0\left(h_{i-1, i} \neq 0\right), i=2, \ldots, n$. And $H$ is said to be tridiagonal if it is both upper and lower Hessenberg, that is, $h_{i j}=0$ when $|i-j|>1$. A tridiagonal matrix is unreduced if $h_{i, i-1} h_{i-1, i} \neq 0, i=2, \ldots, n$.

Let $C$ be a real nonsymmetric tridiagonal matrix,

$$
C=\left[\begin{array}{ccccc}
a_{1} & c_{1} & & &  \tag{2.1}\\
b_{1} & a_{2} & c_{2} & & \\
& \ddots & \ddots & \ddots & \\
& & b_{n-2} & a_{n-1} & c_{n-1} \\
& & & b_{n-1} & a_{n}
\end{array}\right] \in \mathbb{R}^{n \times n} .
$$

Whenever convenient we will adopte the notation

$$
C=\operatorname{tridiag}(\boldsymbol{b}, \boldsymbol{a}, \boldsymbol{c}),
$$

where $\boldsymbol{a}=\left(a_{k}\right), k=1, \ldots, n$, and $\boldsymbol{b}=\left(b_{k}\right), \boldsymbol{c}=\left(c_{k}\right), k=1, \ldots n-1$.
Assume that $C$ has an LU factorization $C=L U$. Next we will consider details of this factorization as well as the LDU factorization.

## LU factorization

A matrix $B=\left(b_{i j}\right)$ is said to be upper bidiagonal (lower bidiagonal) if $b_{i j}=0$ for $i \neq j$ and $i \neq j-1(i \neq j+1)$.

Assuming there are no break-downs, matrices $L$ and $U$ of the factorization of $C$ are lower bidiagonal and upper bidiagonal, respectively, and will be given by

$$
L=\left[\begin{array}{ccccc}
1 & & & &  \tag{2.2}\\
l_{1} & 1 & & & \\
& l_{2} & 1 & & \\
& & \ddots & \ddots & \\
& & & l_{n-2} & 1 \\
& & & & l_{n-1}
\end{array}\right], \quad U=\left[\begin{array}{llllll}
u_{1} & c_{1} & & & \\
& u_{2} & c_{2} & & & \\
& & u_{3} & c_{3} & & \\
& & & \ddots & \ddots & \\
& & & & u_{n-1} & c_{n-1} \\
& & & & & u_{n}
\end{array}\right],
$$

GE for computing $L$ and $U$ is described by the recurrence relations

$$
\left.u_{1}=a_{1}, \quad \begin{array}{l}
l_{i-1}=b_{i-1} / u_{i-1}  \tag{2.3}\\
u_{i}=a_{i}-l_{i-1} c_{i-1}
\end{array}\right\} i=2, \ldots, n
$$

For $A=\left(a_{i j}\right)$, let $|A|$ denote the matrix whose entries are $\left|a_{i j}\right|$. An error analysis says that, in matrix terms, the computed factors $\tilde{L}$ and $\tilde{U}$ satisfy

$$
A=\tilde{L} \tilde{U}+\delta A, \quad|\delta A| \leq \varepsilon|\tilde{L}||\tilde{U}|
$$

where $\varepsilon$ is the unit roundoff. See Higham [27, p.174]

## LDU factorization

Assume that $u_{i} \neq 0, i=1, \ldots, n-1$. Then the factorization $C=L D U$ is immediate: in (2.2) leave $L$ unchanged, define $D:=\operatorname{diag}\left(d_{1}, \ldots, d_{n}\right)$ with $\boldsymbol{d}=\left(u_{i}\right)$, and redefine

$$
U=\left[\begin{array}{cccccc}
1 & \frac{c_{1}}{u_{1}} & & & &  \tag{2.4}\\
& 1 & \frac{c_{2}}{u_{2}} & & & \\
& & 1 & \frac{c_{3}}{u_{3}} & & \\
& & & \ddots & \ddots & \\
& & & & 1 & \frac{c_{n-1}}{u_{n-1}} \\
& & & & & 1
\end{array}\right]
$$

If $C$ is a symmetric matrix we must have $U=L^{T}$ and thus $C=L D L^{T}$.

### 2.2 Representations of tridiagonals

One of the powerful ideas in matrix computations is factorization. And the power of methods come from representations, especially in the tridiagonal case. Naturally, we can represent a tridiagonal matrix by exhibiting explicitly its entries. Alternatively, we may decompose the matrix into a product of certain factors and give the entries of these factors.

### 2.2.1 Normalization

Consider a tridiagonal matrix $C=\operatorname{tridiag}(\boldsymbol{b}, \boldsymbol{a}, \boldsymbol{c})$ as given in (2.1). If for some $k, c_{k}$ (or $b_{k}$ ) is zero then $C$ can be split into two tridiagonal matrices, say $C_{1}$ of order $k$ and $C_{2}$ or order $n-k$. The eigenvalues and eigenvectors of $C$ can be recovered from those of $C_{1}$ and $C_{2}$ and in this way the computation may be reduced to finding the eigenvalues and eigenvectors of smaller tridiagonal submatrices which are unreduced. So, there is no loss of generality in confining attention to the unreduced case, that is, when $b_{k} c_{k} \neq 0, k=1, \ldots, n-1$.

If $C$ is unreduced, the nullity (dimension of the null space) of $C-\lambda I$ can not exceed 1 . Notice that the minor of the $(1, n)$ element of $C-\lambda I$ is $b_{1} b_{2} \ldots b_{n-1} \neq 0$ and, consequently, $\operatorname{rank}(C-\lambda I) \geq n-1$. So, the dimension of the null space of $C-\lambda I$ is either 0 (when $\lambda$ is not an eigenvalue) or 1 (when $\lambda$ is an eigenvalue).

Fact 2.2.1 Unreduced tridiagonal matrices are nonderogatory.

It is known that the eigenvalues are determined by the diagonal elements and the products of the off-diagonal elements.

In the eigenvalue computation context there is no loss of generality in supposing that tridiagonals are normalized so that all entries in positions $(i, i+1)$ are 1 . So we may assume that $c_{k}=1, k=1, \ldots, n-1$. In fact,

Lemma 2.2.1 Any tridiagonal matrix $C$ that does not split (unreduced) is diagonally similar to a form with 1's above the diagonal.

Proof. If we consider the diagonal matrix

$$
D=\operatorname{diag}\left(1, c_{1}, c_{1} c_{2}, \ldots, c_{1} \cdots c_{n-1}\right),
$$

then the similar matrix $D T D^{-1}$ has the form

$$
\left[\begin{array}{ccccc}
a_{1} & 1 & & & \\
b_{1} c_{1} & a_{2} & 1 & & \\
& \ddots & \ddots & \ddots & \\
& & b_{n-2} c_{n-2} & a_{n-1} & 1 \\
& & & b_{n-1} c_{n-1} & a_{n}
\end{array}\right] \cdot \square
$$

Such tridiagonals will be denoted by $J$. Lets say we will always have

$$
J=\left[\begin{array}{ccccc}
a_{1} & 1 & & &  \tag{2.5}\\
b_{1} & a_{2} & 1 & & \\
& \ddots & \ddots & \ddots & \\
& & b_{n-2} & a_{n-1} & 1 \\
& & & b_{n-1} & a_{n}
\end{array}\right] .
$$

This is called the qd form.

### 2.2.2 Products of bidiagonals

Most, but not all, such $J$ permit triangular factorization

$$
J=L U
$$

where $L$ and $U$ are bidiagonal matrices of a special form that will play a leading role. Precise forms are shown below: $L$ is a lower bidiagonal matrix with 1's in positions $(i, i)$ and $U$ is an upper bidiagonal matrix with 1's in positions $(i, i+1)$,

$$
L=\left[\begin{array}{cccccc}
1 & & & & & \\
l_{1} & 1 & & & & \\
& l_{2} & 1 & & & \\
& & \ddots & \ddots & & \\
& & & l_{n-2} & 1 & \\
& & & & l_{n-1} & 1
\end{array}\right], \quad U=\left[\begin{array}{cccccc}
u_{1} & 1 & & & & \\
& u_{2} & 1 & & & \\
& & u_{3} & 1 & & \\
& & & \ddots & \ddots & \\
& & & & u_{n-1} & 1 \\
& & & & & u_{n}
\end{array}\right] .
$$

The attractive feature here is that, in the $n \times n$ case, because the 1 's need not be represented explicitly the factored form of $J$ requires no more storage than $J$ itself: $L$ and $U$ together are defined by $2 n-1$ parameters, exactly the same degree of freedom as in $J$. Later we will argue that the pair $L, U$ is preferable to $J$ itself.

The pair $L, U$ determines two triangular matrices: first

$$
J=L U=\left[\begin{array}{cccccc}
u_{1} & 1 & & & & \\
l_{1} u_{1} & l_{1}+u_{2} & 1 & & & \\
& l_{2} u_{2} & l_{2}+u_{3} & 1 & & \\
& & \ddots & \ddots & \ddots & \\
& & & l_{n-2} u_{n-2} & l_{n-2}+u_{n-1} & 1 \\
& & & & l_{n-1} u_{n-1} & l_{n-1}+u_{n}
\end{array}\right]
$$

and second

$$
J^{\prime}=U L=\left[\begin{array}{cccccc}
u_{1}+l_{1} & 1 & & & & \\
u_{2} l_{1} & u_{2}+l_{2} & 1 & & & \\
& u_{3} l_{2} & u_{3}+l_{3} & 1 & & \\
& & \ddots & \ddots & \ddots & \\
& & & u_{n-1} l_{n-2} & u_{n-1}+l_{n-1} & 1 \\
& & & & u_{n} l_{n-1} & u_{n}
\end{array}\right] .
$$

Note that both tridiagonals have their superdiagonal entries, that is, entries $(i, i+1)$, equal to 1 . Also note that $J^{\prime}=L^{-1} J L$, that is, $J$ and $J^{\prime}$ are similar.

The inertia of a symmetric matrix $A$ is the triple of integers

$$
\operatorname{Inertia}(A) \equiv(\nu, \zeta, \pi),
$$

where $\nu$ is the number of negative eigenvalues of $A, \zeta$ is the number of zero eigenvalues of $A$ and $\pi$ is the number of positive eigenvalues of $A$.

If $X$ is a nonsingular matrix, we say that $X^{T} A X$ and $A$ are congruent.
If $X$ is orthogonal, then $X^{T} A X$ are similar and so have the same eigenvalues. When $X$ is only nonsingular, $X^{T} A X$ will generally not have the same eigenvalues of $A$, but the next theorem tells us that the two sets of eigenvalues will at least have the same signs.

Theorem 2.2.1 Sylvester's inertia theorem. Let $A$ be symmetric and $X$ be nonsingular. Then $A$ and $X^{T} A X$ have the same inertia.

A proof of this result can be found, for instance, in Demmel [11, p. 202].

In the past most attention has been paid to the positive case: $l_{i}>0$, $i=1, \ldots, n-1, u_{j}>0, j=1, \ldots, n-1$. Remember the following standard results.

Lemma 2.2.2 If $l_{i} u_{i}>0, i=1, \ldots, n-1$, then $J=L U$ is symmetrizable by a diagonal similarity and the number of positive (negative) $u_{i}$ is the number of positive (negative) eigenvalues.

Proof. Consider the diagonal matrix

$$
D=\operatorname{diag}\left(1, \sqrt{l_{1} u_{1}}, \sqrt{l_{1} u_{1} l_{2} u_{2}}, \ldots, \sqrt{l_{1} u_{1} l_{2} u_{2} \cdots l_{n-1} u_{n-1}}\right) .
$$

Then the $i$ th subdiagonal and superdiagonal entries of the tridiagonal matrix $T \equiv D^{-1} J D$ are equal to $\sqrt{l_{i} u_{i}}, i=1, \ldots, n-1$.

Now, notice that

$$
T=D^{-1} J D=D^{-1} L U D=\left(D^{-1} L D\right)\left(D^{-1} U D\right)=L^{\prime} U^{\prime}
$$

where $L^{\prime} \equiv D^{-1} L D$ and $U^{\prime} \equiv D^{-1} U D$ is the $L U$ decomposition of $T$. The symmetry of $T$ permits us to write

$$
T=L^{\prime} D^{\prime} L^{\prime T}
$$

where $D^{\prime}=\operatorname{diag}\left(u_{1}, u_{2}, \ldots, u_{n}\right)$. So, we have that $J$ and $T$ are similar and $T$ and $D^{\prime}$ are congruent. By Sylvester's inertia theorem, $J$ and $D^{\prime}$ have the same inertia. The eigenvalues of $D^{\prime}$ are the values $u_{i}$ and, then, the number of positive (negative) $u_{i}$ is the number of positive (negative) eigenvalues of $J$.

Lemma 2.2.3 If $l_{i} u_{i+1}>0, \quad i=1, \ldots, n-1$, then $J^{\prime}=U L$ is symmetrizable by a diagonal similarity and the number of positive (negative) $u_{i}$ is the number of positive (negative) eigenvalues.

The proof of this lemma is entirely analogous to the proof of the previous lemma.

### 2.2.3 Balancing

Ordinarily, balancing improves the conditioning of the initial matrix, enabling more accurate computation of eigenvectors and eigenvalues. It is an attempt to concentrate any ill-conditioning of the eigenvector matrix into a diagonal scaling.

Let $A=\left(a_{i j}\right) \in \mathbb{C}^{n \times n}$ be a $n \times n$ matrix and let $A^{T}=\left(a_{i j}^{\prime}\right)$ be its transpose. We say that $A$ is balanced if $\left\|\boldsymbol{a}_{i}\right\|_{2}=\left\|\boldsymbol{a}^{\prime}\right\|_{2}, i=1, \ldots, n$, that is, $i^{\text {th }}$ column and $i^{\text {th }}$ row of $A$ have the same norm.

Lemma 2.2.4 Any unreduced tridiagonal matrix $C \in \mathbb{C}^{n \times n}$ may be balanced by a diagonal similarity transformation.

Proof. Note that we now consider complex matrices. Let

$$
C=\operatorname{tridiag}(\boldsymbol{b}, \boldsymbol{a}, \boldsymbol{c})
$$

be an unreduced tridiagonal matrix so that

$$
D \equiv \operatorname{diag}\left(1, \sqrt{\frac{c_{1}}{b_{1}}}, \sqrt{\frac{c_{1} c_{2}}{b_{1} b_{2}}}, \ldots, \sqrt{\frac{c_{1} c_{2} \cdots c_{n-1}}{b_{1} b_{2} \cdots b_{n-1}}}\right)
$$

is defined. Then the similar matrix $D C D^{-1}$ is symmetric,

$$
D C D^{-1}=\left[\begin{array}{ccccc}
a_{1} & \sqrt{b_{1} c_{1}} & & & \\
\sqrt{b_{1} c_{1}} & a_{2} & \sqrt{b_{2} c_{2}} & & \\
& \ddots & \ddots & \ddots & \\
& & \sqrt{b_{n-2} c_{n-2}} & a_{n-1} & \sqrt{b_{n-1} c_{n-1}} \\
& & & \sqrt{b_{n-1} c_{n-1}} & a_{n}
\end{array}\right],
$$

and, thus, balanced.
In practice, balancing only attempts to make the norm of each row equal to the norm of the corresponding column and, usually, can not turn a real nonsymmetric matrix into a real symmetric matrix.

So, any unreduced tridiagonal matrix $C$ may be symmetrizable. If $C$ is real and $b_{i} c_{i}>0$, $i=1, \ldots, n$, then $D C D^{-1}$ is also real and, thus, hermitian. In this case, all eigenvalues of $C$ are real. Otherwise, $D C D^{-1}$ will have complex entries and may have complex eigenvalues.

Symmetrizing is not the only way to balance a matrix. If we consider the diagonal similarity defined by

$$
D \equiv \operatorname{diag}\left(1, \sqrt{\left|\frac{c_{1}}{b_{1}}\right|}, \sqrt{\left|\frac{c_{1} c_{2}}{b_{1} b_{2}}\right|}, \cdots, \sqrt{\left|\frac{c_{1} c_{2} \cdots c_{n-1}}{b_{1} b_{2} \cdots b_{n-1}}\right|}\right)
$$

the matrix $D C D^{-1}$ will also be balanced. It will be equal to

$$
\left[\begin{array}{ccccc}
a_{1} & \mathrm{e}^{\mathrm{i} \varphi_{1}} \sqrt{\left|b_{1} c_{1}\right|} & & &  \tag{6}\\
\mathrm{e}^{\mathrm{i} \theta_{1}} \sqrt{\left|b_{1} c_{1}\right|} & a_{2} & \mathrm{e}^{\mathrm{i} \varphi_{2}} \sqrt{\left|b_{2} c_{2}\right|} & & \\
& \ddots & \ddots & \ddots & \\
& & & & \\
& & \mathrm{e}^{\mathrm{i} \theta_{n-2}} \sqrt{\left|b_{n-2} c_{n-2}\right|} & a_{n-1} & \mathrm{e}^{\mathrm{i} \varphi_{n-1}} \sqrt{\left|b_{n-1} c_{n-1}\right|} \\
& & & \mathrm{e}^{\mathrm{i} \theta_{n-1}} \sqrt{\left|b_{n-1} c_{n-1}\right|} & a_{n}
\end{array}\right]
$$

where $b_{i}=\mathrm{e}^{\mathrm{i} \theta_{i}}\left|b_{i}\right|$ and $c_{i}=\mathrm{e}^{\mathrm{i} \varphi_{i}}\left|c_{i}\right|, i=1, \ldots, n-1$, with $\theta_{i}, \varphi_{i} \in[0,2 \pi[$, and $\mathbf{i}$ is the imaginary unit.

Lemma 2.2.5 Any unreduced tridiagonal matrix $C \in \mathbb{R}^{n \times n}$ is diagonally similar to a form $\Delta T$ where $T$ is symmetric and $\Delta=\operatorname{diag}\left(\delta_{1}, \delta_{2}, \ldots, \delta_{n}\right), \delta_{i}= \pm 1$.

Proof. Consider the diagonal similarity shown in (2.6) and let $B=D C D^{-1}$. Since $C$ is real, $\varphi_{i}$ and $\theta_{i}$ are either $\pi$ or 0 , for all $i$. So, $\mathrm{e}^{\mathrm{i} \theta_{i}}=\operatorname{sign}\left(b_{i}\right)$ and $\mathrm{e}^{\mathbf{i} \varphi_{i}}=\operatorname{sign}\left(c_{i}\right)$. Let

$$
\begin{aligned}
& \delta_{1}=1, \quad \delta_{k}=\delta_{k-1} \operatorname{sign}\left(b_{k-1} c_{k-1}\right), \quad k=2, \ldots, n . \\
& \gamma_{1}=\operatorname{sign}\left(c_{1}\right), \quad \gamma_{k}=\delta_{k} \operatorname{sign}\left(c_{k}\right), \quad k=2, \ldots, n-1 .
\end{aligned}
$$

Notice that $\operatorname{sign}\left(\delta_{k} \gamma_{k}\right)=\operatorname{sign}\left(c_{k}\right), k=1, \ldots, n-1$, and $\operatorname{sign}\left(\delta_{k} \gamma_{k-1}\right)=\operatorname{sign}\left(b_{k-1}\right)$, $k=2, \ldots, n$.

Now we can write $B$ as $B=\Delta T$ where $\Delta=\operatorname{diag}\left(\delta_{1}, \delta_{2}, \ldots, \delta_{n}\right)$, and

$$
\begin{aligned}
& T=\left[\begin{array}{ccccc}
\delta_{1} a_{1} & \gamma_{1} \sqrt{\left|b_{1} c_{1}\right|} & & \\
\gamma_{1} \sqrt{\left|b_{1} c_{1}\right|} & \delta_{2} a_{2} & \gamma_{2} \sqrt{\left|b_{2} c_{2}\right|} & \\
& \ddots & & \ddots & \\
& & \ddots & & \\
& & & \gamma_{n-2} \sqrt{\left|b_{n-2} c_{n-2}\right|} & \delta_{n-1} a_{n-1}
\end{array}\right] \gamma_{n-1} \sqrt{\left|b_{n-1} c_{n-1}\right|} \\
& \text { And } C=D^{-1} \Delta T D . \\
& \\
& \\
&
\end{aligned}
$$

Matrix $\Delta$ is called a signature matrix.
Consider the following $5 \times 5$ example.

Example 2.2.1 We have

$$
\left[\begin{array}{ccccc}
1 & -2 & & & \\
3 & 2 & 2 & & \\
& -1 & 3 & -2 & \\
& & 3 & 4 & 2 \\
& & & -1 & \\
& -1 & & & \\
& & 1 & & \\
& & & -1 & \\
& & & & 1
\end{array}\right]=D^{-1}\left[\begin{array}{ccccc}
1 & & -\sqrt{6} & & \\
-\sqrt{6} & -2 & -\sqrt{2} & & \\
& & -\sqrt{2} & 3 & -\sqrt{6}
\end{array}\right]
$$

where

$$
D=\left[\begin{array}{lllll}
1 & & & & \\
& \sqrt{\frac{2}{3}} & & & \\
& & \sqrt{\frac{4}{3}} & & \\
& & & \sqrt{\frac{8}{9}} & \\
& & & & \sqrt{\frac{16}{9}}
\end{array}\right] . \diamond
$$

### 2.3 Generalized eigenproblem

A generalized eigenvalue problem is an eigenproblem involving more than one matrix. In this section we will only give the definition and a result relating this problem to the standard one.

The standard eigenvalue problem asks for which scalars $\lambda$ the matrix $A-\lambda I$ is singular; these scalars are the eigenvalues. This notion generalizes in several important ways.

If $A$ and $B$ are $m \times n$ matrices, $A-\lambda B$ is called a matrix pencil or just a pencil.
Definition 2.3.1 If $A$ and $B$ are square matrices and $\operatorname{det}(A-\lambda B)$ is not identically zero, the pencil $A-\lambda B$ is called regular. Otherwise it is called singular. When $A-\lambda B$ is regular, $p(\lambda) \equiv \operatorname{det}(A-\lambda B)$ is called the characteristic polynomial of $A-\lambda B$ and the eigenvalues of $A-\lambda B$ are defined to be

1. the roots of $p(\lambda)=0$,
2. $\infty$ (with multiplicity $n-\operatorname{deg} p$ ) if $\operatorname{deg} p<n$.

Matrix pencils arise naturally in many mathematical models of physical systems and the generalized eigenvalue problem is, in principle, more difficult than the standard one. The next fact relates the eigenvalues of a regular pencil $A-\lambda B$ to the eigenvalues of a single matrix.

Fact 2.3.1 Let $A-\lambda B$ be regular. If $B$ is nonsingular, all eigenvalues of $A-\lambda B$ are finite and the same as the eigenvalues of $A B^{-1}$ or $B^{-1} A$. If $B$ is singular, $A-\lambda B$ has eigenvalue $\infty$ with multiplicity $n-\operatorname{rank}(B)$. If $A$ is nonsingular, the eigenvalues of $A-\lambda B$ are the same as the reciprocals of the eigenvalues of $A^{-1} B$ or $B A^{-1}$, where a zero eigenvalue of $A^{-1} B$ corresponds to an infinite eigenvalue of $A-\lambda B$.

For a proof see [11, p. 174].
For a balanced matrix $\Delta T$ the standard eigenvalue problem $(\Delta T-\lambda I) \boldsymbol{v}=0$ is equivalent to the generalized eigenvalue problem

$$
(T-\lambda \Delta) \boldsymbol{v}=0
$$

In fact, we have

$$
(\Delta T-\lambda I) \boldsymbol{v}=0 \Longleftrightarrow \Delta(\Delta T-\lambda I) \boldsymbol{v}=0 \Longleftrightarrow(T-\lambda \Delta) \boldsymbol{v}=0
$$

Both $T$ and $\Delta$ are real and symmetric and, so, we are in the presence of a symmetric generalized eigenvalue problem. If either $T$ or $\Delta$ is positive definite then the eigenvalues must be real. Usually $T$ is not symmetric positive definite. Thus, the case of interest here is when neither $T$ nor $\Delta$ is positive definite (both indefinite) and so there may be complex eigenvalues. See Parlett [38, Chapter 15].

### 2.4 Measures of sensitivity

Nonsymmetric matrices can have a mixture of poorly and well conditioned eigenvalues. Small perturbations in the matrix, such as roundoff errors, can lead to large changes in the eigenvalues.

In this section we present several measures, some new, for the sensitivity of eigenvalues of a tridiagonal matrix. Tridiagonal matrices may be represented as products of bidiagonals in various ways depending on properties such as symmetry and positive definiteness, and different representations lead to different measures of sensitivity. These are relative condition numbers (relative perturbation results) for a simple eigenvalue $\lambda \neq 0$ of products of bidiagonals. The constraint that $\lambda \neq 0$ is always a weakness of a relative approach. When $\lambda=0$ the relative approach looses its meaning and we go back to the absolute approach.

Recall that for $A=\left(a_{i j}\right),|A|$ denotes the matrix whose entries are $\left|a_{i j}\right|$. No use will be made of norms.

### 2.4.1 Balanced product of bidiagonals $\mathrm{T} \Delta=\mathrm{LDU}$

As shown above, any unreduced tridiagonal matrix $C$ may be transformed into a balanced form by a diagonal similarity transformation. So, in this section we will restrict our attention to balanced tridiagonal matrices in the form $B=\Delta T$ (specially relevant for the computation
of eigenvectors) and such that symmetric tridiagonal matrix $T$ admits the product

$$
T=L D L^{T}
$$

where $L=I+\stackrel{\circ}{L}$,

$$
\stackrel{\circ}{L}=\left[\begin{array}{ccccc}
0 & & & & \\
l_{1} & 0 & & & \\
& \ddots & \ddots & & \\
& & l_{n-2} & 0 & \\
& & & l_{n-1} & 0
\end{array}\right]
$$

and $D=\operatorname{diag}\left(d_{1}, \ldots, d_{n}\right)$.
Since $\Delta^{-1} B \Delta=\Delta^{-1} \Delta T \Delta=T \Delta$, matrices $B$ and $T \Delta$ are similar. Thus we will be interested in the factorization of $T \Delta$ that will be

$$
T \Delta=L D U
$$

where $U=\left(I+\stackrel{\circ}{L}^{T}\right) \Delta$. As observed before, when $T$ is real, $\Delta=\operatorname{diag}(1, \pm 1, \ldots, \pm 1)$.
In detail, we are just saying that an unreduced tridiagonal matrix $C$ is always diagonally similar to a form $T \Delta$ where $T$ is symmetric,

$$
C=D_{1}^{-1} B D_{1}=D_{1}^{-1} \Delta T D_{1}=D_{1}^{-1} \Delta T \Delta \Delta^{-1} D_{1}=\left(D_{1}^{-1} \Delta\right)(T \Delta)\left(D_{1}^{-1} \Delta\right)^{-1}
$$

Now, let us concentrate on the factorization $T \Delta=L D U$ for the real case. The offdiagonal entries of $T \Delta$ are $\pm l_{j} d_{j}$. The perturbations will be of the form

$$
l_{i} \rightarrow l_{i}\left(1+\alpha_{i}\right), \quad d_{i} \rightarrow d_{i}\left(1+\beta_{i}\right)
$$

with $\left|\alpha_{i}\right| \leq \eta,\left|\beta_{i}\right| \leq \eta$. In matrix terms,

$$
L \rightarrow L+\delta L, \quad D \rightarrow D+\delta D
$$

with $\delta L=\operatorname{diag}\left(\alpha_{1}, \ldots, \alpha_{n}\right) \stackrel{\circ}{L}$ and $\delta D=\operatorname{diag}\left(\beta_{1}, \ldots, \beta_{n}\right) D$. However the $\alpha_{i}$ and $\beta_{i}$ play no role because we majorize such terms by $\eta$.

Then, element by element, we write

$$
|\delta L| \leq \eta|\check{L}|, \quad|\delta D| \leq \eta|D|, \quad|\delta U| \leq \eta\left|\check{L}^{T} \Delta\right|=\eta\left|\check{L}^{T}\right| .
$$

Matrix $\Delta$ does not change.
Consider the change in a simple eigenvalue $\lambda \neq 0$ satisfying

$$
\begin{equation*}
L D U \boldsymbol{x}=\boldsymbol{x} \lambda, \quad \boldsymbol{y}^{*} L D U=\lambda \boldsymbol{y}^{*}, \quad \boldsymbol{y}^{*} \boldsymbol{x} \neq 0 . \tag{2.7}
\end{equation*}
$$

The perturbed values satisfy

$$
\begin{equation*}
(L+\delta L)(D+\delta D)(U+\delta U)(\boldsymbol{x}+\delta \boldsymbol{x})=(\boldsymbol{x}+\delta \boldsymbol{x})(\lambda+\delta \lambda) . \tag{2.8}
\end{equation*}
$$

If $\eta$ is small enough then, after multiplying out the factors in (2.8) and using (2.7), we obtain

$$
\begin{equation*}
(L D \delta U+L \delta D U+\delta L D U) \boldsymbol{x}+L D U \delta \boldsymbol{x}+\mathcal{O}\left(\eta^{2}\right)=\boldsymbol{x} \delta \lambda+\delta \boldsymbol{x} \lambda+\mathcal{O}\left(\eta^{2}\right) . \tag{2.9}
\end{equation*}
$$

Multiply (2.9) on the left by $\boldsymbol{y}^{*}$ to get

$$
\boldsymbol{y}^{*}(L D \delta U+L \delta D U+\delta L D U) \boldsymbol{x}+\boldsymbol{y}^{*} L D U \delta \boldsymbol{x}=\boldsymbol{y}^{*} \boldsymbol{x} \delta \lambda+\boldsymbol{y}^{*} \delta \boldsymbol{x} \lambda+\mathcal{O}\left(\eta^{2}\right) .
$$

Now use (2.7) to see that $\boldsymbol{y}^{*} L D U \delta \boldsymbol{x}=\boldsymbol{y}^{*} \delta \boldsymbol{x} \lambda$ so that

$$
\begin{equation*}
\delta \lambda \boldsymbol{y}^{*} \boldsymbol{x}=\boldsymbol{y}^{*}(L D \delta U+L \delta D U+\delta L D U) \boldsymbol{x}+\mathcal{O}\left(\eta^{2}\right) . \tag{2.10}
\end{equation*}
$$

There are two ways to proceed. First insert $\delta U, \delta D$, and $\delta L$ and majorize to find

$$
\left|\delta \lambda \| \boldsymbol{y}^{*} \boldsymbol{x}\right| \leq \eta|\boldsymbol{y}|^{T}\left(|L D|\left|\dot{L}^{T}\right|+|L||D||U|+|\Sigma ̊||D U|\right)|\boldsymbol{x}|+\mathcal{O}\left(\eta^{2}\right) .
$$

Use the fact that $L=I+\dot{L}, U=I+\dot{L}^{T} \Delta$ to find

$$
\left|\delta \lambda \| \boldsymbol{y}^{*} \boldsymbol{x}\right| \leq \eta|\boldsymbol{y}|^{T}\left(|D|+3\left|\AA^{\circ} D \check{L}^{T}\right|+2\left|D \grave{L}^{T}\right|+2|亡 ̊ D|\right)|\boldsymbol{x}|+\mathcal{O}\left(\eta^{2}\right) .
$$

Defining $G_{1} \equiv|D|+3\left|\check{L} D \stackrel{\circ}{L}^{T}\right|+2\left|D \grave{L}^{T}\right|+2|\stackrel{\circ}{L} D|$ we have

$$
\frac{|\delta \lambda|}{|\lambda|} \leq \eta \frac{|\boldsymbol{y}|^{T} G_{1}|\boldsymbol{x}|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right||\lambda|}+\mathcal{O}\left(\eta^{2}\right) .
$$

This gives the first relative condition number

$$
\operatorname{relcond}_{1}(\lambda ; L D U):=\frac{|\boldsymbol{y}|^{T} G_{1}|\boldsymbol{x}|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right||\lambda|}
$$

and

$$
G_{1}=\left[\begin{array}{rrrrr}
\left|d_{1}\right| & 2\left|d_{1}\right|\left|l_{1}\right| & & & \\
2\left|d_{1}\right|\left|l_{1}\right| & \left|d_{2}\right|+3\left|d_{1}\right|\left|l_{1}\right|^{2} & 2\left|d_{2}\right|\left|l_{2}\right| & & \\
\ddots & \ddots & & \ddots & \\
& 2\left|d_{n-2}\right|\left|l_{n-2}\right| & \left|d_{n-1}\right|+3\left|d_{n-2}\right|\left|l_{n-2}\right|^{2} & 2\left|d_{n-1}\right|\left|l_{n-1}\right| \\
& & 2\left|d_{n-1}\right|\left|l_{n-1}\right| & \left|d_{n}\right|+3\left|d_{n-1}\right|\left|l_{n-1}\right|^{2}
\end{array}\right]
$$

The second approach avoids the presence of $|\lambda|$ in the denominator. The division by $|\lambda|$ will be done implicitly and, thus, more elegantly. So, return to (2.10) and use the relations (the constraint $\lambda \neq 0$ implies that $U$ is invertible)

$$
\boldsymbol{y}^{*} L D=\lambda \boldsymbol{y}^{*} U^{-1}, \quad \boldsymbol{y}^{*} L=\lambda \boldsymbol{y}^{*} U^{-1} D^{-1}, \quad U \boldsymbol{x}=D^{-1} L^{-1} \lambda \boldsymbol{x} \quad \text { and } \quad D U \boldsymbol{x}=L^{-1} \lambda \boldsymbol{x}
$$

to find

$$
\begin{equation*}
\delta \lambda \boldsymbol{y}^{*} \boldsymbol{x}=\lambda \boldsymbol{y}^{*}\left(U^{-1} \delta U+U^{-1} D^{-1} \delta D D^{-1} L^{-1} \lambda+\delta L L^{-1}\right) \boldsymbol{x}+\mathcal{O}\left(\eta^{2}\right) . \tag{2.11}
\end{equation*}
$$

In order to majorize each term we will use the fact that $\grave{L}$ is a nilpotent matrix. We have $\stackrel{\circ}{L}^{n}=\left(\dot{L}^{T}\right)^{n}=O$, and then

$$
\begin{gather*}
(I+\check{L})^{-1}=\sum_{i=0}^{\infty}(-\stackrel{\circ}{L})^{i}=\sum_{i=0}^{n-1}(-\stackrel{\circ}{L})^{i}  \tag{2.12}\\
\left(I+\stackrel{\circ}{L}^{T}\right)^{-1}=\sum_{i=0}^{\infty}\left(-\stackrel{\circ}{L}^{T}\right)^{i}=\sum_{i=0}^{n-1}\left(-\stackrel{\circ}{L}^{T}\right)^{i} . \tag{2.13}
\end{gather*}
$$

Thus

$$
\begin{equation*}
\left|U^{-1}\right|=\left|\left[\left(I+\dot{L}^{T}\right) \Delta\right]^{-1}\right|=\left|\left(I+\check{L}^{T}\right)^{-1}\right| \leq \sum_{i=0}^{n-1}\left|\dot{L}^{T}\right|^{i}=\left(I-\left|\dot{L}^{T}\right|\right)^{-1} \tag{2.14}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|L^{-1}\right|=\left|(I+\check{L})^{-1}\right| \leq \sum_{i=0}^{n-1}|\stackrel{L}{L}|^{i}=(I-|\check{L}|)^{-1} \tag{2.15}
\end{equation*}
$$

Since $|\delta U| \leq \eta|\check{L}|^{T}$,

$$
\begin{equation*}
\left|U^{-1} \delta U\right| \leq \eta|\check{L}|^{T}\left(I-|\dot{L}|^{T}\right)^{-1} \tag{2.16}
\end{equation*}
$$

and, since $|\delta L| \leq \eta|\Sigma ̊|$,

$$
\begin{equation*}
\left|\delta L L^{-1}\right| \leq \eta \mid 亡(L)(I-|\check{L}|)^{-1} \tag{2.17}
\end{equation*}
$$

Finally, since $|\delta D| \leq \eta|D|$, we have

$$
\begin{align*}
\left|U^{-1} D^{-1} \delta D D^{-1} L^{-1}\right| & \leq \eta\left(I-|\check{L}|^{T}\right)^{-1}\left|D^{-1}\right||D|\left|D^{-1}\right|(I-|\check{L}|)^{-1} \\
& =\eta\left(I-|\check{L}|^{T}\right)^{-1}|D|^{-1}(I-|\check{L}|)^{-1} . \tag{2.18}
\end{align*}
$$

Insert these three values, (2.16), (2.17) and (2.18), in (2.11) to find

$$
\begin{align*}
\left|\frac{\delta \lambda}{\lambda}\right| \leq \frac{\eta|\boldsymbol{y}|^{T}}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}\left[|\check{L}|^{T}\left(I-|\check{L}|^{T}\right)^{-1}+\left(I-|\AA|^{T}\right)^{-1}\right. & |D|^{-1}(I-|\check{L}|)^{-1}|\lambda|+ \\
& \left.+|\check{L}|(I-|\check{L}|)^{-1}\right]|\boldsymbol{x}|+\mathcal{O}\left(\eta^{2}\right) . \tag{2.19}
\end{align*}
$$

It is easy to see that $|\check{L}|^{T}$ and $\left(I-|\check{L}|^{T}\right)^{-1}$ comute, so we can write

$$
\begin{align*}
&\left|\frac{\delta \lambda}{\lambda}\right| \leq \frac{\eta|\boldsymbol{y}|^{T}}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}\left(I-|\check{L}|^{T}\right)^{-1}\left[|\check{L}|^{T}(I-|\check{L}|)+|D|^{-1}|\lambda|+\right. \\
&\left.+\left(I-|\check{L}|^{T}\right)|\check{L}|\right](I-|\check{L}|)^{-1}|\boldsymbol{x}|+\mathcal{O}\left(\eta^{2}\right) . \tag{2.20}
\end{align*}
$$

Now, using the fact that
and that

$$
\left(I-|\check{L}|^{T}\right)^{-1}=\left[(I-|\check{L}|)^{-1}\right]^{T}
$$

we define

$$
G_{2} \equiv|D|^{-1}|\lambda|-2|\check{L}|^{T}|\check{L}|+|\check{L}|^{T}+|\check{L}| .
$$

Note that $G_{2}$ is scaling invariant because it is a sum of scaling invariant matrices. This is important because the relative error is scaling invariant and, then, the formulae should also be scaling invariant. When we multiply the matrix $T$ by a nonzero scalar $\alpha$, the eigenvalue changes to $\alpha \lambda$ and the factor $D$ to $\alpha D$,

$$
\alpha T=L(\alpha D) L^{T} .
$$

Factor $L$ does not change neither does $|D|^{-1}|\lambda|$, i.e., $|\alpha D|^{-1}|\alpha \lambda|=|D|^{-1}|\lambda|$.
So (2.20) can be written as

$$
\frac{|\delta \lambda|}{|\lambda|} \leq \eta \frac{|\boldsymbol{y}|^{T}\left[(I-|\Sigma ̊|)^{-1}\right]^{T} G_{2}(I-|\Sigma ̊|)^{-1}|\boldsymbol{x}|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}+\mathcal{O}\left(\eta^{2}\right)
$$

Our second relative condition number is

$$
\operatorname{relcond}_{2}(\lambda ; L D U):=\frac{v^{T} G_{2} w}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}
$$

where $v$ and $w$ are the solutions of the lower bidiagonal systems

$$
(I-|\check{L}|) v=|\boldsymbol{y}| \quad \text { and } \quad(I-|\check{L}|) w=|\boldsymbol{x}|
$$

and

$$
G_{2}=\left[\begin{array}{ccccc}
|\lambda|\left|d_{1}\right|^{-1}-2 l_{1}^{2} & \left|l_{1}\right| & & & \\
\left|l_{1}\right| & & |\lambda|\left|d_{2}\right|^{-1}-2 l_{2}^{2} & \left|l_{2}\right| & \\
& \ddots & \ddots & \ddots & \\
& & \left|l_{n-2}\right| & & |\lambda|\left|d_{n-1}\right|^{-1}-2 l_{n-1}^{2} \\
& & & \left|l_{n-1}\right| \\
& & & \left|l_{n-1}\right| & |\lambda|\left|d_{n}\right|^{-1}
\end{array}\right] .
$$

Both condition numbers are computable when $\boldsymbol{x}$ and $\boldsymbol{y}$ are known. The former requires only the use of a tridiagonal matrix while the second also requires the solution of two bidiagonal systems with no divisions.

### 2.4.2 Product of bidiagonals $\mathrm{J}=\mathrm{LU}$

Consider a simple eigenvalue of a tridiagonal matrix $J$ represented by $L U$ (relevant for the computation of eigenvalues) where

$$
\begin{equation*}
L=I+\stackrel{\circ}{L} \quad \text { and } \quad U=N+\operatorname{diag}(\boldsymbol{u}) \tag{2.21}
\end{equation*}
$$

with $\boldsymbol{u}=\left(u_{1}, u_{2}, \ldots, u_{n}\right)^{T}$,

$$
\stackrel{L}{L}=\left[\begin{array}{ccccc}
0 & & & & \\
l_{1} & 0 & & & \\
& \ddots & \ddots & & \\
& & l_{n-2} & 0 & \\
& & & l_{n-1} & 0
\end{array}\right] \quad \text { and } \quad N=\left[\begin{array}{ccccc}
0 & 1 & & & \\
& 0 & 1 & & \\
& & \ddots & \ddots & \\
& & & 0 & 1 \\
& & & & 0
\end{array}\right] .
$$

Both $N$ and $\stackrel{\circ}{L}$ are nilpotent matrices, $N^{n}=\dot{L}^{n}=O$.
Without loss of generality we assume that $u_{i} \neq 0, i=1,2, \ldots, n-1$. We also assume that $\lambda \neq 0$.

A simple eigenvalue $\lambda$ has right and left eigenvectors satisfying

$$
\begin{equation*}
L U \boldsymbol{x}=\boldsymbol{x} \lambda, \quad \boldsymbol{y}^{*} L U=\lambda \boldsymbol{y}^{*}, \quad \boldsymbol{y}^{*} \boldsymbol{x} \neq 0 . \tag{2.22}
\end{equation*}
$$

There is no need to normalize $\boldsymbol{x}$ and $\boldsymbol{y}$.
The perturbations of interest are relative:

$$
l_{i} \rightarrow l_{i}\left(1+\alpha_{i}\right), \quad u_{i} \rightarrow u_{i}\left(1+\beta_{i}\right),
$$

with $\left|\alpha_{i}\right| \leq \eta,\left|\beta_{i}\right| \leq \eta$. Early in the analysis we will majorize the perturbations $\left|\alpha_{i}\right|$ and $\left|\beta_{i}\right|$ by $\eta \ll 1$.

In terms of matrices, the perturbations are

$$
L \rightarrow L+\delta L, \quad U \rightarrow U+\delta U
$$

with $\delta L=\operatorname{diag}\left(\alpha_{1}, \ldots, \alpha_{n}\right) \stackrel{\circ}{L}$ and $\delta U=\operatorname{diag}\left(\beta_{1}, \ldots, \beta_{n}\right) \operatorname{diag}(\boldsymbol{u})$.

Then, element by element, we write

$$
\begin{equation*}
|\delta L|=\left|\delta \circ \circ^{L}\right| \leq \eta|\stackrel{\circ}{L}|, \quad|\delta U|=|\delta \operatorname{diag}(\boldsymbol{u})| \leq \eta|\operatorname{diag}(\boldsymbol{u})| . \tag{2.23}
\end{equation*}
$$

We are now ready to make a first order analysis in terms of $\eta$. The perturbed values satisfy

$$
\begin{equation*}
(L+\delta L)(U+\delta U)(\boldsymbol{x}+\delta \boldsymbol{x})=(\boldsymbol{x}+\delta \boldsymbol{x})(\lambda+\delta \lambda) \tag{2.24}
\end{equation*}
$$

If $\eta$ is small enough then, after multiplying out the factors in (2.24) and using (2.22), we obtain

$$
\begin{equation*}
L \delta U \boldsymbol{x}+\delta L U \boldsymbol{x}+L U \delta \boldsymbol{x}+\mathcal{O}\left(\eta^{2}\right)=\boldsymbol{x} \delta \lambda+\delta \boldsymbol{x} \lambda+\mathcal{O}\left(\eta^{2}\right) \tag{2.25}
\end{equation*}
$$

Premultiply by $\boldsymbol{y}^{*}$ to get

$$
\boldsymbol{y}^{*} L \delta U \boldsymbol{x}+\boldsymbol{y}^{*} \delta L U \boldsymbol{x}+\boldsymbol{y}^{*} L U \delta \boldsymbol{x}=\boldsymbol{y}^{*} \boldsymbol{x} \delta \lambda+\boldsymbol{y}^{*} \delta \boldsymbol{x} \lambda+\mathcal{O}\left(\eta^{2}\right)
$$

Now use (2.22) to find

$$
\begin{equation*}
\delta \lambda \boldsymbol{y}^{*} \boldsymbol{x}=\boldsymbol{y}^{*}(L \delta U+\delta L U) \boldsymbol{x}+\mathcal{O}\left(\eta^{2}\right) \tag{2.26}
\end{equation*}
$$

There are two ways to proceed. To first order, after using (2.23),

$$
\begin{equation*}
\left|\delta \lambda \| \boldsymbol{y}^{*} \boldsymbol{x}\right| \leq \eta|\boldsymbol{y}|^{T}(|L||\operatorname{diag}(\boldsymbol{u})|+|\stackrel{\circ}{L}||U|)|\boldsymbol{x}|+\mathcal{O}\left(\eta^{2}\right) \tag{2.27}
\end{equation*}
$$

Use the fact that $L=I+\stackrel{\circ}{L}, U=N+\operatorname{diag}(\boldsymbol{u})$ to find

$$
\left|\delta \lambda \| \boldsymbol{y}^{*} \boldsymbol{x}\right| \leq \eta|\boldsymbol{y}|^{T}(|\operatorname{diag}(\boldsymbol{u})|+2|\stackrel{\circ}{L}||\operatorname{diag}(\boldsymbol{u})|+|\stackrel{\circ}{L}| N)|\boldsymbol{x}|+\mathcal{O}\left(\eta^{2}\right)
$$

Defining $M_{1} \equiv|\operatorname{diag}(\boldsymbol{u})|+2|\stackrel{\circ}{L}||\operatorname{diag}(\boldsymbol{u})|+|\stackrel{\circ}{L}| N$ we have

$$
\frac{|\delta \lambda|}{|\lambda|} \leq \eta \frac{|\boldsymbol{y}|^{T} M_{1}|\boldsymbol{x}|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right||\lambda|}+\mathcal{O}\left(\eta^{2}\right)
$$

This gives the first relative condition number for the form $J=L U$,

$$
\operatorname{relcond}_{1}(\lambda ; L U):=\frac{|\boldsymbol{y}|^{T} M_{1}|\boldsymbol{x}|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right||\lambda|}
$$

and

$$
M_{1}=\left[\begin{array}{rrrr}
\left|u_{1}\right| & & & \\
2\left|l_{1}\right|\left|u_{1}\right| & \left|l_{1}\right|+\left|u_{2}\right| & & \\
\ddots & & \ddots & \\
& 2\left|l_{n-2}\right|\left|u_{n-2}\right| & \left|l_{n-2}\right|+\left|u_{n-1}\right| & \\
& & & 2\left|l_{n-1}\right|\left|u_{n-1}\right| \\
& \left|l_{n-1}\right|+\left|u_{n}\right|
\end{array}\right] .
$$

The second approach tries to bring $\lambda$ into the right side. Return to (2.26) and, from (2.22), use the relations

$$
U \boldsymbol{x}=L^{-1} \boldsymbol{x} \lambda \quad \text { and } \quad \boldsymbol{y}^{*} L=\lambda \boldsymbol{y}^{*} U^{-1}
$$

to find

$$
\begin{equation*}
\delta \lambda \boldsymbol{y}^{*} \boldsymbol{x}=\lambda \boldsymbol{y}^{*}\left(U^{-1} \delta U+\delta L L^{-1}\right) \boldsymbol{x}+\mathcal{O}\left(\eta^{2}\right) . \tag{2.28}
\end{equation*}
$$

After using (2.23), we have

$$
\begin{equation*}
\left|\delta \lambda \| \boldsymbol{y}^{*} \boldsymbol{x}\right| \leq \eta|\lambda||\boldsymbol{y}|^{T}\left(\left|U^{-1}\right||\operatorname{diag}(\boldsymbol{u})|+\left|\Sigma{ }^{\circ}\right|\left|L^{-1}\right|\right)|\boldsymbol{x}|+\mathcal{O}\left(\eta^{2}\right) . \tag{2.29}
\end{equation*}
$$

In order to majorize the term $\left|\AA^{\circ}\right|\left|L^{-1}\right|$ we will use the fact that $\dot{L}$ is a nilpotent matrix. We have $L^{n}=\left(\dot{L}^{T}\right)^{n}=O$, and then

$$
\begin{equation*}
L^{-1}=(I+\stackrel{\circ}{L})^{-1}=\sum_{i=0}^{n-1}(-\stackrel{\circ}{L})^{i} . \tag{2.30}
\end{equation*}
$$

Note that the supports of the powers of $L$ are disjoint. Hence, term by term,

$$
\begin{equation*}
\left|L^{-1}\right| \leq \sum_{i=0}^{n-1}|\check{L}|^{i}=(I-|\dot{L}|)^{-1} . \tag{2.31}
\end{equation*}
$$

Then

$$
\begin{equation*}
|\check{L}|\left|L^{-1}\right| \leq|\check{L}|(I-|\dot{L}|)^{-1}=(I-|\check{L}|)^{-1}-I . \tag{2.32}
\end{equation*}
$$

To majorize $\left|U^{-1}\right||\operatorname{diag}(\boldsymbol{u})|=\left|U^{-1} \operatorname{diag}(\boldsymbol{u})\right|$, note first that $U^{-1} \operatorname{diag}(\boldsymbol{u})=\left(\operatorname{diag}(\boldsymbol{u})^{-1} U\right)^{-1}$. Hence

$$
\operatorname{diag}(\boldsymbol{u})^{-1} U=\left[\begin{array}{ccccc}
1 & u_{1}^{-1} & & & \\
& 1 & u_{2}^{-1} & & \\
& & \ddots & \ddots & \\
& & & 1 & u_{n-1}^{-1} \\
& & & & 1
\end{array}\right]
$$

It is convenient to define the nilpotent matrix

$$
\grave{U}=\left[\begin{array}{ccccc}
0 & u_{1}^{-1} & & & \\
& 0 & u_{2}^{-1} & & \\
& & \ddots & \ddots & \\
& & & 0 & u_{n-1}^{-1} \\
& & & & 0
\end{array}\right]
$$

So, $\dot{U}^{n}=O$ and $\operatorname{diag}(\boldsymbol{u})^{-1} U=I+\dot{U}$. There is no addition or subtraction of scalars in expanding the inverse of $I+{ }^{\circ}$ and from that it follows that

$$
\begin{equation*}
\left(\operatorname{diag}(\boldsymbol{u})^{-1} U\right)^{-1}=(I+\dot{U})^{-1}=\sum_{i=0}^{n-1}(-\dot{U})^{i} \tag{2.33}
\end{equation*}
$$

And, term by term,

$$
\begin{equation*}
\left|U^{-1} \operatorname{diag}(\boldsymbol{u})\right| \leq \sum_{i=0}^{n-1}|\dot{U}|^{i}=(I-|\dot{U}|)^{-1} \tag{2.34}
\end{equation*}
$$

Now insert majorations (2.34) and (2.32) into (2.29) to find

$$
\begin{equation*}
\left|\frac{\delta \lambda}{\lambda}\right|\left|\boldsymbol{y}^{*} \boldsymbol{x}\right| \leq \eta|\boldsymbol{y}|^{T}\left\{(I-|U ْ|)^{-1}+\left[(I-|\dot{L}|)^{-1}-I\right]\right\}|\boldsymbol{x}|+\mathcal{O}\left(\eta^{2}\right) . \tag{2.35}
\end{equation*}
$$

But matrix $(I-|\check{L}|)^{-1}-I$ is not invertible.
We define

$$
M_{2} \equiv(I-|\dot{U}|)^{-1}+\left[(I-|\check{L}|)^{-1}-I\right] .
$$

So (2.35) can be written as

$$
\frac{|\delta \lambda|}{|\lambda|} \leq \eta \frac{|\boldsymbol{y}|^{T} M_{2}|\boldsymbol{x}|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}+\mathcal{O}\left(\eta^{2}\right)
$$

Our second condition number is

$$
\operatorname{relcond}_{2}(\lambda ; L U):=\frac{|\boldsymbol{y}|^{T} M_{2}|\boldsymbol{x}|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}
$$

where

$$
M_{2}=\left[\begin{array}{ccccccc}
1 & \left|u_{1}\right|^{-1} & \left|u_{1} u_{2}\right|^{-1} & \left|u_{1} u_{2} u_{3}\right|^{-1} & \ldots & . & \left|u_{1} u_{2} \ldots u_{n-1}\right|^{-1} \\
\left|l_{1}\right| & 1 & \left|u_{2}\right|^{-1} & \left|u_{2} u_{3}\right|^{-1} & \ldots & . & \left|u_{2} \ldots u_{n-1}\right|^{-1} \\
\left|l_{1} l_{2}\right| & \left|l_{2}\right| & 1 & \left|u_{3}\right|^{-1} & \ldots & . & \left|u_{3} \ldots u_{n-1}\right|^{-1} \\
\vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\
\left|l_{1} \ldots l_{n-2}\right| & \left|l_{2} \ldots l_{n-2}\right| & \left|l_{3} \ldots l_{n-2}\right| & \cdot & \ldots & 1 & \left|u_{n-1}\right|^{-1} \\
\left|l_{1} \ldots l_{n-1}\right| & \left|l_{2} \ldots l_{n-1}\right| & \left|l_{3} \ldots l_{n-1}\right| & \cdot & \ldots & \left|l_{n-1}\right| & 1
\end{array}\right]
$$

We do not need this matrix explicitly. Instead we can solve two bidiagonal systems. Going back to (2.35) we can write

$$
\begin{aligned}
\frac{|\delta \lambda|}{|\lambda|} & \leq \eta \frac{|\boldsymbol{y}|^{T}(I-|\stackrel{\circ}{ }|)^{-1}|\boldsymbol{x}|+|\boldsymbol{y}|^{T}\left[(I-|\stackrel{\circ}{L}|)^{-1}-I\right]|\boldsymbol{x}|}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}+\mathcal{O}\left(\eta^{2}\right) \\
& =\eta \frac{|\boldsymbol{y}|^{T}(v+w)}{\left|\boldsymbol{y}^{*} \boldsymbol{x}\right|}+\mathcal{O}\left(\eta^{2}\right)
\end{aligned}
$$

where $v$ is the solution of the system

$$
(I-|\stackrel{\circ}{U}|) v=|\boldsymbol{x}|
$$

and $w$ satisfies

$$
\left[(I-\mid \circ 口)^{-1}-I\right]|\boldsymbol{x}|=w \Longleftrightarrow(I-|\stackrel{\circ}{L}|)(w+|\boldsymbol{x}|)=|\boldsymbol{x}| .
$$

If we solve the system

$$
(I-|\stackrel{\circ}{L}|) z=|x|
$$

then $w=z-|\boldsymbol{x}|$.
We may be bothered by the fact that $\dot{U}$ is not scaling invariant. This defect is a consequence of the normalization in the representation $L U$ that keeps the superdiagonal entries at the value 1 . In other words the eigenvectors $\boldsymbol{x}$ and $\boldsymbol{y}$ have to change along with $L$ and $U$. In more detail, the matrix $10 L U$ has representation

$$
D^{-1}(10 L U) D=(I+10 \check{L})(10 \operatorname{diag}(\boldsymbol{u})+N)
$$

where $D=\operatorname{diag}\left(1,10^{-1}, 10^{-2}, \ldots, 10^{1-n}\right)$.
The qd representation is not scaling invariant. The factors $L$ and $U$ will change if we multiply the matrix $J$ by a scalar; and the condition number will change too. We need a diagonal scaling.

### 2.4.3 Derivatives from $\Delta \mathrm{T}=\Delta \mathrm{L} \Omega \mathrm{L}^{\mathrm{T}}$

In this section, from Parlett [42], we will describe sensitivity in terms of the generalized singular value decomposition. We will present formulas that state the relative sensitivity of the eigenvalues of $\Delta L \Omega L^{T}$ to $L$ 's entries.

Any real tridiagonal is diagonally similar to $\Delta T$, where $\Delta T$ is real symmetric and $\Delta=\operatorname{diag}\left(\delta_{1}, \delta_{2}, \ldots, \delta_{n}\right), \delta_{i}= \pm 1$. This similarity transformation is an instance of balancing a nonsymmetric matrix. We assume that the eigenvalues $\lambda_{i}$ are distinct but we allow them to be complex. Thus

$$
\begin{align*}
& T s_{i}=\Delta s_{i} \lambda_{i}, \quad i=1, \ldots, n \\
& S=\left[s_{1}, \ldots, s_{n}\right], \text { possibly complex. } \tag{2.36}
\end{align*}
$$

Suppose such matrix $T$ may be written as

$$
\begin{equation*}
T=L \Omega L^{T} \tag{2.37}
\end{equation*}
$$

where $L$ is lower bidiagonal and $\Omega$ is diagonal.

If $T$ is positive definite then $\Omega=I$ and $T=L L^{T}$ is the Cholesky factorization of $T$. If $T$ is indefinite and has the factorization

$$
T=L_{1} D L_{1}^{T},
$$

with $D=\operatorname{diag}\left(d_{1}, \ldots, d_{n}\right)$ and

$$
L_{1}=\left[\begin{array}{ccccc}
1 & & & & \\
l_{1} & 1 & & & \\
& \ddots & \ddots & & \\
& & l_{n-2} & 1 & \\
& & & l_{n-1} & 1
\end{array}\right]
$$

then, defining

$$
D_{1}:=\operatorname{diag}\left(\sqrt{\left|d_{1}\right|}, \ldots, \sqrt{\left|d_{n}\right|}\right) \text { and } \Omega:=\operatorname{diag}\left(\omega_{1}, \ldots, \omega_{n}\right), \quad \omega_{i}=\operatorname{sign}\left(d_{i}\right),
$$

we can obtain the factorization (2.37). We will have

$$
\begin{equation*}
T=L_{1} D_{1} \Omega D_{1} L_{1}^{T}=L \Omega L^{T} \tag{2.38}
\end{equation*}
$$

with $L \equiv L_{1} D_{1}$. This is the closest factorization to the Cholesky factorization that we can get.

A standard normalization for the eigenvectors is

$$
\begin{equation*}
S^{T} \Delta S=I \tag{2.39}
\end{equation*}
$$

Let $\Lambda=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right)$ and denote by $\Lambda^{1 / 2}$ the principal square root of $\Lambda$. That is, if $\lambda_{i}=\rho_{i} \mathrm{e}^{\mathbf{i} 2 \theta_{i}}$ then $\lambda_{i}^{1 / 2}=\rho_{i}^{1 / 2} \mathrm{e}^{\mathbf{i} \theta_{i}},-\frac{\pi}{2}<\theta_{i} \leq \frac{\pi}{2}$ and

$$
\Lambda^{1 / 2}=\operatorname{diag}\left(\lambda_{1}^{1 / 2}, \ldots, \lambda_{n}^{1 / 2}\right)
$$

As in the symmetric case there is an indefinite singular value decomposition of $L$,

$$
\begin{equation*}
L=\Delta S \Lambda^{1 / 2} P^{T} \tag{2.40}
\end{equation*}
$$

defining $P$ which is complex in general. This is a generalization of the familiar singular value decomposition $U \Sigma V^{T}$ of $L$.

Note first that from (2.36), we have

$$
\begin{equation*}
T S=\Delta S \Lambda \Leftrightarrow(\Delta S)^{-1} T=\Lambda S^{-1} \tag{2.41}
\end{equation*}
$$

Also note that

$$
\begin{array}{rlr}
P^{T} \Omega P & =\Lambda^{-1 / 2}(\Delta S)^{-1} L \Omega L^{T}(\Delta S)^{-T} \Lambda^{-1 / 2}, & \text { by }(2.40), \\
& =\Lambda^{-1 / 2}(\Delta S)^{-1} T(\Delta S)^{-T} \Lambda^{-1 / 2}, & \text { by }(2.38), \\
& =\Lambda^{-1 / 2}\left[\Lambda S^{-1}(\Delta S)^{-T}\right] \Lambda^{-1 / 2}, & \text { by }(2.41) \\
& =I, \quad \text { by }(2.39) .
\end{array}
$$

So, $P$ is defined by (2.40) as

$$
P=L^{T}(\Delta S)^{-T} \Lambda^{-1 / 2}
$$

With $\Delta$ and $\Omega$ fixed we will study how $\Lambda$ depends on $L$. Write

$$
L=\left[\begin{array}{ccccc}
\mathfrak{a}_{1} & & & & \\
\mathfrak{b}_{1} & \mathfrak{a}_{2} & & & \\
& \ddots & \ddots & & \\
& & \mathfrak{b}_{\mathfrak{n}-2} & \mathfrak{a}_{\mathfrak{n}-1} & \\
& & & \mathfrak{b}_{\mathfrak{n}-1} & \mathfrak{a}_{\mathfrak{n}}
\end{array}\right]
$$

and let $P:=\left[\boldsymbol{p}_{1}, \ldots, \boldsymbol{p}_{n}\right]$. So, $P^{T} \Omega P=I$ and, from (2.40),

$$
L(\Omega P)=\Delta S \Lambda^{1 / 2} P^{T} \Omega P=\Delta S \Lambda^{1 / 2}
$$

Also, using (2.39),

$$
S^{T} L=(\Delta S)^{-1} L=\Lambda^{1 / 2} P^{T}
$$

Consider a typical singular triple $\left(\lambda, s, \boldsymbol{p}^{T}\right)$

$$
\begin{align*}
& L(\Omega \boldsymbol{p})=(\Delta \boldsymbol{s}) \lambda^{1 / 2}  \tag{2.42}\\
& \boldsymbol{s}^{T} L=\lambda^{1 / 2} \boldsymbol{p}^{T}
\end{align*}
$$

Recall that we use to note $\boldsymbol{p}=\left(p_{k}\right)$ and $\boldsymbol{s}=\left(s_{k}\right), k=1, \ldots, n$.
We can now state the relative sensitivity of the eigenvalues of $\Delta T=\Delta L \Omega L^{T}$ to $L$ 's entries. The expressions may be complex.

Theorem 2.4.1 Let $\Delta=\operatorname{diag}\left(\delta_{1}, \delta_{2}, \ldots, \delta_{n}\right), \delta_{i}= \pm 1$, and $\Omega=\operatorname{diag}\left(\omega_{1}, \ldots, \omega_{n}\right), \omega_{i}= \pm 1$. If $\Delta T$ has distinct eigenvalues and (2.37) holds then, for $\lambda \neq 0$,

$$
\begin{align*}
& \frac{1}{2} \frac{\mathfrak{a}_{j}}{\lambda} \frac{\partial \lambda}{\partial \mathfrak{a}_{j}}=\sum_{k=1}^{j} \delta_{k} s_{k}^{2}-\sum_{m=1}^{j-1} \omega_{m} p_{m}^{2}=\sum_{m=j}^{n} \omega_{m} p_{m}^{2}-\sum_{k=j+1}^{n} \delta_{k} s_{k}^{2},  \tag{2.43}\\
& \frac{1}{2} \frac{\mathfrak{b}_{j}}{\lambda} \frac{\partial \lambda}{\partial \mathfrak{b}_{j}}=\sum_{m=1}^{j}\left[\omega_{m} p_{m}^{2}-\delta_{m} s_{m}^{2}\right]=\sum_{m=j+1}^{n}\left[\delta_{m} s_{m}^{2}-\omega_{m} p_{m}^{2}\right] . \tag{2.44}
\end{align*}
$$

For a proof see [42, p. 430$]$.

### 2.4.4 Derivatives from $\Delta \mathrm{T}=\Delta \mathrm{LDL}^{\mathrm{T}}$

From Z. Wu [64]. In this section we will give theoretical evidence that eigenvalues may be determined to high relative accuracy by the bidiagonals. Consider a tridiagonal matrix written in the balanced form $\Delta T$, where $T$ is symmetric and $\Delta=\operatorname{diag}\left(\delta_{1}, \delta_{2}, \ldots, \delta_{n}\right)$ with $\delta_{i}= \pm 1$. Consider a typical eigenpair $(\lambda, \boldsymbol{v})$ of $\Delta T$ with the normalization $\boldsymbol{v}^{*} \boldsymbol{v}=1$. That is,

$$
\begin{equation*}
\Delta T \boldsymbol{v}=\lambda \boldsymbol{v} \Longleftrightarrow(T-\lambda \Delta) \boldsymbol{v}=0 \tag{2.45}
\end{equation*}
$$

Lemma 2.4.1 If $\Im \lambda \neq 0$ then $\boldsymbol{v}^{*} \Delta \boldsymbol{v}=\boldsymbol{v}^{*} T \boldsymbol{v}=0$

Proof. Since $(T-\lambda \Delta) \boldsymbol{v}=0$, we have

$$
\begin{equation*}
\boldsymbol{v}^{*}(T-\lambda \Delta) \boldsymbol{v}=0 . \tag{2.46}
\end{equation*}
$$

Also, since $\boldsymbol{v}^{*} T \boldsymbol{v}$ is a scalar and $T$ is real and symmetric, we have

$$
\overline{\boldsymbol{v}^{*} T \boldsymbol{v}}=\boldsymbol{v}^{T} T \overline{\boldsymbol{v}}=\left(\boldsymbol{v}^{T} T \overline{\boldsymbol{v}}\right)^{T}=\boldsymbol{v}^{*} T \boldsymbol{v}
$$

Using the same arguments, we have that $\boldsymbol{v}^{*} \Delta \boldsymbol{v}$ is also real. So, from (2.46),

$$
\lambda \boldsymbol{v}^{*} \Delta \boldsymbol{v}=\boldsymbol{v}^{*} T \boldsymbol{v}
$$

that is, $\lambda \boldsymbol{v}^{*} \Delta \boldsymbol{v}$ is also real. Let $\lambda=a+b \mathbf{i}$, where $a=\Re \lambda$ and $b=\Im \lambda$. Then

$$
\lambda \boldsymbol{v}^{*} \Delta \boldsymbol{v}=a \boldsymbol{v}^{*} \Delta \boldsymbol{v}+\left(b \boldsymbol{v}^{*} \Delta \boldsymbol{v}\right) \mathbf{i}
$$

and, therefore, if $b \neq 0$, we must have $\boldsymbol{v}^{*} \Delta \boldsymbol{v}=0$.
First consider the two following basic facts.
Fact 2.4.1 Let $V=\left[\boldsymbol{v}_{1}, \boldsymbol{v}_{2}, \ldots, \boldsymbol{v}_{n}\right]$ be the matrix of eigenvectors of $\Delta T$. If $\lambda_{i} \neq \lambda_{j}, i \neq j$, then $V$ is invertible.

The proof of this fact rests on the invertibility of the Vandermonde matrix (eigenvectors corresponding to distinct eigenvalues are linearly independent).

Fact 2.4.2 If $\lambda_{i} \neq \lambda_{j}$ for $i \neq j$, then $\boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{j}=0$.

Proof. From (2.45), suppose that $\lambda_{i} \neq \lambda_{j}$ and vectors $\boldsymbol{v}_{i}$ and $\boldsymbol{v}_{j}$ satisfy

$$
T \boldsymbol{v}_{j}=\Delta \boldsymbol{v}_{j} \lambda_{j} \quad \text { and } \quad T \boldsymbol{v}_{i}=\Delta \boldsymbol{v}_{i} \lambda_{i}
$$

Then

$$
\begin{equation*}
\boldsymbol{v}_{i}^{T} T \boldsymbol{v}_{j}=\boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{j} \lambda_{j} \quad \text { and } \quad \boldsymbol{v}_{j}^{T} T \boldsymbol{v}_{i}=\boldsymbol{v}_{j}^{T} \Delta \boldsymbol{v}_{i} \lambda_{i} \tag{2.47}
\end{equation*}
$$

Note that, since we are dealing with scalars and $T$ is symmetric,

$$
\boldsymbol{v}_{i}^{T} T \boldsymbol{v}_{j}=\left(\boldsymbol{v}_{i}^{T} T \boldsymbol{v}_{j}\right)^{T}=\boldsymbol{v}_{j}^{T} T \boldsymbol{v}_{i} \quad \text { and } \quad \boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{j}=\left(\boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{j}\right)^{T}=\boldsymbol{v}_{j}^{T} \Delta \boldsymbol{v}_{i}
$$

Now, subtracting in (2.47), we get

$$
0=\boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{j}\left(\lambda_{j}-\lambda_{i}\right)
$$

and, since $\lambda_{i} \neq \lambda_{j}$, we obtain that

$$
\boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{j}=0
$$

Lemma 2.4.2 Let $\boldsymbol{v}_{i}, i=1, \ldots, n$ be eigenvectors of $\Delta T$. Assume that $\lambda_{i} \neq \lambda_{j}$ for $i \neq j$. Then $\boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{i} \neq 0, i=1, \ldots, n$.

Proof. Using both facts above, we have that $V^{T} \Delta V$ is a diagonal matrix with diagonal entries $\boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{i}, i=1, \ldots, n$. Thus

$$
\operatorname{det}\left(V^{T} \Delta V\right)=\prod_{i=1}^{n} \boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{i}
$$

Since both $V$ and $\Delta$ are invertible, we must have $\operatorname{det}\left(V^{T} \Delta V\right) \neq 0$ and, therefore,

$$
\boldsymbol{v}_{i}^{T} \Delta \boldsymbol{v}_{i} \neq 0, \quad i=1, \ldots, n
$$

Theorem 2.4.2 Assume that $T$ admits the factorization $T=L D L^{T}$, where $L$ is unit bidiagonal with subdiagonal elements $l_{1}, l_{2}, \ldots, l_{n}$ and $D=\operatorname{diag}\left(d_{1}, d_{2}, \ldots, d_{n}\right)$. Suppose that all the eigenvalues of $\Delta T$ are distinct. Then for $\lambda \neq 0$ and corresponding eigenvector $\boldsymbol{v}=\left(v_{j}\right)$ we have, for $k=1, \ldots, n$,

$$
\frac{d_{k}}{\lambda} \frac{\partial \lambda}{\partial d_{k}}=\frac{\left(\boldsymbol{v}^{T} L \boldsymbol{e}_{k}\right)\left(\boldsymbol{e}_{k}^{T} L^{-1} \Delta \boldsymbol{v}\right)}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}}
$$

and, for $k=1, \ldots, n-1$,

$$
\frac{1}{\lambda} \frac{\partial \lambda}{\partial l_{k}}=\frac{2 v_{k+1}\left(\boldsymbol{e}_{k}^{T} L^{-1} \Delta \boldsymbol{v}\right)}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}}
$$

And

$$
\begin{aligned}
& \boldsymbol{e}_{k}^{T} L^{-1} \Delta \boldsymbol{v}=\delta_{k} v_{k}-\delta_{k-1} v_{k-1} l_{k-1}+(-1)^{2} \delta_{k-2} v_{k-2} l_{k-1} l_{k-2}+\ldots+(-1)^{k-1} \delta_{1} v_{1} l_{k-1} l_{k-2} \ldots l_{1}, \\
& k=1, \ldots, n, \\
& \boldsymbol{v}^{T} L \boldsymbol{e}_{k}=v_{k}+v_{k+1} l_{k}, \quad k=1, \ldots, n-1, \\
& \boldsymbol{v}^{T} L \boldsymbol{e}_{n}=v_{n} .
\end{aligned}
$$

Proof. Differentiate

$$
\begin{equation*}
T \boldsymbol{v}=\lambda \Delta \boldsymbol{v} \tag{2.48}
\end{equation*}
$$

with respect to $d_{k}$ to obtain

$$
\frac{\partial}{\partial d_{k}}(T \boldsymbol{v})=\frac{\partial}{\partial d_{k}}(\lambda \Delta \boldsymbol{v}) .
$$

Since we have $T=L D L^{T}$ it follows that

$$
\begin{equation*}
\frac{\partial\left(L D L^{T}\right)}{\partial d_{k}} \boldsymbol{v}+T \frac{\partial \boldsymbol{v}}{\partial d_{k}}=\frac{\partial \lambda}{\partial d_{k}} \Delta \boldsymbol{v}+\lambda \frac{\partial(\Delta \boldsymbol{v})}{\partial d_{k}} \tag{2.49}
\end{equation*}
$$

Since $\frac{\partial D}{\partial d_{k}}=\boldsymbol{e}_{k} \boldsymbol{e}_{k}^{T}, \frac{\partial L}{\partial d_{k}}=\frac{\partial L^{T}}{\partial d_{k}}=O$ and $\lambda \frac{\partial(\Delta \boldsymbol{v})}{\partial d_{k}}=\lambda \Delta \frac{\partial \boldsymbol{v}}{\partial d_{k}}$, from (2.49) we get

$$
L \boldsymbol{e}_{k} \boldsymbol{e}_{k}^{T} L^{T} \boldsymbol{v}+(T-\lambda \Delta) \frac{\partial \boldsymbol{v}}{\partial d_{k}}=\frac{\partial \lambda}{\partial d_{k}} \Delta \boldsymbol{v}
$$

Now, multiplying by $\boldsymbol{v}^{T}$ on the left,

$$
\boldsymbol{v}^{T} L \boldsymbol{e}_{k} \boldsymbol{e}_{k}^{T} L^{T} \boldsymbol{v}+\boldsymbol{v}^{T}(T-\lambda \Delta) \frac{\partial \boldsymbol{v}}{\partial d_{k}}=\frac{\partial \lambda}{\partial d_{k}} \boldsymbol{v}^{T} \Delta \boldsymbol{v}
$$

From this we obtain

$$
\frac{\partial \lambda}{\partial d_{k}}=\frac{\left(\boldsymbol{v}^{T} L \boldsymbol{e}_{k}\right)\left(\boldsymbol{e}_{k}^{T} L^{T} \boldsymbol{v}\right)}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}}
$$

because $\boldsymbol{v}^{T}(T-\lambda \Delta)=[(T-\lambda \Delta) \boldsymbol{v}]^{T}=0$, since $T$ and $\Delta$ are symmetric, and $\boldsymbol{v}^{T} \Delta \boldsymbol{v} \neq 0$ by lemma (2.4.2). And multiplying both sides by $d_{k}$,

$$
d_{k} \frac{\partial \lambda}{\partial d_{k}}=\frac{\left(\boldsymbol{v}^{T} L \boldsymbol{e}_{k}\right)\left(d_{k} \boldsymbol{e}_{k}^{T} L^{T} \boldsymbol{v}\right)}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}}
$$

To make $\lambda$ appear in the denominator, note that, using (2.48) we can write

$$
\begin{equation*}
d_{k} \boldsymbol{e}_{k}^{T} L^{T} \boldsymbol{v}=\boldsymbol{e}_{k}^{T} D L^{T} \boldsymbol{v}=\boldsymbol{e}_{k}^{T} L^{-1} L D L^{T} \boldsymbol{v}=\boldsymbol{e}_{k}^{T} L^{-1} T \boldsymbol{v}=\boldsymbol{e}_{k} L^{-1} \Delta \boldsymbol{v} \lambda \tag{2.50}
\end{equation*}
$$

Then when $\lambda \neq 0$,

$$
\frac{d_{k}}{\lambda} \frac{\partial \lambda}{\partial d_{k}}=\frac{\left(\boldsymbol{v}^{T} L \boldsymbol{e}_{k}\right)\left(\boldsymbol{e}_{k}^{T} L^{-1} \Delta \boldsymbol{v}\right)}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}}
$$

The second relation is obtained similarly. Differentiating (2.48) with respect to $l_{k}$ we obtain

$$
\begin{equation*}
\frac{\partial\left(L D L^{T}\right)}{\partial l_{k}} \boldsymbol{v}+T \frac{\partial \boldsymbol{v}}{\partial l_{k}}=\frac{\partial \lambda}{\partial l_{k}} \Delta \boldsymbol{v}+\lambda \frac{\partial(\Delta \boldsymbol{v})}{\partial l_{k}} . \tag{2.51}
\end{equation*}
$$

Since $\frac{\partial D}{\partial l_{k}}=O, \frac{\partial L}{\partial l_{k}}=\boldsymbol{e}_{k+1} \boldsymbol{e}_{k}^{T}, \frac{\partial L^{T}}{\partial l_{k}}=\boldsymbol{e}_{k} \boldsymbol{e}_{k+1}^{T}$ and $\lambda \frac{\partial(\Delta \boldsymbol{v})}{\partial d_{k}}=\lambda \Delta \frac{\partial \boldsymbol{v}}{\partial d_{k}}$, it follows that

$$
\boldsymbol{e}_{k+1} \boldsymbol{e}_{k}^{T} D L^{T} \boldsymbol{v}+L D \boldsymbol{e}_{k} \boldsymbol{e}_{k+1}^{T} \boldsymbol{v}+(T-\lambda \Delta) \frac{\partial \boldsymbol{v}}{\partial l_{k}}=\frac{\partial \lambda}{\partial l_{k}} \Delta \boldsymbol{v}
$$

Again, pre-multiplying by $\boldsymbol{v}^{T}$ and using the same arguments as before,

$$
\begin{aligned}
\frac{\partial \lambda}{\partial l_{k}} & =\frac{\boldsymbol{v}^{T} \boldsymbol{e}_{k+1}\left(\boldsymbol{e}_{k}^{T} D L^{T} \boldsymbol{v}\right)+\left(\boldsymbol{v}^{T} L D \boldsymbol{e}_{k}\right) \boldsymbol{e}_{k+1}^{T} \boldsymbol{v}}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}} \\
& =\frac{v_{k+1}\left(\boldsymbol{e}_{k}^{T} D L^{T} \boldsymbol{v}\right)+\left(\boldsymbol{e}_{k}^{T} D L^{T} \boldsymbol{v}\right)^{T} v_{k+1}}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}}
\end{aligned}
$$

Using again (2.50) it follows

$$
\frac{1}{\lambda} \frac{\partial \lambda}{\partial l_{k}}=\frac{2 v_{k+1}\left(\boldsymbol{e}_{k}^{T} L^{-1} \Delta \boldsymbol{v}\right)}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}}
$$

Finally, the result

$$
\begin{aligned}
& \boldsymbol{e}_{k}^{T} L^{-1} \Delta \boldsymbol{v}=\delta_{k} v_{k}-\delta_{k-1} v_{k-1} l_{k-1}+(-1)^{2} \delta_{k-2} v_{k-2} l_{k-1} l_{k-2}+\ldots+(-1)^{k-1} \delta_{1} v_{1} l_{k-1} l_{k-2} \ldots l_{1}, \\
& k=1, \ldots, n, \\
& \boldsymbol{v}^{T} L \boldsymbol{e}_{k}=v_{k}+v_{k+1} l_{k}, \quad k=1, \ldots, n-1, \\
& \boldsymbol{v}^{T} L \boldsymbol{e}_{n}=v_{n} .
\end{aligned}
$$

follows from direct calculations.
Notice that $d_{k}$, like $\lambda$, is proportional to $\|T\|$ while $l_{k}$ is not. That is, if we multiply $T$ by $\mu \neq 0$, then $L(\mu D) L^{T}=\mu T$. Consequently, it is natural to consider relative changes to $d_{k}$ but absolute changes to $l_{k}$. In particular, when $k=1$,

$$
\begin{aligned}
\frac{d_{1}}{\lambda} \frac{\partial \lambda}{\partial d_{1}} & =\frac{\left(v_{1}+v_{2} l_{1}\right) \delta_{1} v_{1}}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}} \\
\frac{1}{\lambda} \frac{\partial \lambda}{\partial l_{1}} & =\frac{2 v_{2} \delta_{1} v_{1}}{\boldsymbol{v}^{T} \Delta \boldsymbol{v}}
\end{aligned}
$$

Theorem (2.4.2) permits us to define the following condition numbers

$$
\operatorname{relcond}\left(\lambda, d_{k} ; \Delta L D L^{T}\right):=\frac{\left|\boldsymbol{v}^{T} L \boldsymbol{e}_{k}\right|\left|\boldsymbol{e}_{k}^{T} L^{-1} \Delta \boldsymbol{v}\right|}{\left|\boldsymbol{v}^{T} \Delta \boldsymbol{v}\right|}
$$

and

$$
\operatorname{relcond}\left(\lambda, l_{k} ; \Delta L D L^{T}\right):=\frac{2\left|v_{k+1}\right|\left|\boldsymbol{e}_{k}^{T} L^{-1} \Delta \boldsymbol{v}\right|}{\left|\boldsymbol{v}^{T} \Delta \boldsymbol{v}\right|}
$$

A comparative study of all these measures of sensitivity we just presented with the more general ones would be useful and is part of the next stage of our study. Classical BauerFike theorem states one condition number for all the eigenvalues and must certainly be too crude. Wilkinson's condition number gives one condition number,

$$
\frac{\left\|\boldsymbol{y}^{*}\right\|\|\boldsymbol{x}\|}{\left|\lambda \| \boldsymbol{y}^{*} \boldsymbol{x}\right|},
$$

for each eigenvalue but must also be pessimistic. It is still too general. The condition numbers we gave are more refined because they exploit the tridiagonal form. The hope is that they will be realistic estimates of sensitivity, not just bounds.

## Chapter 3

## LR and dqds

The focus of this chapter is the description of LR and qds algorithms showing the connection between them. The dqds version demands a little more of arithmetic effort than qds but has compensating advantages - it enjoys the property of mixed high relative stability.

### 3.1 LR algorithm

For more than forty years the standard algorithm for calculating eigenvalues has been the QR algorithm of Francis [18]. But historically the LR algorithm preceded the QR algorithm. The whole field had its genesis in 1957 with Rutishauser's quotient-difference algorithm [46, 47], which Rutishauser then generalized to the LR algorithm [48]. The QR algorithm followed shortly thereafter. Surprisingly, the quotient-difference algorithm has had a recent revival. In 1994, Fernando and Parlett [16] introduced new versions for finding singular values of bidiagonal matrices and eigenvalues of symmetric tridiagonal matrices but it may also be used in the unsymmetric tridiagonal case.

Contrary to QR, the LR algorithm preserves bandwith and thus it is one of the most efficient methods for calculating all the eigenvalues of a nonsymmetric tridiagonal matrix. Most of the improvements that have been incorporated into QR over the years such as implicit double-shift iterations, deflation, splitting, and arbitrary shifts, can also be used in
the context of LR iteration.

## LR algorithm

The LR transformation gives a reduction of a general matrix to triangular form by means of non-unitary similarity transformations. It is based on the LU decomposition of a matrix. Under mild conditions a square matrix admits triangular factorization, $A=L R$, where $L$ is unit lower-triangular and $R$ is upper triangular.

The method bases essentially on the fact that by starting with the given matrix $A=: A_{1}$, the infinite sequence of similar matrices $A_{1}, A_{2}, A_{3}, \ldots$, generated by

## LR algorithm

$$
\begin{align*}
& A_{1}=A \\
& \text { for } i=1,2, \ldots  \tag{3.1}\\
& \quad \text { Factor } A_{i}=L_{i} R_{i} \quad \text { (LU factorization) } \\
& \quad A_{i+1}=R_{i} L_{i} \\
& \text { end }
\end{align*}
$$

converges to a triangular or diagonal matrix $A_{\text {inf }}$, under reasonable assumptions.
In words, we decompose $A_{i}$ and multiply the factors in reverse order to obtain a matrix $A_{i+1}$ that is similar to $A$. This process is repeated until convergence occurs. In fact,

$$
A_{i+1}=R_{i} L_{i}=L_{i}^{-1}\left(L_{i} R_{i}\right) L_{i}=L_{i}^{-1} A_{i} L_{i}
$$

and then, by induction, $A_{i+1}$ is similar to $A_{1}$.
If $A$ is symmetric and positive definite, and if the decomposition of $A_{i}$ into $L_{i} R_{i}$ is such that $L_{i}$ is lower triangular and matrix $R_{i}$ is the transpose of $L_{i}$ for any $i$ (Cholesky decomposition), then $\lim _{i \rightarrow \infty} A_{i}$ exists and is diagonal. A proof of this is given in Rutishauser [48].

But not every matrix has an LU decomposition and difficulties arise when we are close to a matrix that has no LU decomposition. The LR algorithm may not be stable. Stability
can be improved markedly by the introduction of pivoting but this destroys the bandwith structure. Wilkinson [60] uses the LR method with row and column interchanges to produce an algorithm which is stable and requires no more arithmetic than the basic LR. In practice, this variation seems both fast and accurate but no convergence proof has been published yet nor is it apparent how to adapt the technique to give, economically, the complex eigenvalues of real matrices.

### 3.2 QR algorithm

A standard approach in computing the eigenvalues of a general square matrix is to reduce the matrix first to Hessenberg form by a sequence of orthogonal transformations, and then to determine the eigenvalues of the Hessenberg matrix through an iterative process known as the QR algorithm. The reduction to Hessenberg form requires $\mathcal{O}\left(n^{3}\right)$ operations, where $n$ is the order of the matrix, and the subsequent iterative phase also requires $\mathcal{O}\left(n^{3}\right)$ operations. The function eig of Matlab uses this scheme to compute all of the eigenvalues and eigenvectors of a general matrix.

If the original matrix is symmetric, then the symmetry can be preserved in the initial reduction, so the reduced matrix is tridiagonal. Although the reduction to tridiagonal form still requires $\mathcal{O}\left(n^{3}\right)$ operations, the subsequent iterations preserve the tridiagonal form and, hence, are much less expensive, so that the total cost of the iterative phase is reduced to $\mathcal{O}\left(n^{2}\right)$ operations.

The attractively low operation count obtained when iterating with a tridiagonal matrix suggests that the tridiagonal form would be extremely beneficial in the nonsymmetric case as well. Such an approach presents two difficulties, however. First, QR iteration does not preserve the structure of a nonsymmetric tridiagonal matrix. This problem can be overcome by using LR iteration, which preserves the tridiagonal form. Second, it is difficult to reduce a nonsymmetric matrix to tridiagonal form by similarity in a numerically stable manner. But promising methods to improve the stability of this reduction have been developed recently. See Sidje [50]. See also Dongarra [13] and Geist [20].

## QR algorithm

As implemented in the late 1950's and early 1960's, LR proved insufficiently reliable and was displaced by the QR algorithm in the mid 1960's. The QR algorithm, developed by Francis [18], is closely related to the LR algorithm. Based on the use of unitary transformations, the QR algorithm avoids the possible instability of the original LR transformation but it is slower. However, in many respects, this has proved to be the most effective of known methods for the solution of the general algebraic eigenvalue problem. In comparison with current reliable methods, the QR technique is fast and highly satisfactory.

## QR algorithm

$$
\begin{aligned}
& A_{1}=A \\
& \text { for } i=1,2, \ldots \\
& \quad \text { Factor } A_{i}=Q_{i} R_{i} \quad \text { (QR factorization) } \\
& \quad A_{i+1}=R_{i} Q_{i} \\
& \text { end }
\end{aligned}
$$

Orthogonal similarities give stability to the QR algorithm and preserve symmetry as well as the Hessenberg form; thus, the symmetric tridiagonal form is also preserved.

Lemma 3.2.1 Let $A=Q R$ be the QR decomposition of $A$ and

$$
\hat{A}=Q^{T} A Q
$$

be its QR transformation. If $A$ is tridiagonal, symmetric and non-singular, then $\hat{A}$ is also tridiagonal, symmetric and non-singular.

The QR algorithm does not preserve the tridiagonal form in the non-symmetric case. If $A$ is non-symmetric tridiagonal, then $\hat{A}$ will be upper Hessenberg. But it never breaks down.

Convergence of the basic QR algorithm may be too poor for pratical computations. Convergence rates depend on the ratios of eigenvalues and to speed up convergence, shifts of origin should be incorporated.

## QR algorithm with a shift

$$
\begin{aligned}
& A_{1}=A \\
& \text { for } i=1,2, \ldots \\
& \quad \text { Choose a shift } \sigma_{i} \\
& \quad \text { Factor } A_{i}-\sigma_{i} I=Q_{i} R_{i} \quad \text { (QR factorization) } \\
& \quad A_{i+1}=Q_{i} R_{i}+\sigma_{i} I \\
& \text { end }
\end{aligned}
$$

Ideally we want to shift by the eigenvalues. If $\sigma_{i}$ is an exact eigenvalue of $A_{i}$, then the QR iteration converges in one step. See Demmel [11, p.162]. So, shifts $\sigma_{i}$ close to eigenvalues hasten convergence. Note that near convergence to a real eigenvalue, the last diagonal element of $A_{i}=\left(a_{k j}^{(i)}\right)$ is close to that eigenvalue, so $\sigma_{i}=a_{n n}^{(i)}$ is a good choice for a shift.

### 3.3 Shifted LR algorithm

So, in order to accelerate convergence (and also to be able to treat indefinite symmetric matrices) the simple decomposition-recombination procedure (3.1) of LR algorithm must be modified in the following way:

## LR algorithm with a shift

$$
\begin{aligned}
& A_{1}=A \\
& \text { for } i=1,2, \ldots \\
& \quad \text { Choose a shift } \sigma_{i} \\
& \quad \text { Factor } A_{i}-\sigma_{i} I=L_{i} R_{i} \quad \text { (LU factorization) } \\
& A_{i+1}=R_{i} L_{i}+\sigma_{i} I \\
& \text { end }
\end{aligned}
$$

Notice that matrices $A_{i}$ and $A_{i+1}$ continue to be similar,

$$
\begin{equation*}
A_{i+1}=R_{i} L_{i}+\sigma_{i} I=L_{i}^{-1}\left(L_{i} R_{i}+\sigma_{i} I\right) L_{i}=L_{i}^{-1} A_{i} L_{i}, \tag{3.2}
\end{equation*}
$$

and we call the transformation shift restoring.
By proper choice of $\sigma_{i}$, the last diagonal element $a_{n n}^{(i)}$ of $A_{i}$ will, in general, converge to the smallest eigenvalue of $A$. If $\sigma_{i}$ is chosen to be close to an eigenvalue of $A_{i}$ then $A_{i}-\sigma_{i} I$ has an eigenvalue close to zero.

There are no simple expressions for convergence rates in shifted LR algorithm because it depends on all shifts $\sigma_{i}$ 's.

Lemma 3.3.1 If $A-\sigma I$ is singular and $A-\sigma I=L R$ is the LU decomposition of $A-\sigma I$, then the transformation $\hat{A}=L^{-1} A L$ deflates in one step.

Proof. Let $A$ be of order $n$. Since $A-\sigma I$ is singular, if $L=\left(l_{i j}\right)$ and $R=\left(r_{i j}\right)$ are the factors of its LU factorization, then, by theorem 2.1.1 (page 27), we must have $r_{n n}=0$. Therefore,

$$
L^{-1} A L=L^{-1}(L R+\sigma I) L=R L+\sigma I .
$$

But $\boldsymbol{e}_{n}^{T} R=\mathbf{0}^{T}$ and then

$$
\boldsymbol{e}_{n}^{T}(R L+\sigma I)=\mathbf{0}^{T}+\sigma \boldsymbol{e}_{n}^{T}=\sigma \boldsymbol{e}_{n}^{T} .
$$

To say that $\hat{A}=L^{-1} A L$ deflates in one step, means that the first eigenvalue $\sigma$ is obtained and the remaining eigenvalues of $A$ are the eigenvalues of the submatrix that results from discarding the last row and the last column of $\hat{A}$. This reduction of the order of the eigenproblem is called deflation.

In the general case, with LR algorithm there is still much work to be done in finding clever shift strategies that approach an eigenvalue in a stable way. The development of reliable shift strategies is an open domain.

### 3.4 Implicit double shifted LR algorithm

When $A=: A_{1}$ is a real unsymmetric matrix and has complex eigenvalues, to speed up the convergence to a complex eigenvalue it is necessary to choose a shift close to an eigenvalue and hence to choose a complex shift. If on iteration $i$ we take a complex shift $\sigma_{i}$, then the resulting matrix $A_{i+1}=R_{i} L_{i}+\sigma_{i} I$ will be complex. This means that all arithmetic has to be complex, increasing the cost by a factor of about 4 . Since complex eigenvalues of real matrices occur in complex conjugate pairs, we can shift by $\sigma_{i}$ and $\overline{\sigma_{i}}$ at the same time. It turns out that this will permit us to maintain real arithmetic. This double shift technique was developed by Francis for the QR algorithm: it uses real arithmetic and converges to real Schur canonical form. For the details on the implicit double shifted QR algorithm see Demmel [10, pp.170-172]. In this section we will describe a similar technique for the LR algorithm.

## Implicit double shifted LR algorithm

We will see in detail how to combine two consecutive iterations of LR algorithm choosing successive complex shifts $\sigma$ and $\bar{\sigma}$ such that the result after this double shift is again real.

Considering the first two steps, the result of shifting by $\sigma$ and $\bar{\sigma}$ in succession are

$$
\begin{align*}
A_{1}-\sigma I & =L_{1} R_{1} \\
A_{2} & =R_{1} L_{1}+\sigma I \text { so } A_{2}=L_{1}^{-1} A_{1} L_{1}  \tag{3.3}\\
A_{2}-\bar{\sigma} I & =L_{2} R_{2} \\
A_{3} & =R_{2} L_{2}+\bar{\sigma} I \text { so } A_{3}=L_{2}^{-1} A_{2} L_{2}=L_{2}^{-1} L_{1}^{-1} A_{1} L_{1} L_{2} \tag{3.4}
\end{align*}
$$

The real part of $\sigma$ will be denoted by $\Re \sigma$ and the imaginary part by $\Im \sigma$.

Lemma 3.4.1 Consider the application of two steps of LR algorithm with successive shifts $\sigma$ and $\bar{\sigma}$. Then

$$
\left(L_{1} L_{2}\right)\left(R_{2} R_{1}\right)=A_{1}^{2}-2(\Re \sigma) A_{1}+|\sigma|^{2} I
$$

Proof. Since $L_{2} R_{2}=A_{2}-\bar{\sigma} I$ and $L_{1} R_{1}=A_{1}-\sigma I$, we get

$$
\begin{aligned}
L_{1} L_{2} R_{2} R_{1} & =L_{1}\left(A_{2}-\bar{\sigma} I\right) R_{1} \\
& =L_{1}\left(L_{1}^{-1} A_{1} L_{1}-\bar{\sigma} I\right) R_{1} \\
& =\left(A_{1}-\bar{\sigma} I\right) L_{1} R_{1} \\
& =\left(A_{1}-\bar{\sigma} I\right)\left(A_{1}-\sigma I\right) \\
& =A_{1}^{2}-2(\Re \sigma) A_{1}+|\sigma|^{2} I .
\end{aligned}
$$

Define $M \equiv A_{1}^{2}-2(\Re \sigma) A_{1}+|\sigma|^{2} I$. Thus $\left(L_{1} L_{2}\right)\left(R_{2} R_{1}\right)$ is the LU decomposition of the real matrix $M$ and so ( $L_{1} L_{2}$ ) and ( $R_{2} R_{1}$ ) are both real. This means that

$$
A_{3}=\left(L_{1} L_{2}\right)^{-1} A_{1}\left(L_{1} L_{2}\right)
$$

is also real.
Moreover, the first column of $L_{1} L_{2}$ is proportional to the first column of $M$ :

$$
\begin{aligned}
M \boldsymbol{e}_{1} & =\left(L_{1} L_{2} R_{2} R_{1}\right) \boldsymbol{e}_{1} \\
& =r\left(L_{1} L_{2}\right) \boldsymbol{e}_{1},
\end{aligned}
$$

where $r$ is the element in position $(1,1)$ of upper triangular matrix $R_{2} R_{1}$ (remember that $\boldsymbol{e}_{1}$ is the first unit vector).

Our implementation of double shifted LR iteration will depend on the following theorem, known as the Implicit L Theorem. We will suppose that a reduction of the initial matrix $A$ to Hessenberg form has been performed. Thus, we will be considering double shifted LR on an Hessenberg matrix which is also a form that is preserved by this algorithm.

Theorem 3.4.1 Implicit L theorem Suppose that $L$ is unit lower triangular and $L^{-1} A L=H$ is unreduced upper Hessenberg. Then columns 2 through $n$ of $L$ are determined uniquely by the first column of $L$.

Proof. Let $\boldsymbol{l}_{j}$ and $\boldsymbol{h}_{j}, 1 \leq j \leq n$, be the columns of unit lower triangular matrix $L=\left(l_{i j}\right)_{1 \leq i, j \leq n}$ and upper unreduced Hessenberg matrix $H=\left(h_{i j}\right)_{1 \leq i, j \leq n}$, respectively.

Thus,

$$
\left.\begin{array}{rlrl}
L & =\left[\begin{array}{lllll}
\boldsymbol{l}_{1} & \boldsymbol{l}_{2} & \boldsymbol{l}_{3} & \cdots & \boldsymbol{l}_{n}
\end{array}\right] & & \text { and }
\end{array} \begin{array}{cccc}
1 & & & \\
l_{21} & 1 & & \\
& =\left[\begin{array}{llllll}
\boldsymbol{h}_{1} & \boldsymbol{h}_{2} & \boldsymbol{h}_{3} & \cdots & \boldsymbol{h}_{n}
\end{array}\right] \\
l_{31} & l_{32} & 1 & \\
\vdots & \vdots & \ddots & \ddots \\
l_{n 1} & l_{n 2} & \cdots & l_{n, n-1}
\end{array}\right] .\left[\begin{array}{lllll}
h_{11} & h_{12} & h_{13} & \cdots & h_{1 n} \\
h_{21} & h_{22} & h_{23} & \cdots & h_{2 n} \\
& h_{32} & h_{33} & \cdots & h_{3 n} \\
& & \ddots & \ddots & \vdots \\
& & & h_{n, n-1} & h_{n n}
\end{array}\right] .
$$

Equating entries on each side of the equation

$$
A L=L H
$$

shows, for the first column, that

$$
\begin{equation*}
A \boldsymbol{l}_{1}=L \boldsymbol{h}_{1}=\boldsymbol{l}_{1} h_{11}+\boldsymbol{l}_{2} h_{21} \tag{3.5}
\end{equation*}
$$

Then, looking at the first and second elements of $A l_{1}$, we get

$$
\boldsymbol{e}_{1}^{T} A \boldsymbol{l}_{1}=h_{11} \quad \text { and } \quad \boldsymbol{e}_{2}^{T} A \boldsymbol{l}_{1}=l_{21} h_{11}+h_{21}
$$

So, given a matrix $A$ and the first column $\boldsymbol{l}_{1}$ of $L$, elements $h_{11}$ and $h_{21}$ are uniquely determined. That is, the first column $\boldsymbol{h}_{1}$ of $H$ is uniquely determined by $\boldsymbol{l}_{1}$.

Now we can use (3.5) to obtain the second column of $L$,

$$
\boldsymbol{l}_{i 2}=\left(A \boldsymbol{l}_{1}-\boldsymbol{l}_{1} h_{11}\right)_{i 2} / h_{21}, \quad i=3, \ldots, n
$$

since, by assumption, $h_{21} \neq 0$.
To obtain columns $\boldsymbol{h}_{j}$ and $\boldsymbol{l}_{j+1}, 2 \leq j \leq n-1$, notice that

$$
\begin{equation*}
A \boldsymbol{l}_{j}=L \boldsymbol{h}_{j}=\sum_{i=1}^{j} \boldsymbol{l}_{i} h_{i j}+\boldsymbol{l}_{j+1} h_{j+1, j} \tag{3.6}
\end{equation*}
$$

Then, the first $j+1$ elements of $A l_{j}$ are

$$
\begin{align*}
e_{1}^{T} A l_{j} & =h_{1 j} \\
e_{2}^{T} A l_{j} & =h_{1 j} l_{21}+h_{2 j} \\
& \vdots  \tag{3.7}\\
e_{j}^{T} A l_{j} & =h_{1 j} l_{j 1}+h_{2 j} l_{j 2}+\cdots+h_{j j} \\
e_{j+1}^{T} A l_{j} & =h_{1 j} l_{j+1,1}+h_{2 j} l_{j+1,2}+\cdots+h_{j j} l_{j+1, j}+h_{j+1, j} .
\end{align*}
$$

Given columns $\boldsymbol{l}_{1}, \ldots, \boldsymbol{l}_{j}$, these equations determine, in turn, elements $h_{1 j}, \ldots, h_{j+1, j}$, that is, column $\boldsymbol{h}_{j}$ of $H$. And, since $h_{j+1, j} \neq 0$, from (3.6) we get column $\boldsymbol{l}_{j+1}$,

$$
\boldsymbol{l}_{i, j+1}=\frac{1}{h_{j+1, j}}\left(A \boldsymbol{l}_{j}-\sum_{i=1}^{j} \boldsymbol{l}_{i} h_{i j}\right)_{i, j+1}, \quad i=j+2, \ldots, n .
$$

Note that if $j=n-1$ we have $\boldsymbol{l}_{j+1}=\boldsymbol{e}_{n}$ and there is no need for any calculation.
Finally, for $j=n$, we obtain the last column of $H, \boldsymbol{h}_{n}=\left[h_{1 n} \ldots h_{n n}\right]^{T}$, the same way, but using only the first $j$ equations in (3.7).

The implicit L theorem implies that in the double shifted LR algorithm, to compute $A_{3}$ in (3.4) from $A_{1}$ we will only need to:

1. compute the first column of $L_{1} L_{2}$ which is proportional to the first column of $M$ and so can be gotten just by normalizing this column vector.
2. compute other columns of $L_{1} L_{2}$ implicitly using the implicit L theorem through elementary transformations.

Next we will describe a technique called bulge chasing that, after performing 1 , allow us to achieve goal 2 .

### 3.4.1 Bulge chasing

The technique of bulge chasing is justified by the implicit L theorem. This theorem plays the same role as the Implicit $Q$ Theorem [10, p.168] in the implementation of the implicit double shifted QR algorithm.

When $A=\left(a_{i j}\right)_{1 \leq i, j \leq n}$ is upper Hessenberg, then the first column of $M$ is

$$
\left[\begin{array}{c}
a_{11}^{2}+a_{12} a_{21}-2(\Re \sigma) a_{11}+|\sigma|^{2}  \tag{3.8}\\
a_{21}\left(a_{11}+a_{22}-2(\Re \sigma)\right) \\
a_{21} a_{32} \\
0 \\
\vdots \\
0
\end{array}\right] .
$$

We have

$$
\begin{equation*}
A_{3}=\mathbf{L}^{-1} A_{1} \mathbf{L} \tag{3.9}
\end{equation*}
$$

where $\mathbf{L} \equiv L_{1} L_{2}$ and the first column of $\mathbf{L}$ is proportional to (3.8). We will present a $6 \times 6$ example of bulge chasing to show how we obtain $A_{3}$ and $\mathbf{L}$ from $A_{1}$.

## Example 3.4.1

If we consider $n=6$, the transformation of $A_{1}$ to $A_{3}$ occurs accordingly to the following steps.

1. Let $\mathbf{I}_{1}=\left[\begin{array}{llllll}1 & * & * & 0 & 0 & 0\end{array}\right]^{T}$ be the first column of $\mathbf{L}$, obtained just by normalizing (3.8). Consider the elementary matrix $\mathbf{L}_{1}=I+\mathbf{l}_{1}^{\prime} e_{1}^{T}$, where $\mathbf{I}_{1}^{\prime}=\left[\begin{array}{llllll}0 & * & * & 0 & 0 & 0\end{array}\right]^{T}$, that is $\mathbf{I}_{1}$ with 0 in first position instead of 1 . So, since $\mathbf{L}_{1}^{-1}=I-\mathbf{I}_{1}^{\prime} \boldsymbol{e}_{1}^{T}$,

$$
\mathbf{L}_{1}^{-1} A_{1}=\left[\begin{array}{llllll}
x & x & x & x & x & x \\
x & x & x & x & x & x \\
+ & x & x & x & x & x \\
& x & x & x & x \\
& & & x & x & x \\
& & & & x & x
\end{array}\right] \text { and } A^{(1)} \equiv \mathbf{L}_{1}^{-1} A_{1} \mathbf{L}_{1}=\left[\begin{array}{llllll}
x & x & x & x & x & x \\
x & x & x & x & x & x \\
+ & x & x & x & x & x \\
+ & x & x & x & x \\
& & & x & x & x \\
& & & & x & x
\end{array}\right] .
$$

We see that there is a $2 \times 1$ bulge, indicated by plus signs.

Next we will apply a sequence of elementary similarity transformations such that each transformation pushes the bulge one row down and one column to the right. Finally the bulge is gotten rid of to restore upper Hessenberg form.
2. Form an elementary matriz (Gauss matriz) $\mathbf{L}_{2}^{-1}$ which affects only rows 3 and 4 of $\mathbf{L}_{2}^{-1} A^{(1)}$, zeroing out entries $(3,1)$ and $(4,1)$ of $A^{(1)}$. So, $\mathbf{L}_{2}=I+\mathbf{I}_{2}^{\prime} \boldsymbol{e}_{2}^{T}$ with $\mathbf{I}_{2}^{\prime}=\left[\begin{array}{llllll}0 & 0 & * & * & 0 & 0\end{array}\right]^{T}, \mathbf{L}_{2}^{-1}=I-\mathbf{I}_{2}^{\prime} \mathbf{e}_{2}^{T}$ and

$$
\mathbf{L}_{2}^{-1} A^{(1)}=\left[\begin{array}{llllll}
x & x & x & x & x & x \\
x & x & x & x & x & x \\
& x & x & x & x & x \\
& + & x & x & x & x \\
& & & x & x & x \\
& & & & x & x
\end{array}\right] \text { and } A^{(2)} \equiv \mathbf{L}_{2}^{-1} A^{(1)} \mathbf{L}_{2}=\left[\begin{array}{llllll}
x & x & x & x & x & x \\
x & x & x & x & x & x \\
& x & x & x & x & x \\
+ & x & x & x & x \\
+ & x & x & x \\
+ & & & x & x
\end{array}\right] .
$$

3. Form an elementary matriz $\mathbf{L}_{3}^{-1}$ which affects only rows 4 and 5 of $\mathbf{L}_{3}^{-1} A^{(2)}$, zeroing out entries $(4,2)$ and $(5,2)$ of $A^{(2)}$. We will have $\mathbf{L}_{3}=I+\mathbf{I}_{3}^{\prime} \boldsymbol{e}_{3}^{T}$ with $\mathbf{I}_{3}^{\prime}=\left[\begin{array}{llllll}0 & 0 & 0 & * & * & 0\end{array}\right]^{T}, \mathbf{L}_{3}^{-1}=I-\mathbf{I}_{3}^{\prime} \mathbf{e}_{3}^{T}$ and

$$
\mathbf{L}_{3}^{-1} A^{(2)}=\left[\begin{array}{llllll}
x & x & x & x & x & x \\
x & x & x & x & x & x \\
& x & x & x & x & x \\
& & x & x & x & x \\
& & + & x & x & x \\
& & & & x & x
\end{array}\right] \text { and } A^{(3)} \equiv \mathbf{L}_{3}^{-1} A^{(2)} \mathbf{L}_{3}=\left[\begin{array}{llllll}
x & x & x & x & x & x \\
x & x & x & x & x & x \\
& x & x & x & x & x \\
& x & x & x & x \\
& & & x & x & x \\
& & + & & x & x
\end{array}\right] .
$$

4. Form an elementary matriz $\mathbf{L}_{4}^{-1}$ which affects only rows 5 and 6 of $\mathbf{L}_{4}^{-1} A^{(3)}$, zeroing out entries $(5,3)$ and $(6,3)$ of $A^{(3)}$. We will have $\mathbf{L}_{4}=I+\mathbf{I}_{4}^{\prime} \boldsymbol{e}_{4}^{T}$ with

$$
\begin{gathered}
\mathbf{I}_{4}^{\prime}=\left[\begin{array}{llllll}
0 & 0 & 0 & 0 & * & *
\end{array}\right]^{T}, \mathbf{L}_{4}^{-1}=I-\mathbf{I}_{4}^{\prime} \mathbf{e}_{4}^{T} \text { and } \\
A^{(4)} \equiv \mathbf{L}_{4}^{-1} A^{(3)} \mathbf{L}_{4}=\left[\begin{array}{cccccc}
x & x & x & x & x & x \\
x & x & x & x & x & x \\
& x & x & x & x & x \\
& x & x & x & x \\
& & & x & x & x \\
& & & & x & x
\end{array}\right] .
\end{gathered}
$$

5. Finally to zero out entry $(6,4)$ of $A^{(4)}$ we form $\mathbf{L}_{5}=I+\mathbf{I}_{5}^{\prime} \mathbf{e}_{5}^{T}$ with $\mathbf{I}_{5}^{\prime}=\left[\begin{array}{llllll}0 & 0 & 0 & 0 & 0 & *\end{array}\right]^{T}$. We have $\mathbf{L}_{5}^{-1}=I-\mathbf{I}_{5}^{\prime} \boldsymbol{e}_{\mathbf{5}}{ }^{T}$,

$$
A^{(5)} \equiv \mathbf{L}_{5}^{-1} A^{(4)} \mathbf{L}_{5}=\left[\begin{array}{llllll}
x & x & x & x & x & x \\
x & x & x & x & x & x \\
& x & x & x & x & x \\
& & x & x & x & x \\
& & & x & x & x \\
& & & & x & x
\end{array}\right]
$$

and we are back to upper Hessenberg form.

In the end we have

$$
A^{(5)}=\left(\mathbf{L}_{1} \mathbf{L}_{2} \mathbf{L}_{3} \mathbf{L}_{4} \mathbf{L}_{5}\right)^{-1} A_{1} \mathbf{L}_{1} \mathbf{L}_{2} \mathbf{L}_{3} \mathbf{L}_{4} \mathbf{L}_{5} .
$$

In exact arithmetic the implicit $L$ theorem ensures that

$$
A_{3}=A^{(5)} \quad \text { and } \quad \mathbf{L}=L_{1} L_{2}=\mathbf{L}_{1} \mathbf{L}_{2} \mathbf{L}_{3} \mathbf{L}_{4} \mathbf{L}_{5} \diamond \diamond
$$

The rest of the this chapter is concerned with qd algorithms and its relation to LR algorithm.

### 3.5 The qd algorithms

Consider a tridiagonal matrix normalized into $J$ form as presented earlier in the previous chapter (page 30), that is, such that all superdiagonal entries are equal to 1 . When $L R$ is rewritten in the LU representation for a $J$ matrix one obtains the (progressive) qd algorithm. The letters q and d stand for quotient and difference. When we incorporate shifts we have the qds algorithm.

The differential qd algorithm (dqd) is a variant of qd that requires a little more of arithmetic effort than qd itself. Shifted version is named dqds. They are new to the scene of matrix computations and one feature that makes them attractive is that they seem to be more accurate than their rivals (see Parlett [40]). In particular, in the positive case, all eigenvalues can be found to high relative accuracy as long as the shifts preserve positivity.

### 3.5.1 Stationary qd algorithms

Triangular factors change in a complicated way under translation. Given $L$ and $U$ of the factorization $J=L U$ as presented in section 2.2.2 (page 31),

$$
L=\left[\begin{array}{cccccc}
1 & & & & & \\
l_{1} & 1 & & & & \\
& l_{2} & 1 & & & \\
& & \ddots & \ddots & & \\
& & & l_{n-2} & 1 & \\
& & & & l_{n-1} & 1
\end{array}\right], \quad U=\left[\begin{array}{cccccc}
u_{1} & 1 & & & & \\
& u_{2} & 1 & & & \\
& & u_{3} & 1 & & \\
& & & \ddots & \ddots & \\
& & & & u_{n-1} & 1 \\
& & & & & u_{n}
\end{array}\right]
$$

the task here is to compute $\bar{L}$ and $\bar{U}$ such that

$$
J-\sigma I=L U-\sigma I=\bar{L} \bar{U}
$$

for a given suitable shift $\sigma$. Equating entries on each side shows that

$$
\begin{aligned}
l_{i}+u_{i+1}-\sigma & =\bar{l}_{i}+\bar{u}_{i+1}, & & i=0, \ldots, n-1, \quad l_{0}=0, \quad \bar{l}_{0}=0, \\
l_{i} u_{i} & =\bar{l}_{i} \bar{u}_{i}, & & i=1, \ldots, n-1 .
\end{aligned}
$$

These relations yield the so-called stationary qd algorithm with shift:

$$
\begin{aligned}
\operatorname{stqds}(\sigma): & \bar{u}_{1}=u_{1}-\sigma \\
& \text { for } i=1, \ldots, n-1 \\
& \bar{l}_{i}=l_{i} u_{i} / \bar{u}_{i} \\
& \bar{u}_{i+1}=l_{i}+u_{i+1}-\sigma-\bar{l}_{i}
\end{aligned}
$$

end for.

Naturally it fails if $\bar{u}_{i}=0$ for some $i<n$.
An alternative algorithm for $\bar{L}$ and $\bar{U}$ involves more arithmetic effort and an auxiliary variable but has some advantages in accuracy for finite-precision arithmetic. To derive this algorithm we define a variable $t_{i}$ by

$$
t_{i+1} \equiv \bar{u}_{i+1}-u_{i+1}=l_{i}-\bar{l}_{i}-\sigma .
$$

Observe that

$$
\begin{aligned}
t_{i+1} & =l_{i}-l_{i} u_{i} / \bar{u}_{i}-\sigma \\
& =l_{i}\left(\bar{u}_{i}-u_{i}\right) / \bar{u}_{i}-\sigma \\
& =t_{i} l_{i} / \bar{u}_{i}-\sigma .
\end{aligned}
$$

The associated algorithm is called the differential form of the stationary qds algorithm. We will name it as dstqds.

$$
\begin{array}{ll}
\operatorname{dstqds}(\sigma): & t_{1}=-\sigma \\
& \text { for } i=1, \ldots, n-1 \\
& \bar{u}_{i}=u_{i}+t_{i} \\
& \bar{l}_{i}=u_{i}\left(l_{i} / \bar{u}_{i}\right) \\
& t_{i+1}=t_{i}\left(l_{i} / \bar{u}_{i}\right)-\sigma \\
& \text { end for } \\
& \bar{u}_{n}=u_{n}+t_{n} .
\end{array}
$$

In practice the $t$-values may be written over each other in a single variable $t$. If the common subexpression $l_{i} / \bar{u}_{i}$ is recognized then only one division is needed. Thus the dstqds exchanges a subtraction for a multiplication, so the extra cost is not excessive.

### 3.5.2 Progressive qd algorithms

Recall that the product $U L=J^{\prime}$ has the superdiagonal entries also equal to 1 . This section seeks the triangular factorization of $J^{\prime}-\sigma I$, not $J-\sigma I$ :

$$
J^{\prime}-\sigma I=U L-\sigma I=\hat{L} \hat{U}
$$

for a suitable shift $\sigma$. Equating entries on each side of the defining equation gives:

$$
\begin{aligned}
u_{i+1}+l_{i+1}-\sigma & =\hat{l}_{i}+\hat{u}_{i+1}, & & i=0, \ldots, n-1, \quad \hat{l}_{0}=0, \quad l_{n}=0, \\
l_{i} u_{i+1} & =\hat{l}_{i} \hat{u}_{i}, & & i=1, \ldots, n-1 .
\end{aligned}
$$

These relations give the so-called progressive qd algorithm with shift which we will call qds.

$$
\begin{array}{ll}
\operatorname{qds}(\sigma): & \hat{u}_{1}=u_{1}+l_{1}-\sigma \\
& \text { for } i=1, \ldots, n-1 \\
& \hat{l}_{i}=l_{i} u_{i+1} / \hat{u}_{i} \\
& \hat{u}_{i+1}=u_{i+1}+l_{i+1}-\sigma-\hat{l}_{i} \\
& \text { end for. }
\end{array}
$$

The algorithm qds fails when $\hat{u}_{i}=0$ for some $i<n$. When $\sigma=0$ we write simply qd, not qds.

There is an alternative implementation of qds that is slightly slower than qds but has compensating advantages. Lets define an auxiliary variable

$$
d_{i+1} \equiv \hat{u}_{i+1}-l_{i+1}=u_{i+1}-\hat{l}_{i}-\sigma .
$$

Observe that

$$
\begin{aligned}
d_{i+1} & =u_{i+1}-l_{i} u_{i+1} / \hat{u}_{i}-\sigma \\
& =u_{i+1}\left(\hat{u}_{i}-l_{i}\right) / \hat{u}_{i}-\sigma \\
& =d_{i} u_{i+1} / \hat{u}_{i}-\sigma .
\end{aligned}
$$

Rutishauser seems to have discovered the unshifted version some years after discovering qd but he did not make much use of it. He called it the differential qd algorithm (dqd) and the new shifted version will be called dqds.

$$
\begin{array}{cl}
\operatorname{dqds}(\sigma): & d_{1}=u_{1}-\sigma \\
& \text { for } i=1, \ldots, n-1 \\
\hat{u}_{i}=d_{i}+l_{i} \\
\hat{l}_{i}=l_{i}\left(u_{i+1} / \hat{u}_{i}\right) \\
& d_{i+1}=d_{i}\left(u_{i+1} / \hat{u}_{i}\right)-\sigma \\
\text { end for } \\
\hat{u}_{n}=d_{n}
\end{array}
$$

By definition, $\mathrm{dqd}=\mathrm{dqds}(0)$. In practice each $d_{i+1}$ may be written over its predecessor in a single variable $d$.

In the positive case $\left(l_{i}>0, i=1, \ldots, n-1 ; u_{i}>0, i=1, \ldots, n\right)$ dqd requires no subtractions and enjoys very high relative stability. In the symmetric case, dqds, even with the current simple shift strategies, achieves good accuracy in all eigenvalues and is faster then QR. See Fernando and Parlett [16]. In fact, this algorithm finds the singular values of a bidiagonal matrix in $\mathcal{O}\left(n^{2}\right)$ time, but as accurately and rather more efficiently than QR algorithm. It is the sequential algorithm of choice for singular values and is implemented in LAPACK ${ }^{1}$ [2].

In what concerns to an error analysis of dqds, it was proved in [16] that dqds yields mixed high relative stability. Given matrices $L$ and $D$ and shift $\sigma$, suppose that the dqds algorithm in finite precision produces representable output $\tilde{L}$ and $\tilde{D}$. We introduce ideal matrices $\hat{L}, \hat{D}, \breve{L}$ and $\breve{D}$ such that $\breve{L}$ and $\breve{D}$ is the output of dqds acting on $\hat{L}$ and $\hat{D}$ in exact arithmetic. Moreover $\hat{L}$ and $\hat{D}$ are small relative perturbations of $L$ and $D$, and $\breve{L}$

[^2]and $\breve{D}$ are small relative perturbations of $\tilde{L}$ and $\tilde{D}$. See figure bellow. This property is called mixed stability in [9] but note that the perturbations are relative ones. And dqds enjoys high mixed relative stability even with element growth.


Figure 3.1: Effects of roundoff for dqds

The diagram shows that, in the absence of division by zero, underflow or overflow, the diagram commutes and, for all $k, \hat{l}_{k}$ and $\hat{u}_{k}$ differ from $l_{k}$ and $u_{k}$ by 3 and 1 ulps ${ }^{2}$, respectively, and $\tilde{l}_{k}$ and $\tilde{u}_{k}$ differ from $\breve{l}_{k}$ and $\breve{u}_{k}$, respectively, by 2 ulps, at most.

### 3.6 Relation of LR algorithm for $J$ matrices to qds

The LR transform of $J$ is $J^{\prime}$ and the LR transform of $J^{\prime}$ is the matrix $J^{\prime \prime}$ defined in two steps by

$$
J^{\prime}=L^{\prime} U^{\prime}, \quad J^{\prime \prime}=U^{\prime} L^{\prime}
$$

In fact, the first step of the LR algorithm

$$
J=L U, \quad J^{\prime}=U L,
$$

defines $J^{\prime}$ and the second step defines $J^{\prime \prime}$,

$$
J^{\prime}=L^{\prime} U^{\prime}, \quad J^{\prime \prime}=U^{\prime} L^{\prime}
$$

[^3]Now qd applied to $L$ and $U$ yields $L^{\prime}$ and $U^{\prime}$ and so defines $J^{\prime \prime}$ implicitly. There is no need to form $J^{\prime}$ or $J^{\prime \prime}$.

When shifts are employed the situation is a little more complicated. It is necessary to look at two successive steps with shifts $\sigma_{1}$ and $\sigma_{2}$.

In shifted LR we have

$$
\begin{aligned}
J_{1}-\sigma_{1} I & =L_{1} U_{1} \\
J_{2} & =U_{1} L_{1}+\sigma_{1} I \\
J_{2}-\sigma_{2} I & =L_{2} U_{2} \\
J_{3} & =U_{2} L_{2}+\sigma_{2} I .
\end{aligned}
$$

In other words, the shifts are restored so that $J_{1}, J_{2}, J_{3}$ are similar. Note that

$$
\begin{aligned}
J_{2} & =U_{1} L_{1}+\sigma_{1} I \\
& =L_{1}^{-1} L_{1} U_{1} L_{1}+\sigma_{1} I \\
& =L_{1}^{-1}\left(J_{1}-\sigma_{1} I\right) L_{1}+\sigma_{1} I \\
& =L_{1}^{-1} J_{1} L_{1},
\end{aligned}
$$

and

$$
J_{3}=L_{2}^{-1} J_{2} L_{2}
$$

However, if $J_{2}$ is not to be formed one cannot explicitly add $\sigma_{1}$ back to the diagonal. On the other hand,

$$
J_{2}-\sigma_{2} I=U_{1} L_{1}-\left(\sigma_{2}-\sigma_{1}\right) I=L_{2} U_{2}
$$

In general, we have

$$
\begin{aligned}
J_{i+1}-\sigma_{i+1} I & =U_{i} L_{i}+\sigma_{i} I-\sigma_{i+1} I \\
& =U_{i} L_{i}-\left(\sigma_{i+1}-\sigma_{i}\right) I=L_{i+1} U_{i+1} .
\end{aligned}
$$

Thus, to find $L_{2}$ and $U_{2}$ from $L_{1}$ and $U_{1}$ it is only necessary to apply qds $\left(\sigma_{2}-\sigma_{1}\right)$. In other words, to get qds equivalent to LR with shifts $\left\{\sigma_{i}\right\}_{i=1}^{\infty}$ it is necessary to use the differences ( $\sigma_{i}-\sigma_{i-1}$ ) with qds.

In LR the shifts should converge to an eigenvalue of the original $J$ or $J^{\prime}$. In qds the shifts should converge to 0 and $u_{n} \rightarrow 0, l_{n-1} \rightarrow 0$ too and all shifts must be accumulated. See diagram in Figure 3.2.


Figure 3.2: Relation of LR to qds

In the positive case both LR and qds are accurate and efficient.
In practice, the LR algorithm for $J$ matrices avoids explicit calculation of the $L$ 's and $U$ 's and the transformation $J_{i} \rightarrow J_{i+1}$ is effected via a sequence of elementary similarity transformations. For comparison purposes, two implementations of LR, the explicit and the implicit shifted versions, can be found in [40, pp.469-471].

We end this section summarizing some advantages and disadvantages of the factored form $L U$.

## Advantages of the factored form

1. $L, U$ determines the entries of $J$ to greater than working-precision accuracy because the addition and multiplication of $l$ 's and $u$ 's is implicit. Thus, for instance, the ( $i, i$ ) entry of $J$ is given by $l_{i-1}+u_{i}$ implicitly but $f l\left(l_{i-1}+u_{i}\right)$ explicitly.
2. The mapping $L, U \rightarrow J$ is naturally parallel; for example, if $\boldsymbol{l}=\left(l_{i}\right)$ and $\boldsymbol{u}=\left(u_{i}\right)$, then $\boldsymbol{l} * \boldsymbol{u}$ gives the off-diagonal entries of $J$. In contrast, the mapping $J \rightarrow L, U$, that is, Gaussian elimination, is intrinsically sequential.
3. Singularity of $J$ is detectable by inspection when $L$ and $U$ are given, but only by calculation from $J$. So, $L U$ reveals singularity, $J$ does not.
4. First experiments with the measures of sensitivity presented in Chapter 7 show that $L U$ defines the eigenvalues better than $J$ does (usually).
5. Solution of $J x=b$ takes half the time when $L$ and $U$ are available.

## Disadvantages of the factored form

The mapping $J \mapsto L, U$ is not everywhere defined. Even when the factorization exists it can happen that $\|L\|$ and $\|U\|$ greatly exceed $\|J\|$. This is very bad for applying the LR algorithm but harmless when eigenvectors are to be calculated. So, some care is needed to consider the goal before stigmatizing a process as unstable. Moreover, in the eigenvalue context we are free to replace $J$ by $J-\sigma I=L U$ for some suitable chosen shift $\sigma$ that gives acceptable $L$ and $U$. And there is a lot that can be done in improving existing shift strategies so that element growth can be monitored.

## Chapter 4

## Convergence results for LR

In this chapter we establish a new convergence result for the basic $L R$ algorithm on a real unreduced tridiagonal matrix with a one-point spectrum - the Jordan form is only one big Jordan block. First we show the classical convergence results and summarize eigenvector properties of real unreduced unsymmetric tridiagonal matrices.

### 4.1 Classical results for the convergence of LR algorithm

The $k^{\text {th }}$ iteration of the basic LR method is based on the LU decomposition $A_{k}=L_{k} U_{k}$ and on the multiplication of the factors $L_{k}$ and $U_{k}$ in reverse order to get the matrix $A_{k+1}$. The sequence of matrices $A=: A_{1}, A_{2}, A_{3}, \ldots$ is then generated by

## LR algorithm

$$
\begin{aligned}
& A_{1}=A \\
& \text { for } i=1,2, \ldots \\
& \quad \text { Factor } A_{i}=L_{i} R_{i} \quad \text { (LU factorization) } \\
& \quad A_{i+1}=R_{i} L_{i} \\
& \text { end }
\end{aligned}
$$

It is easy to see that $A_{k+1}$ is similar to $A_{k}$,

$$
\begin{equation*}
A_{k+1}=R_{k} L_{k}=L_{k}^{-1}\left(L_{k} R_{k}\right) L_{k}, \tag{4.1}
\end{equation*}
$$

and then, by an inductive argument, $A_{k+1}$ is similar to $A_{1}$.
Lemma 4.1.1 Let $A_{1}:=A$ and $\left\{A_{i}\right\}_{i=1}^{\infty}$ be the sequence of matrices generated by LR algorithm, $i=1,2, \ldots$. Then

$$
A_{i+1}=\left(L_{1} L_{2} \ldots L_{i}\right)^{-1} A_{1}\left(L_{1} L_{2} \ldots L_{i}\right), \quad i=1,2, \ldots \ldots
$$

Proof. Use (4.1) repeatedly.

Lemma 4.1.2 Let matrices $\mathcal{L}_{i}$ and $\mathcal{U}_{i}$ be defined by

$$
\mathcal{L}_{i} \equiv L_{1} L_{2} \ldots L_{i} \quad \text { and } \quad \mathcal{U}_{i} \equiv R_{i} R_{i-1} \ldots R_{1} .
$$

Then $\mathcal{L}_{i} \mathcal{U}_{i}$ is the LU decomposition of $A_{1}^{i}$.

Proof. From lemma 4.1.1 we have

$$
\begin{equation*}
L_{1} L_{2} \ldots L_{i-1} A_{i}=A_{1} L_{1} L_{2} \ldots L_{i-1} \tag{4.2}
\end{equation*}
$$

Notice that matrices $\mathcal{L}_{i}$ and $\mathcal{U}_{i}$ are unit-lower triangular and upper-triangular, respectively. Consider the product $\mathcal{L}_{i} \mathcal{U}_{i}$. For $i=2,3, \ldots$, we have

$$
\begin{aligned}
\mathcal{L}_{i} \mathcal{U}_{i} & =L_{1} L_{2} \ldots L_{i-1}\left(L_{i} R_{i}\right) R_{i-1} \ldots R_{1} \\
& =L_{1} L_{2} \ldots L_{i-1} A_{i} R_{i-1} \ldots R_{1} \\
& =A_{1} L_{1} L_{2} \ldots L_{i-1} R_{i-1} \ldots R_{1} \quad(\text { from (4.2)) } \\
& =A_{1} \mathcal{L}_{i-1} \mathcal{U}_{i-1}
\end{aligned}
$$

Repeated application of this result shows that

$$
\begin{equation*}
\mathcal{L}_{i} \mathcal{U}_{i}=A_{1}^{i}, \tag{4.3}
\end{equation*}
$$

so that $\mathcal{L}_{i} \mathcal{U}_{i}$ gives the LU decomposition of $A_{1}^{i}$.

Thus, we have

$$
A_{i+1}=\mathcal{L}_{i}^{-1} A_{1} \mathcal{L}_{i} \quad \text { and } \quad \mathcal{L}_{i} \mathcal{U}_{i}=A_{1}^{i},
$$

that is, $i$ steps of LR applied to $A_{1}$ are equivalent to a similarity given by a factorization of $A_{1}^{i}$.

### 4.1.1 Convergence of LR algorithm in the simplest case

We now restrict ourselves to the case when $A_{1}$ is nonsingular and has eigenvalues $\lambda_{i}, i=1, \ldots, n$, of distinct modulus, so that it necessarily has linear elementary divisors. We may write

$$
\begin{equation*}
A_{1}=X \operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right) X^{-1}=: X D Y, \tag{4.4}
\end{equation*}
$$

where the columns of $X$ are the right eigenvectors of $A_{1}$ and the rows of $Y$ are the row eigenvectors of $A_{1}$.

Suppose we order the eigenvalues $\lambda_{i}, i=1, \ldots, n$, of $A_{1}$ so that they satisfy

$$
\begin{equation*}
\left|\lambda_{1}\right|>\left|\lambda_{2}\right|>\cdots>\left|\lambda_{n}\right| . \tag{4.5}
\end{equation*}
$$

Then, under certain restrictions, we have

$$
L_{i} \rightarrow I \quad \text { and } \quad R_{i} \rightarrow A_{i} \rightarrow\left[\begin{array}{cccc}
\lambda_{1} & * & * & * \\
& \lambda_{2} & * & * \\
& & \ddots & \vdots \\
& & & \lambda_{n}
\end{array}\right] \quad \text { as } \quad i \rightarrow \infty
$$

Theorem 4.1.1 If $A=X \operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right) X^{-1}$ and

1. $\left|\lambda_{1}\right|>\left|\lambda_{2}\right|>\cdots>\left|\lambda_{n}\right|$,
2. the leading principal minors of $X$ and $X^{-1}$ are nonzero,
then

$$
\lim A_{i}=A_{\infty}
$$

exists and is upper triangular with the $i^{\text {th }}$ diagonal entry of $A_{\infty}$ equal to $\lambda_{i}$.

Before writing the proof of this result, observe two simple properties of triangular decomposition. The triangular decomposition of a nonsingular matrix $X$, if it exists, is unique. Further, if

$$
X=I+F
$$

then, if $\|F\|$ is sufficiently small, the triangular decomposition of $X$ exists, and if

$$
X=I+F=L U
$$

then

$$
L, U \rightarrow I \quad \text { as } \quad\|F\| \rightarrow 0 .
$$

This can be verified if we think in terms of Gaussian elimination.
The proof of theorem 4.1.1 is rather technical (see Wilkinson [59] and [60, pp.487-492]) and it was first given by Rutishauser [49]. It is based on the fact that if matrix $X$ of right eigenvectors of $A$ admits triangular factorization

$$
X=L_{X} U_{X}
$$

then

$$
\begin{aligned}
\mathcal{L}_{i} \rightarrow L_{X} \quad \text { and } \quad A_{i}=\mathcal{L}_{i-1}^{-1} A_{1} \mathcal{L}_{i-1} \rightarrow L_{X}^{-1} A_{1} L_{X} & =U_{X} X^{-1} A X U_{X}^{-1} \\
& =U_{X} \operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right) U_{X}^{-1},
\end{aligned}
$$

showing that the limiting $A_{i}$ is upper-triangular with diagonal elements $\lambda_{1}, \ldots, \lambda_{n}$.
Proof. From (4.3) we have

$$
\begin{equation*}
\mathcal{L}_{i} \mathcal{U}_{i}=A_{1}^{i}=X D^{i} Y \tag{4.6}
\end{equation*}
$$

and hence $\mathcal{L}_{i} \mathcal{U}_{i}$ is the LU decomposition of $X D^{i} Y$.

Since we are assuming that the leading principal minors of $X$ and $Y$ are non-zero, then $X$ and $Y$ have LU decompositions. We write

$$
X=L_{X} U_{X} \quad \text { and } \quad Y=L_{Y} U_{Y},
$$

and attempt to construct the LU decomposition of $X D^{i} Y$. We have

$$
\begin{equation*}
X D^{i} Y=X D^{i} Y=X D^{i} L_{Y} U_{Y}=X\left(D^{i} L_{Y} D^{-i}\right)\left(D^{i} U_{Y}\right) \tag{4.7}
\end{equation*}
$$

If we write

$$
\begin{equation*}
D^{i} L_{Y} D^{-i}=I+F_{i}, \tag{4.8}
\end{equation*}
$$

with $F_{i}=\left(f_{k j}^{(i)}\right)$ and $L_{Y}=\left(l_{k j}\right)$, we have

$$
f_{k j}^{(i)}= \begin{cases}l_{k j}\left(\frac{\lambda_{k}}{\lambda_{j}}\right)^{i}, & k>j  \tag{4.9}\\ 0, & k \leq j\end{cases}
$$

From relations (4.5) and (4.9) we see that $F_{i} \rightarrow O$ as $i \rightarrow \infty$. Hence

$$
\begin{aligned}
X D^{i} Y=X\left(I+F_{i}\right) D^{i} U_{Y} & =L_{X} U_{X}\left(I+F_{i}\right) D^{i} U_{Y} \\
& =L_{X}\left(I+U_{X} F_{i} U_{X}^{-1}\right) U_{X} D^{i} U_{Y} \\
& =L_{X}\left(I+G_{i}\right) U_{X} D^{i} U_{Y}
\end{aligned}
$$

where $G_{i} \rightarrow O$ as $i \rightarrow \infty$. Thus, for sufficiently large $i, I+G_{i}$ has an LU decomposition which we may write in the form

$$
I+G_{i}=\left(I+L_{G}^{(i)}\right)\left(I+U_{G}^{(i)}\right)
$$

where $L_{G}^{(i)}, U_{G}^{(i)} \rightarrow O$ as $i \rightarrow \infty$.
Finally from (4.6) we have

$$
\mathcal{L}_{i} \mathcal{U}_{i}=L_{X}\left(I+L_{G}^{(i)}\right)\left(I+U_{G}^{(i)}\right) U_{X} D^{i} U_{Y}
$$

which, from the uniqueness of the LU decomposition, gives

$$
\mathcal{L}_{i}=L_{X}\left(I+L_{G}^{(i)}\right) \rightarrow L_{X} \quad \text { as } \quad i \rightarrow \infty .
$$

Hence, since $A_{i+1}=\mathcal{L}_{i}{ }^{-1} A_{1} \mathcal{L}_{i}$, we have

$$
\begin{equation*}
A_{i+1}=\mathcal{L}_{i}^{-1} X D X^{-1} \mathcal{L}_{i} \rightarrow L_{X}^{-1} L_{X} U_{X} D U_{X}^{-1} L_{X}^{-1} L_{X}=U_{X} D U_{X}^{-1} \tag{4.10}
\end{equation*}
$$

The matrix $U_{X} D U_{X}^{-1}$ is upper-triangular and has diagonal elements $\lambda_{1}, \ldots, \lambda_{n}$ in this order.

According to [59, p.80], we should note that it was assumed that $A_{i}$ has a triangular decomposition at all stages. This is not assured by the non-vanishing of the principal minors of $X$ and $Y$ and hence even when these conditions are satisfied, the LR method can break down. Such a failure corresponds to the non-existence of a decomposition of $I+G_{i}$. Since $I+G_{i} \rightarrow O$ this cannot happen at a late stage in the process.

If we remove the condition that $Y$ has non-vanishing principal minors, a phenomenon called disorder of latent roots occurs - $A_{i+1}$ tends to an upper triangular matrix having as its diagonal the eigenvalues but no longer in monotonic decreasing order. The eigenvalues are therefore disordered in the limiting matrix and this phenomenon happens to be unstable in practice.

The speed of convergence of the LR algorithm is determined essentially by the speed at which the elements of $F_{i}$ tend to zero and hence on the quantities $\frac{\lambda_{k}}{\lambda_{j}}, k>j$. If $\left|\frac{\lambda_{k}}{\lambda_{j}}\right|$ is close to 1 , that is, if the separation of the eigenvalues is poor, convergence may be slow. So, the speed of convergence of $A_{i}$ to upper-triangular form depends, particularly, on the ratios $\left|\frac{\lambda_{k+1}}{\lambda_{k}}\right|$ which are the highest.

If we let $L_{i}=\left(l_{k j}^{(i)}\right)$, from the relation $L_{i}=\mathcal{L}_{i-1}^{-1} \mathcal{L}_{i}$ it can be proved that

$$
l_{k j}^{(i)}=\mathcal{O}\left(\frac{\lambda_{k}}{\lambda_{j}}\right)^{i} \quad \text { as } \quad i \rightarrow \infty \quad(k>j)
$$

So if we have $\lambda_{n} \approx 0$ then

$$
l_{n, j}^{(i)}=\mathcal{O}\left(\frac{\lambda_{n}}{\lambda_{j}}\right)^{i}, \quad j<n
$$

converges quickly to 0 , that is, the last line of $L_{i}$ converges quickly to $\boldsymbol{e}_{n}^{T}$. Thus $A_{i+1}=R_{i} L_{i}$ will have the last line converging quickly to $\lambda_{n} \boldsymbol{e}_{n}^{T} \approx \mathbf{0}$. This explains why it is more
efficient to apply LR algorithm to $A-\sigma I$ with $\sigma$ close $\lambda_{n}$ and how shifts of origin accelerate convergence.

Note that in establishing this result of convergence we made the assumption that all the eigenvalues are of different magnitude and in this case $A$ can not be a real matrix with complex conjugate eigenvalues.

### 4.1.2 Eigenvalues of coincident absolute value

If $\left|\frac{\lambda_{k}}{\lambda_{j}}\right|=1$ it appears at first sight that we no longer have convergence. But this is not the case. If the $\left|\lambda_{i}\right|$ are not distinct then, loosely speaking, $A_{i}$ may be said to converge to block triangular form. See Parlett [41]. More precisely, suppose

$$
\left|\lambda_{1}\right|=\left|\lambda_{2}\right|=\ldots=\left|\lambda_{p}\right|>\left|\lambda_{p+1}\right| .
$$

Strictly speaking, $A_{\infty}$ may not exist but, as $i \rightarrow \infty, A_{\infty}$ becomes reduced. The elements in the first $p$ columns may not converge but the characteristic polynomial of the leading principal submatrix of order $p$ does converge to the monic polynomial with roots $\lambda_{1}, \ldots, \lambda_{p}$. The complementary submatrix has eigenvalues which converge to $\lambda_{p+1}, \ldots, \lambda_{n}$. If any of these eigenvalues have equal modulus then this submatrix will become reduced also. Thus, as $i \rightarrow \infty$, the submatrix blocks which become isolated along the diagonal correspond to groups of eigenvalues of equal modulus.

In the important case of real matrices with complex conjugate pairs, $A_{i}$ may be expected to have along the diagonal, for large $i$, isolated $2 \times 2$ real submatrix which yield the eigenvalues very conveniently.

Theorem 4.1.2 Let $A=: A_{1}$ be a diagonalizable real square matrix. Suppose that the LU factorization of $A_{i}$ exists at every step $i=1,2, \ldots$ of LR algorithm. Then $A_{i}$ tends to a block upper triangular form as $i \rightarrow \infty$. The diagonal blocks, say $X_{1}^{i}, X_{2}^{i}, \ldots, X_{p}^{i}$, need not converge but the eigenvalues of $X_{j}^{i}, j=1, \ldots, p$, converge to the set of eigenvalues of $j^{\text {th }}$ largest magnitude. In particular, if the eigenvalues have distinct magnitude, except for
complex conjugate pairs, then the blocks $X_{j}^{i}$ are either $1 \times 1$ for real eigenvalues, or $2 \times 2$, for a complex conjugate pair.

In [59] Wilkinson separates the proof of this result into two parts: equal eigenvalues having linear elementary divisors and unequal eigenvalues of equal modulus. That is, it is still assumed that the matrix $A$ is diagonalizable.

For the case of non-linear divisors it is given in the last section of [59] a simple counterexample that shows immediately that the LR algorithm does not necessarily give convergence to an upper-triangular matrix. The matrix

$$
A=\left[\begin{array}{ccc}
a & 0 & 0 \\
1 & a & 0 \\
0 & 1 & a
\end{array}\right], a \in \mathbb{R},
$$

is LR invariant, that is, $A_{i}=A$ for all $i$.

### 4.2 Eigenvector properties of an unreduced tridiagonal

Eigenvector matrices of real unreduced symmetric tridiagonal matrices have several attractive properties and have been studied widely in the literature. See Parlett [38, Chapter 7]. The eigenvalues are real and distinct and key properties are

- the first and last entries cannot vanish; there are very elegant formulae for the squares of entries of normalized eigenvectors.
- When the off-diagonal entries are all of the same sign, the eigenvector for the rightmost (largest) eigenvalue has no sign changes, for the second largest eigenvalue, the eigenvector has one sign change, and so on. The eigenvector for the leftmost (smallest) eigenvalue has the maximal number of sign changes, namely $n-1$ for a $n \times n$ matrix. See Gantmakher and Krein [19] and Fiedler [17].

Our interest is in the real unsymmetric case and we expect the matrix spectrum to have a mixture of real and complex eigenvalues. Of the properties above only the first extends to
the case of our interest. The proof is identical to the symmetric case and will be omitted. A new difficulty in our case is that the eigenvalues need not be simple, so the Jordan form may not be diagonal. In such cases the eigenvector matrix must be filled out with the so-called generalized eigenvectors with the property that

$$
(C-\lambda I)^{j} v=0, \quad(C-\lambda I)^{j-1} v \neq 0 .
$$

We say $v$ is an eigenvector of grade $j$ (see page 10).
In what follows we shall present some properties of eigenvector matrices that are sufficient to guarantee convergence of the basic LR algorithm without invoking the extra hypotheses needed by Rutishauser and Wilkinson for the general case. To the best of our knowledge these results are new.

Our convergence theory for LR algorithm requires that certain matrices $X$ permit triangular factorization $X=L D U$ or $X=L(D U)$. This property plays a prominent role in Linear Systems Theory.

We will say that $X$ is strongly (or completely) regular with the meaning that $X$ and all its leading principal submatrices are invertible. We shall use the terms "strongly regular" and "permits LU" interchangeably. To be precise, we note that a singular matrix may permit triangular factorization but in our work all the matrices of interest will be invertible.

We proceed from the easier cases to the more difficult in stages.
Most of our results extend directly to complex unreduced tridiagonal matrices but we focus on real matrices for simplicity and because it is the most frequent case in applications.

Consider an unreduced real tridiagonal matrix $C=\operatorname{tridiag}(\boldsymbol{b}, \boldsymbol{a}, \boldsymbol{c})$,

$$
C=\left[\begin{array}{ccccc}
a_{1} & c_{1} & & & \\
b_{1} & a_{2} & c_{2} & & \\
& \ddots & \ddots & \ddots & \\
& & b_{n-2} & a_{n-1} & c_{n-1} \\
& & & b_{n-1} & a_{n}
\end{array}\right] \in \mathbb{R}^{n \times n}
$$

with $b_{i} c_{i} \neq 0, i=1, \ldots, n-1$.

Define monic polynomials $p_{0}, p_{1}, \ldots, p_{n}$ by

$$
p_{0}(\tau)=1, \quad p_{j}(\tau):=\operatorname{det}\left(\tau I_{j}-C_{j}\right), \quad j=1, \ldots, n,
$$

where $I_{j}$ represents the $j \times j$ identity matrix and $C_{j}$ the $j^{\text {th }}$ leading principal submatrix of C.

### 4.2.1 All eigenvalues distinct

Suppose all eigenvalues of $C$ are distinct and let the spectrum be

$$
\Lambda=\left\{\lambda_{1}, \lambda_{2}, \ldots, \lambda_{n}\right\}
$$

The following matrix plays a key role in our results,

$$
P=P_{C}=\left[p_{i-1}\left(\lambda_{j}\right)\right]_{i, j=1}^{n} .
$$

The notation means that the $(i, j)$ element of $P$ is equal to $p_{i-1}\left(\lambda_{j}\right)$. So the $j^{\text {th }}$ column of $P$ is given by the column vector

$$
\boldsymbol{p}\left(\lambda_{j}\right):=\left[\begin{array}{llll}
p_{0}\left(\lambda_{j}\right) & p_{1}\left(\lambda_{j}\right) & \ldots & p_{n-1}\left(\lambda_{j}\right)
\end{array}\right]^{T}, \quad j=1, \ldots, n,
$$

that is,

$$
P=\left[\begin{array}{cccc}
1 & 1 & \ldots & 1  \tag{4.11}\\
p_{1}\left(\lambda_{1}\right) & p_{1}\left(\lambda_{2}\right) & \ldots & p_{1}\left(\lambda_{n}\right) \\
\vdots & \vdots & & \vdots \\
p_{n-1}\left(\lambda_{1}\right) & p_{n-1}\left(\lambda_{2}\right) & \ldots & p_{n-1}\left(\lambda_{n}\right)
\end{array}\right] .
$$

$P$ and $P^{T}$ are called polynomial Vandermond matrices. The standard Vandermond matrix $V$ is defined by

$$
V=V_{\Lambda}=\left[\lambda_{i}^{j-1}\right]_{i, j=1}^{n},
$$

that is,

$$
V=\left[\begin{array}{ccccc}
1 & \lambda_{1} & \lambda_{1}^{2} & \ldots & \lambda_{1}^{n-1}  \tag{4.12}\\
1 & \lambda_{2} & \lambda_{2}^{2} & \ldots & \lambda_{2}^{n-1} \\
\vdots & \vdots & & \vdots & \\
1 & \lambda_{n} & \lambda_{n}^{2} & \ldots & \lambda_{n}^{n-1}
\end{array}\right] .
$$

The valuable property of $V$ is that

$$
\begin{equation*}
\operatorname{det}(V)=\prod_{i>j}\left(\lambda_{i}-\lambda_{j}\right) \tag{4.13}
\end{equation*}
$$

where the product extends over all pairs $(i, j)$ with $n \geq i>j \geq 1 ; n(n-1) / 2$ terms in all.
When the $\lambda_{i}$ are distinct then $V$ is strongly regular because each leading principal submatrix of $V$ is also a Vandermond built from a subset of the eigenvalues.

Lemma 4.2.1 If all the eigenvalues $\lambda_{i}, i=1, \ldots, n$, are distinct then $V$ is strongly regular.
Moreover $\Pi V$ is strongly regular for any permutation matrix $\Pi$ because no particular ordering of the eigenvalues was specified in the definition.

Lemma 4.2.2 $\operatorname{det}(P)=\operatorname{det}\left(V^{T}\right)=\operatorname{det}(V)$.
Proof. $P$ may be reduced to $V^{T}$ by elementary row operations that leave the determinant unchanged. For example, if $p_{1}(\tau)=\tau-k$ then add $k$ times row 1 to row 2 of $P$ and the second row becomes $\left[\begin{array}{lll}\lambda_{1} & \ldots & \lambda_{n}\end{array}\right]=\boldsymbol{e}_{2}^{T} V^{T}$. And so on.

With the aid of $P=\left[\begin{array}{llll}\boldsymbol{p}\left(\lambda_{1}\right) & \boldsymbol{p}\left(\lambda_{2}\right) & \cdots & \boldsymbol{p}\left(\lambda_{n}\right)\end{array}\right]$ we can find simple forms for the column and row eigenvectors of $C$.

We will denote

$$
\begin{aligned}
& D_{b}:=\operatorname{diag}\left(1, b_{1}, b_{1} b_{2}, b_{1} b_{2} b_{3}, \ldots, b_{1} b_{2} \cdots b_{n-1}\right), \\
& D_{c}:=\operatorname{diag}\left(1, c_{1}, c_{1} c_{2}, c_{1} c_{2} c_{3}, \ldots, c_{1} c_{2} \cdots c_{n-1}\right) .
\end{aligned}
$$

Lemma 4.2.3 With the notation given above

$$
C\left(D_{c}^{-1} P\right)=\left(D_{c}^{-1} P\right) \Lambda, \quad\left(P^{T} D_{b}^{-1}\right) C=\Lambda\left(P^{T} D_{b}^{-1}\right)
$$

Proof. The result is a reformulation of the celebrated three term recurrence (3TR) associated with $C$ :

$$
\begin{align*}
p_{j+1}(\tau) & =\left(\tau-a_{j+1}\right) p_{j}(\tau)-b_{j} c_{j} p_{j-1}(\tau), & & j=1,2, \ldots, n-1  \tag{4.14}\\
p_{1}(\tau) & =\left(\tau-a_{1}\right)=\left(\tau-a_{1}\right) p_{0}(\tau), \text { since } & & p_{0}(\tau)=1 .
\end{align*}
$$

If we define $p_{-1}(\tau)=0$ then we may use the 3 TR for $j=0$.
Rewrite the 3 TR in the form

$$
b_{j} c_{j} p_{j-1}(\tau)+\left(a_{j+1}-\tau\right) p_{j}(\tau)+p_{j+1}(\tau)=0, \quad j=0, \ldots, n-1 .
$$

The key step is to divide through by $c_{1} \cdots c_{j}(\neq 0)$ and rearrange coefficients,

$$
b_{j} \frac{p_{j-1}(\tau)}{c_{0} c_{1} \cdots c_{j-1}}+\left(a_{j+1}-\tau\right) \frac{p_{j}(\tau)}{c_{0} c_{1} \cdots c_{j}}+c_{j+1} \frac{p_{j+1}(\tau)}{c_{1} \cdots c_{j+1}}=0, \quad j=0,1, \ldots, n-1,
$$

where $b_{0}=0, b_{n}=c_{n}=1$ and $c_{0}=1$. This set of equations may be written as

$$
(C-\tau I) D_{c}^{-1} \boldsymbol{p}(\tau)=-\boldsymbol{e}_{n} \frac{p_{n}(\tau)}{c_{1} \cdots c_{n-1} c_{n}}
$$

Thus, since for each $\tau=\lambda_{k}$ we have

$$
p_{n}\left(\lambda_{k}\right)=0,
$$

it follows that

$$
\begin{equation*}
\left(C-\lambda_{k} I\right) D_{c}^{-1} \boldsymbol{p}\left(\lambda_{k}\right)=-\boldsymbol{e}_{n} \frac{p_{n}\left(\lambda_{k}\right)}{c_{1} \cdots c_{n-1} c_{n}}=\mathbf{0} \quad, k=1, \ldots, n . \tag{4.15}
\end{equation*}
$$

Thus, in matrix terms,

$$
C D_{c}^{-1} P=D_{c}^{-1} P \Lambda .
$$

Note that $C^{T}=D_{c} D_{b}^{-1} C D_{c}^{-1} D_{b}$ and then

$$
C^{T} D_{b}^{-1} P=D_{c} D_{b}^{-1}\left(C D_{c}^{-1} P\right)=D_{c} D_{b}^{-1}\left(D_{c}^{-1} P \Lambda\right)=D_{b}^{-1} P \Lambda .
$$

These row and column eigenvectors are not scaled properly to be inverses of each other. Since the row eigenvectors for $\lambda_{j}$ annihilates all the column eigenvectors for different eigenvalues we may define a special diagonal matrix $\Delta=\Delta_{C}$ by

$$
\begin{aligned}
\left(P^{T} D_{b}^{-1}\right)\left(D_{c}^{-1} P\right) & =\Delta:=\operatorname{diag}\left(\delta_{1}, \ldots, \delta_{n}\right) \\
& \delta_{j}=\boldsymbol{p}\left(\lambda_{j}\right)^{T}\left(D_{b} D_{c}\right)^{-1} \boldsymbol{p}\left(\lambda_{j}\right) \neq 0
\end{aligned}
$$

Note that, if $\delta_{j}=0$ then the row eigenvector for $\lambda_{j}$ annihilates all the column eigenvectors and would be $\mathbf{0}^{T}$ which contradicts the definition of an eigenvector.

The matrix $\Delta^{-1}$ may be attached to either $D_{c}^{-1} P$ or $D_{b}^{-1} P$ or shared between them. In general, $\Delta$ will be indefinite. We will have

$$
\left(P^{T} D_{b}^{-1}\right)\left(D_{c}^{-1} P \Delta^{-1}\right)=\left(D_{c}^{-1} P \Delta^{-1}\right)\left(P^{T} D_{b}^{-1}\right)=I
$$

and

$$
\left(\Delta^{-1} P^{T} D_{b}^{-1}\right)\left(D_{c}^{-1} P\right)=\left(D_{c}^{-1} P\right)\left(\Delta^{-1} P^{T} D_{b}^{-1}\right)=I
$$

Theorem 4.2.1 The spectral decomposition of $C$ with simple eigenvalues may be written

$$
C=\left(D_{c}^{-1} P\right) \Lambda\left(\Delta^{-1} P^{T} D_{b}^{-1}\right)=\left(D_{c}^{-1} P \Delta^{-1}\right) \Lambda\left(P^{T} D_{b}^{-1}\right)
$$

Theorem 4.2.2 When $C$ has simple eigenvalues then both column and row eigenvector matrices $D_{c}^{-1} P \Delta^{-1}$ and $P^{T} D_{b}^{-1}$, respectively, are strongly regular.

The proof needs only the following easy results.

Lemma 4.2.4 For any invertible diagonal matrices $D^{\prime}$ and $D^{\prime \prime}, M$ is strongly regular if, and only if, $D^{\prime} M D^{\prime \prime}$ is strongly regular.

Proof. Let $\left(D^{\prime} M D^{\prime \prime}\right)_{j}$ be the $j^{\text {th }}$ leading principal submatrix of $D^{\prime} M D^{\prime \prime}$. We have

$$
\operatorname{det}\left(\left(D^{\prime} M D^{\prime \prime}\right)_{j}\right)=\operatorname{det}\left(D_{j}^{\prime}\right) \operatorname{det}\left(M_{j}\right) \operatorname{det}\left(D_{j}^{\prime \prime}\right) \neq 0, \quad j=1, \ldots, n
$$

since all the determinants are nonzero.

Lemma 4.2.5 When all the eigenvalues $\lambda_{i}, i=1, \ldots, n$, are distinct, both $P$ and $P^{T}$ are strongly regular.

Proof. Use lemmas 4.2.1 and 4.2.2
Next we consider the opposite extreme, a maximal Jordan block of $C$.

### 4.2.2 The one-point spectrum

Suppose now that $C$ 's spectrum consists of a single nonzero point $\lambda$ and such that its Jordan form is

$$
J=\lambda I+N
$$

where $N$ is the nilpotent matrix

$$
N=\left[\begin{array}{ccccc}
0 & 1 & & & \\
& 0 & 1 & & \\
& & \ddots & \ddots & \\
& & & 0 & 1 \\
& & & & 0
\end{array}\right]
$$

From the previous section we know that the relation (4.15) is verified only for $\lambda$,

$$
(C-\lambda I) D_{c}^{-1} \boldsymbol{p}(\lambda)=\mathbf{0} .
$$

So, the only column eigenvector of $C$ is $D_{c}^{-1} \boldsymbol{p}(\lambda)$ and its single row eigenvector is $\boldsymbol{p}(\lambda)^{T} D_{b}^{-1}$.
The most elegant way to find eigenvectors of higher grade is to differentiate the 3 TR as many times as is necessary. Thus, from (4.14),

$$
p_{j+1}^{\prime}(\tau)=\left(\tau-a_{j+1}\right) p_{j}^{\prime}(\tau)+p_{j}(\tau)-b_{j} c_{j} p_{j-1}^{\prime}(\tau), \quad j=1,2, \ldots, n-1
$$

and rearranging the terms we have

$$
\begin{equation*}
b_{j} c_{j} p_{j-1}^{\prime}(\tau)+\left(a_{j+1}-\tau\right) p_{j}^{\prime}(\tau)+p_{j+1}^{\prime}(\tau)=p_{j}(\tau) . \tag{4.16}
\end{equation*}
$$

Then, dividing by $c_{1} \cdots c_{j}$,
$b_{j} \frac{p_{j-1}^{\prime}(\tau)}{c_{0} c_{1} \cdots c_{j-1}}+\left(a_{j+1}-\tau\right) \frac{p_{j}^{\prime}(\tau)}{c_{0} c_{1} \cdots c_{j}}+c_{j+1} \frac{p_{j+1}^{\prime}(\tau)}{c_{1} \cdots c_{j+1}}=\frac{p_{j}(\tau)}{c_{0} c_{1} \cdots c_{j}}, \quad j=0,1, \ldots, n-1$,
where we define $p_{-1}^{\prime}(\tau)=0, b_{n}=c_{n}=1$ and $c_{0}=1$. In matrix terms

$$
(C-\tau I) D_{c}^{-1} \boldsymbol{p}^{\prime}(\tau)=D_{c}^{-1}\left(\boldsymbol{p}(\tau)-\boldsymbol{e}_{n} p_{n}^{\prime}(\tau)\right)
$$

Hence, when $\tau=\lambda$,

$$
(C-\lambda I) D_{c}^{-1} \boldsymbol{p}^{\prime}(\lambda)=D_{c}^{-1} \boldsymbol{p}(\lambda), \quad \text { since } p_{n}^{\prime}(\lambda)=0
$$

and

$$
(C-\lambda I)^{2} D_{c}^{-1} \boldsymbol{p}^{\prime}(\lambda)=(C-\lambda I) D_{c}^{-1} \boldsymbol{p}(\lambda)=\mathbf{0}
$$

since $p_{n}(\lambda)=(\tau-\lambda)^{n}$ and $\lambda$ is a multiple zero of $p_{n}$.
To obtain the next vector differentiate (4.16) again,

$$
b_{j} c_{j} p_{j-1}^{\prime \prime}(\tau)+\left(a_{j+1}-\tau\right) p_{j}^{\prime \prime}(\tau)+1 \cdot p_{j+1}^{\prime \prime}(\tau)=2 p_{j}^{\prime}(\tau)
$$

In order to keep the superdiagonal entries in the Jordan form at the value 1 we divide through by 2 to obtain

$$
(C-\tau I) \frac{1}{2} D_{c}^{-1} \boldsymbol{p}^{\prime \prime}(\tau)=D_{c}^{-1}\left(\boldsymbol{p}^{\prime}(\tau)-\boldsymbol{e}_{n} \frac{1}{2} p_{n}^{\prime \prime}(\tau)\right)
$$

Putting $\tau=\lambda$ and since $p^{\prime \prime}(\lambda)=0$,

$$
(C-\lambda I) \frac{1}{2} D_{c}^{-1} \boldsymbol{p}^{\prime \prime}(\lambda)=D_{c}^{-1} \boldsymbol{p}^{\prime}(\lambda)
$$

and

$$
(C-\lambda I)^{3} \frac{1}{2} D_{c}^{-1} \boldsymbol{p}^{\prime \prime}(\lambda)=(C-\lambda I)^{2} D_{c}^{-1} \boldsymbol{p}^{\prime}(\lambda)=\mathbf{0}
$$

It may be verified that the appropriate definition of $P$ in the confluent case is the unit lower triangular matrix

$$
P_{\lambda}=P=\left[\begin{array}{lllll}
\boldsymbol{p}(\lambda) & \boldsymbol{p}^{\prime}(\lambda) & \frac{1}{2!} \boldsymbol{p}^{\prime \prime}(\lambda) & \ldots & \frac{1}{(n-1)!} \boldsymbol{p}^{(n-1)}(\lambda)
\end{array}\right]
$$

and so

$$
\begin{equation*}
(C-\lambda I) D_{c}^{-1} P=D_{c}^{-1} P N \tag{4.17}
\end{equation*}
$$

The next lemma summarizes this result.

Lemma 4.2.6 If $C$ has one-point spectrum $\lambda$ and $D_{c}$ and $P$ are as defined above, then

$$
C\left(D_{c}^{-1} P\right)=\left(D_{c}^{-1} P\right)(N+\lambda I)=\left(D_{c}^{-1} P\right) J
$$

So, the matrix of generalized right eigenvectors of $C$ is $D_{c}^{-1} P$. To find the row eigenvectors for $C$ we first use

$$
C^{T}=D_{c} D_{b}^{-1} C D_{c}^{-1} D_{b}
$$

to find, analogously to (4.17),

$$
\left(C^{T}-\lambda I\right) D_{b}^{-1} P=D_{b}^{-1} P N
$$

Now, transposing this equation we get

$$
\begin{equation*}
P^{T} D_{b}^{-1} C=\left(\lambda I+N^{T}\right) P^{T} D_{b}^{-1}=I(\lambda I+N) I P^{T} D_{b}^{-1}, \tag{4.18}
\end{equation*}
$$

where

$$
I=\left[\begin{array}{llll} 
& & & \\
& & & 1 \\
& & & 1 \\
& & \ddots & \\
& 1 & & \\
1 & & &
\end{array}\right]
$$

is the reversal or anti-diagonal matrix (all entries $(i, j)$ are zero except when $i+j=n+1$ ). We used the fact that $I I=I$ and $N^{T}=I N I$. Thus, pre-multiplying (4.18) by $I$, we find that

$$
\left(I P^{T} D_{b}^{-1}\right) C=(\lambda I+N)\left(I P^{T} D_{b}^{-1}\right) .
$$

So, $I P^{T} D_{b}^{-1}$ is the matrix of generalized row eigenvectors of $C$.
Summarizing,
Lemma 4.2.7 If $C$ has one-point spectrum $\lambda$ and $D_{b}$ and $P$ are as defined above, then

$$
\left(I P^{T} D_{b}^{-1}\right) C=J\left(I P^{T} D_{b}^{-1}\right) .
$$

Recall that $D_{c}^{-1} P$ is lower triangular and $P^{T} D_{b}^{-1}$ is upper triangular. Nevertheless, it is not true that the product $\left(I P^{T} D_{b}^{-1}\right)\left(D_{c}^{-1} P\right)$ is diagonal, as was the case for simple eigenvalues. The reason is subtle: for a Jordan block, the eigenvectors of grade higher than 1 are not uniquely defined. The phrase "the Jordan basis" that can be found in some text books is incorrect; it is not unique.

Consider the equation above in (4.17),

$$
C D_{c}^{-1} P=D_{c}^{-1} P(\lambda I+N)
$$

Post-multiply by any invertible matrix $\varphi(N), \varphi$ a polynomial, that comutes with $\lambda I+N$ to find that

$$
C D_{c}^{-1} P \varphi(N)=D_{c}^{-1} P(\lambda I+N) \varphi(N)=D_{c}^{-1} P \varphi(N)(\lambda I+N)
$$

Thus $D_{c}^{-1} P$ is only unique up to post-multiplication by any invertible polynomial $\varphi(N)$. And there is no loss in normalizing $\varphi$ to satisfy $\varphi(O)=I$.

Chosen

$$
\left(I P^{T} D_{b}^{-1}\right)\left(D_{c}^{-1} P\right)=\varphi(N)
$$

we have proved

Theorem 4.2.3 If $C$ has one-point spectrum $\lambda$ and $P, D_{b}, D_{c}$ are as defined above then

$$
C=D_{c}^{-1} P \varphi(N)^{-1}(\lambda I+N) I P^{T} D_{b}^{-1}
$$

for some polynomial $\varphi$ with $\varphi(O)=I$.

The example that follows exhibits this feature.
Before, recall that a square matrix $A$ is Toeplitz when the entries of $A$ are constant down the diagonals parallel to the main diagonal and is Hankel when the entries of $A$ are constant along the diagonals perpendicular to the main diagonal.
Z. S. Liu [31] devised an algorithm to obtain one-point spectrum unreduced tridiagonal matrices of arbitrary dimension $n \times n$. These matrices, that we will call Liu's matrices, have only one eigenvalue, zero with algebraic multiplicity $n$ and geometric multiplicity 1. The Jordan form consists of one big Jordan block. We will represent Liu's matrices as

$$
\operatorname{Liu}_{n}=\operatorname{tridiag}\left(\mathbf{1}^{n}, \boldsymbol{\alpha}^{n}, \gamma^{n}\right)
$$

where $\mathbf{1}^{n}$ always stands for a vector of 1 's of length $n-1$. For $n=6$ we have $\boldsymbol{\alpha}^{6}=\left[\begin{array}{llllll}0 & 0 & -1 & 1 & 0 & 0\end{array}\right]$ and $\boldsymbol{\gamma}^{6}=\left[\begin{array}{lllll}-1 & 1 & -1 & 1 & -1\end{array}\right]$.

## Example 4.2.1

Consider the transpose of $\mathrm{Liu}_{6}$ matrix

$$
C=\left[\begin{array}{cccccc}
0 & 1 & & & & \\
-1 & 0 & 1 & & & \\
& 1 & -1 & 1 & & \\
& & -1 & 1 & 1 & \\
& & & 1 & 0 & 1 \\
& & & & -1 & 0
\end{array}\right]
$$

We have

$$
\begin{aligned}
& p_{0}(\tau)=1, \\
& p_{1}(\tau)=\tau, \\
& p_{2}(\tau)=\tau^{2}+1, \\
& p_{3}(\lambda)=(\tau+1) p_{2}(\tau)-p_{1}(\tau)=\tau^{3}+\tau^{2}+1, \\
& p_{4}(\tau)=(\tau-1) p_{3}(\tau)+p_{2}(\tau)=\tau^{4}+\tau, \\
& p_{5}(\tau)=\tau p_{4}(\tau)-p_{3}(\tau)=\tau^{5}-\tau^{3}-1, \\
& p_{6}(\tau)=\tau p_{5}(\tau)+p_{4}(\tau)=\tau^{6} .
\end{aligned}
$$

Then

$$
P=\left[\begin{array}{cccccc}
1 & 0 & 0 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 & 0 \\
1 & 0 & 1 & 0 & 0 & 0 \\
1 & 0 & 1 & 1 & 0 & 0 \\
0 & 1 & 0 & 0 & 1 & 0 \\
-1 & 0 & 0 & -1 & 0 & 1
\end{array}\right], \quad D_{b}=\operatorname{diag}(1,-1,-1,1,1,-1), \quad D_{c}=I
$$

Now, we have

$$
P^{T} D_{b}^{-1} P=\left[\begin{array}{cccccc}
0 & 0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 & 1 \\
0 & 1 & 0 & 0 & 1 & 0 \\
1 & 0 & 0 & 1 & 0 & -1
\end{array}\right]=I \mathcal{U}^{-1}
$$

defining $\mathcal{U}$.
Thus,

$$
U=\left(I P^{T} D_{b}^{-1} P\right)^{-1}=\left[\begin{array}{cccccc}
1 & 0 & 0 & 1 & 0 & -1 \\
0 & 1 & 0 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 & 0 & 1 \\
0 & 0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 & 0 & 1
\end{array}\right]^{-1}=\left[\begin{array}{cccccc}
1 & 0 & 0 & -1 & 0 & 1 \\
0 & 1 & 0 & 0 & -1 & 0 \\
0 & 0 & 1 & 0 & 0 & -1 \\
0 & 0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 & 0 & 1
\end{array}\right]
$$

that is, $\mathcal{U}=I-N^{3}+N^{5}$ is a polynomial in $N$ and thus commutes with $\lambda I+N$ which is also a polynomial in $N$.

Thus $I P^{T} D_{b}^{-1} P=\mathcal{U}^{-1}=I+N^{3}-N^{5}=\varphi(N)$ is unit upper triangular and Toeplitz. So,

$$
C=P \mathcal{U}(0 I+N) I P^{T} D_{b}^{-1}, \quad I=\left(I P^{T} D_{b}^{-1}\right)(P \mathcal{U})
$$

Note that $P \mathcal{U}$ is in LU form and $P^{T} D_{b}^{-1}$ is upper triangular. It will be shown that $C+\sigma I$, $\sigma \neq 0$, can admit the basic LR algorithm with no breakdowns and will converge to

$$
\mathcal{U}(\sigma I+N) \mathcal{U}^{-1}=\sigma I+N . \diamond
$$

### 4.2.3 The general case

In general, unreduced $C$ will have some simple eigenvalues and some multiple ones. The unreduced property implies that $C$ is nonderogatory, meaning that there is only one Jordan block, and thus one eigenvector, for each eigenvalue.

It follows directly from the previous sections that the matrix $P$ has a column $\boldsymbol{p}$ for each simple eigenvalue $\lambda$ and a block of columns $\left[\begin{array}{lllll}\boldsymbol{p}(\lambda) & \boldsymbol{p}^{\prime}(\lambda) & \frac{1}{2!} \boldsymbol{p}^{\prime \prime}(\lambda) & \ldots & \frac{1}{m!} \boldsymbol{p}^{(m)}(\lambda)\end{array}\right]$ if $\lambda$ has multiplicity $m+1$. The only constraint on ordering of columns is that each block must be treated as a whole.

As shown in the section on one-point spectrum matrices, the order of the row eigenvectors must be reversed within each block.

We state without proof.
Theorem 4.2.4 Let $J$ be the (upper) Jordan form of unreduced matrix $C$. Then the spectral decomposition may be written

$$
C=D_{c}^{-1} P \mathcal{U} J \mathcal{W} P^{T} D_{b}^{-1}
$$

where $\mathcal{U}$ is a unit upper triangular matrix that commutes with $J$ and $\mathcal{W}$ is a symmetric permutation matrix that is a direct sum of reversal matrices $I$ conforming to the block structure of $J . \mathcal{W}$ also commutes with $J . P$, and therefore $P^{T}$, are strongly regular.

### 4.3 Convergence of basic LR algorithm on an unreduced tridiagonal

We will show that the assumptions required by Wilkinson in the general case to ensure the convergence of LR algorithm are no longer needed on an unreduced tridiagonal matrix.

First we recall the essencial facts from the beginning of this chapter. Set $C_{1}=C$ and for $k=1,2, \ldots$ define

$$
\begin{aligned}
C_{k} & =\mathrm{L}_{k} R_{k}, \quad \mathrm{~L}_{k} \text { being unit lower triangular } \\
C_{k+1} & =R_{k} \mathrm{~L}_{k} .
\end{aligned}
$$

The LU factorization of $C^{k}$ is

$$
\mathcal{L}_{k} \mathcal{U}_{k}=C^{k}
$$

with $\mathcal{L}_{k}=\mathrm{L}_{1} \mathrm{~L}_{2} \ldots \mathrm{~L}_{k}$ and $\mathcal{U}_{k}=R_{k} R_{i-1} \ldots R_{1}$. And then

$$
C_{k+1}=\mathcal{L}_{k}{ }^{-1} C \mathcal{L}_{k} .
$$

The L factor of a matrix $M$ will be denoted by $\mathcal{L}(M)$, provided that $M$ is strongly regular. In this new notation we will write

$$
\begin{equation*}
C_{k+1}=\mathcal{L}\left(C^{k}\right)^{-1} C \mathcal{L}\left(C^{k}\right) . \tag{4.19}
\end{equation*}
$$

### 4.3.1 Eigenvalues of distinct absolute value

The result in this section is not entirely new but helps to understand the new case.
Let $C$ be a nonsingular unreduced tridiagonal matrix. Without loss of generality, we may write

$$
\begin{aligned}
& C=X \Lambda X^{-1} \\
& \Lambda=\operatorname{diag}\left(\lambda_{1}, \lambda_{2}, \ldots, \lambda_{n}\right), \quad\left|\lambda_{i}\right|>\left|\lambda_{i+1}\right|, \quad \text { for all } i .
\end{aligned}
$$

With the notation of previous section,

$$
X=D_{c}^{-1} P \Delta^{-1}, \quad X^{-1}=P^{T} D_{b}^{-1}
$$

where $P=P_{\Lambda}$ is the polynomial Vandermond matrix given in (4.11). Since $P$ is strongly regular (and also $X$ and $X^{-1}$ by theorem 4.2.2) we may write

$$
P=L D U,
$$

where $L$ is unit lower triangular, $D$ is diagonal and $U$ is unit upper triangular. Then we can manipulate $C^{k}$ into LU form as follows

$$
\begin{align*}
C^{k} & =X \Lambda^{k} X^{-1} \\
& =D_{c}^{-1} P \Delta^{-1} \Lambda^{k} P^{T} D_{b}^{-1} \\
& =D_{c}^{-1} L D U \Delta^{-1} \Lambda^{k} U^{T} D L^{T} D_{b}^{-1} \\
& =D_{c}^{-1} L D U \Delta^{-1}\left(\Lambda^{k} U^{T} \Lambda^{-k}\right) \Lambda^{k} D L^{T} D_{b}^{-1} \\
& =\left(D_{c}^{-1} L D_{c}\right)\left(D_{c}^{-1} D U \Delta^{-1}\right)\left(I+E_{k}\right)\left(\Lambda^{k} D L^{T} D_{b}^{-1}\right) \\
& =\left(D_{c}^{-1} L D_{c}\right)\left(I+F_{k}\right)\left(D_{c}^{-1} D U \Delta^{-1}\right)\left(\Lambda^{k} D L^{T} D_{b}^{-1}\right) \tag{4.20}
\end{align*}
$$

where

$$
F_{k}=\left(D_{c}^{-1} D U \Delta^{-1}\right) E_{k}\left(D_{c}^{-1} D U \Delta^{-1}\right)^{-1}
$$

Notice that

$$
\begin{align*}
& X=\left(D_{c}^{-1} L D_{c}\right)\left(D_{c}^{-1} D U \Delta^{-1}\right)=: L_{X} U_{X}  \tag{4.21}\\
& Y=U^{T}\left(D L^{T} D_{b}^{-1}\right)=: L_{Y} U_{Y}
\end{align*}
$$

are the LU factorizations of $X$ and $Y$, respectively.
Also, we had written

$$
\Lambda^{k} U^{T} \Lambda^{-k}=I+E_{k}
$$

So, $E_{k}$ is strictly lower triangular and, if we let $U=\left(u_{i j}\right)$ and $E_{k}=\left(e_{i j}^{(k)}\right)$, we have

$$
e_{i j}^{(k)}= \begin{cases}u_{j i}\left(\frac{\lambda_{i}}{\lambda_{j}}\right)^{k}, & i>j  \tag{4.22}\\ 0, & i \leq j\end{cases}
$$

Thus the $(j+m, j)$ entry of $E_{k}$ is $u_{j, j+m}\left(\frac{\lambda_{j+m}}{\lambda_{j}}\right)^{k}$ and tends to zero as $k$ tends to infinity, since $\left|\lambda_{j+m}\right|<\left|\lambda_{j}\right|$. Hence

$$
\left\|F_{k}\right\| \leq \operatorname{cond}\left(D_{c}^{-1} D U \Delta^{-1}\right)\left\|E_{k}\right\| \rightarrow 0 \quad \text { as } \quad k \rightarrow \infty .
$$

It seems likely that $I+F_{k}$ is strongly regular for all $k$ but it certainly holds for large enough $k$, say

$$
I+F_{k}=L_{k} U_{k} \quad \text { with } \quad L_{k}, U_{k} \rightarrow I \quad \text { as } \quad k \rightarrow \infty .
$$

Thus, for large enough $k$, from (4.20),

$$
C^{k}=\left(D_{c}^{-1} L D_{c}\right)\left(L_{k} U_{k}\right)\left(D_{c}^{-1} D U \Delta^{-1}\right)\left(\Lambda^{k} D L^{T} D_{b}^{-1}\right),
$$

and then

$$
\mathcal{L}\left(C^{k}\right)=D_{c}^{-1} L D_{c} L_{k}
$$

since $U_{k}\left(D_{c}^{-1} D U \Delta^{-1}\right)\left(\Lambda^{k} D L^{T} D_{b}^{-1}\right)$ is upper triangular. Thus

$$
\mathcal{L}\left(C^{k}\right) \rightarrow D_{c}^{-1} L D_{c}=L_{X} \quad \text { as } \quad k \rightarrow \infty .
$$

Finally, from (4.19),

$$
C_{k+1}=\mathcal{L}\left(C^{k}\right)^{-1} C \mathcal{L}\left(C^{k}\right)
$$

and thus

$$
\begin{aligned}
C_{k+1} \rightarrow & \left(D_{c}^{-1} L D_{c}\right)^{-1} X \Lambda X^{-1}\left(D_{c}^{-1} L D_{c}\right) \\
& =L_{X}^{-1} L_{X} U_{X} \Lambda U_{X}^{-1} L_{X}^{-1} L_{X} \\
& =U_{X} \Lambda U_{X}^{-1} \quad \text { as } \quad k \rightarrow \infty,
\end{aligned}
$$

showing that $C_{k+1}$ converges to an upper triangular matrix with diagonal elements equal to $\lambda_{1}, \ldots, \lambda_{n}$.

Recalling that the tridiagonal form is preserved by the LR algorithm, we have proved Theorem 4.3.1 Let $C$ be a nonsingular unreduced tridiagonal matrix with distinct eigenvalues $\lambda_{i}, i=1, \ldots, n$. Given the notation above,

$$
\lim C_{k}=C_{\infty}
$$

exists and is upper bidiagonal with the $i^{\text {th }}$ diagonal entry of $C_{\infty}$ equal to $\lambda_{i}$.
We should note that we are assuming that $C_{k}$ permits triangular decomposition at all stages of the LR algorithm.

### 4.3.2 One-point spectrum

Recall that the Vandermond matrix $P$ for this case is unit lower triangular,

$$
P=\left[\begin{array}{lllll}
\boldsymbol{p}(\lambda) & \boldsymbol{p}^{\prime}(\lambda) & \frac{1}{2!} \boldsymbol{p}^{\prime \prime}(\lambda) & \ldots & \frac{1}{(n-1)!} \boldsymbol{p}^{(n-1)}(\lambda)
\end{array}\right]
$$

From theorem 4.2.3,

$$
\begin{equation*}
C=X(\lambda I+N) X^{-1} \tag{4.23}
\end{equation*}
$$

with

$$
X=D_{c}^{-1} P \varphi(N)^{-1} \quad \text { and } \quad X^{-1}=I P^{T} D_{b}^{-1}
$$

Then

$$
C^{k}=D_{c}^{-1} P \varphi(N)^{-1}(\lambda I+N)^{k} I P^{T} D_{b}^{-1}
$$

Note that $P^{T}$ is unit upper triangular. Next we invoke the following lemma.

Lemma 4.3.1 For all $k \geq n,(\lambda I+N)^{k} I$ for $\lambda \neq 0$ is strongly regular and thus admits triangular factorization, say

$$
(\lambda I+N)^{k} I=L_{k} U_{k}=L_{k} D_{k} L_{k}^{T}
$$

and, as $k \rightarrow \infty, L_{k}=I+E_{k}, E_{k} \rightarrow O$. The rate of convergence is low, $\mathcal{O}(1 / k)$.

The proof of this lemma will be given later.
Now, we can factor $C^{k}, k \geq n$ :

$$
\begin{aligned}
C^{k} & =D_{c}^{-1} P D_{c} D_{c}^{-1} \varphi(N)^{-1}\left(I+E_{k}\right) U_{k} P^{T} D_{b}^{-1} \\
& =D_{c}^{-1} P D_{c}\left(I+F_{k}\right) D_{c}^{-1} \varphi(N)^{-1} U_{k} P^{T} D_{b}^{-1}
\end{aligned}
$$

with

$$
F_{k}=\left(\varphi(N) D_{c}\right)^{-1} E_{k}\left(\varphi(N) D_{c}\right) \rightarrow O \quad \text { as } \quad k \rightarrow \infty
$$

since

$$
\left\|F_{k}\right\| \leq \operatorname{cond}\left(\varphi(N) D_{c}\right)\left\|E_{k}\right\| \rightarrow 0 \quad \text { as } \quad k \rightarrow \infty
$$

Thus,

$$
\mathcal{L}\left(C^{k}\right)=D_{c}^{-1} P D_{c} \mathcal{L}\left(I+F_{k}\right) \rightarrow D_{c}^{-1} P D_{c} \quad \text { as } \quad k \rightarrow \infty
$$

since $D_{c}^{-1} \varphi(N)^{-1} U_{k} P^{T} D_{b}^{-1}$ is upper triangular $\left(\varphi(N)^{-1}\right.$ is upper triangular and Toeplitz).
Finally, notice that the LU factorization of $X$ is

$$
X=\left(D_{c}^{-1} P D_{c}\right)\left(D_{c}^{-1} \varphi(N)^{-1}\right):=L_{X} U_{X}
$$

and then

$$
\begin{align*}
C_{k+1} & =\mathcal{L}\left(C^{k}\right)^{-1} C \mathcal{L}\left(C^{k}\right) \\
& =\mathcal{L}\left(C^{k}\right)^{-1} X(\lambda I+N) X^{-1} \mathcal{L}\left(C^{k}\right)  \tag{by4.23}\\
& \rightarrow\left(D_{c}^{-1} P D_{c}\right)^{-1} X(\lambda I+N) X^{-1}\left(D_{c}^{-1} P D_{c}\right) \\
& =L_{X}^{-1} L_{X} U_{X}(\lambda I+N) U_{X}^{-1} L_{X}^{-1} L_{X} \\
& =U_{X}(\lambda I+N) U_{X}^{-1} \\
& =D_{c}^{-1} \varphi(N)^{-1}(\lambda I+N) \varphi(N) D_{c} \\
& =D_{c}^{-1}(\lambda I+N) D_{c}
\end{align*}
$$

since $\varphi(N)$ commutes with $(\lambda I+N)$. Notice that $D_{c}^{-1}(\lambda I+N) D_{c}$ is upper bidiagonal with diagonal entries equal to $\lambda$.

We have just proved

Theorem 4.3.2 Let $C$ be a nonsingular unreduced tridiagonal matrix that permits triangular factorization and has a one-point spectrum $\lambda$. Given the notation above, the basic LR algorithm applied to $C$ produces a sequence of matrices $C_{k}$ that converges (in exact arithmetic) to

$$
D_{c}^{-1}(\lambda I+N) D_{c}
$$

with $D_{c}$ defined above lemma 4.2.3 (page 91).

Proof of lemma 4.3.1:
The proof rests on the form of powers of Jordan blocks. These are Toeplitz matrices and involve binomial coefficients. We have (see [29, p.138])

$$
(\lambda I+N)^{k}=\sum_{i=0}^{k}\binom{k}{i} \lambda^{k-i} N^{i} .
$$

If $J$ is $n \times n$ then all terms $N^{i}=O$ for $i \geq n$. Then, for $k \geq n$,

$$
(\lambda I+N)^{k}=\sum_{i=0}^{n-1}\binom{k}{i} \lambda^{k-i} N^{i} .
$$

So $(\lambda I+N)^{k} I$ is the Hankel matrix

$$
\left[\begin{array}{ccccccc}
\binom{k}{n-1} \lambda^{k-(n-1)} & \left.\begin{array}{c}
k \\
n-2
\end{array}\right) \lambda^{k-(n-2)} & \binom{k}{n-3} \lambda^{k-(n-3)} & \ldots & \binom{k}{2} \lambda^{k-2} & k \lambda^{k-1} & \lambda^{k} \\
\binom{k}{n-2} \lambda^{k-(n-2)} & \binom{k}{n-3} \lambda^{k-(n-3)} & \cdots & \binom{k}{2} \lambda^{k-2} & k \lambda^{k-1} & \lambda^{k} & 0 \\
\binom{k}{n-3} \lambda^{k-(n-3)} & \cdots & \binom{k}{2} \lambda^{k-2} & k \lambda^{k-1} & \lambda^{k} & 0 & 0 \\
\cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\
\binom{k}{2} \lambda^{k-2} & k \lambda^{k-1} & \lambda^{k} & 0 & 0 & \cdots & 0 \\
k \lambda^{k-1} & \lambda^{k} & 0 & 0 & \cdots & 0 & 0 \\
\lambda^{k} & 0 & 0 & \cdots & 0 & 0 & 0
\end{array}\right] .
$$

As long as $\lambda \neq 0$, it can be factored out of all the powers and is just a scalar converging either to 0 or to $\infty$ attached to $U_{k}$ but it does not alter $L_{k}$ (if $A=L U$ then $\alpha A=L(\alpha U)$ ). Thus, let $\Delta_{\lambda}$ be the matrix defined as

$$
\Delta_{\lambda}:=\operatorname{diag}\left(1, \lambda, \lambda^{2}, \ldots, \lambda^{n-1}\right)
$$

and factor out $\lambda^{k-n+1}$ to get

$$
(\lambda I+N)^{k} I=\lambda^{k-n+1} \Delta_{\lambda} \widetilde{H}_{k} \Delta_{\lambda}
$$

where

$$
\widetilde{H}_{k}:=\left[\begin{array}{ccccccc}
\binom{k}{n-1} & \binom{k}{n-2} & \binom{k}{n-3} & \ldots & \binom{k}{2} & k & 1 \\
\binom{k}{n-2} & \binom{k}{n-3} & \ldots & \binom{k}{2} & k & 1 & 0 \\
\binom{k}{n-3} & \ldots & \binom{k}{2} & k & 1 & 0 & 0 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
\binom{k}{2} & k & 1 & 0 & 0 & \ldots & 0 \\
k & 1 & 0 & 0 & \ldots & 0 & 0 \\
1 & 0 & 0 & \ldots & 0 & 0 & 0
\end{array}\right] .
$$

So, factoring out $\lambda^{k-n+1}$ and with a diagonal scaling we are free of the powers of $\lambda$ and will concentrate on $\widetilde{H}_{k}$.

For $i+j \leq n+1$, as a function of $k$, the $(i, j)$ entry of $\widetilde{H}_{k}$ is a polynomial of degree $n+1-i-j$. Since $\widetilde{H}_{k}$ is strongly regular (which will be proved below), it has an LU factorization. If we say $\widetilde{H}_{k}=\widetilde{L}_{k} \widetilde{U}_{k}$, it may be verified that

$$
\widetilde{L}_{k}=\left[\begin{array}{ccccccc}
1 & & & & & & \\
\mathcal{O}\left(\frac{1}{k}\right) & 1 & & & & & \\
\mathcal{O}\left(\frac{1}{k^{2}}\right) & \mathcal{O}\left(\frac{1}{k}\right) & 1 & & & & \\
\cdots & \cdots & \ldots & \ldots & & & \\
\mathcal{O}\left(\frac{1}{k^{n-3}}\right) & \ldots & \mathcal{O}\left(\frac{1}{k^{2}}\right) & \mathcal{O}\left(\frac{1}{k}\right) & 1 & & \\
\mathcal{O}\left(\frac{1}{k^{n-2}}\right) & \mathcal{O}\left(\frac{1}{k^{n-3}}\right) & \ldots & \mathcal{O}\left(\frac{1}{k^{2}}\right) & \mathcal{O}\left(\frac{1}{k}\right) & 1 & \\
\mathcal{O}\left(\frac{1}{k^{n-1}}\right) & \mathcal{O}\left(\frac{1}{k^{n-2}}\right) & \mathcal{O}\left(\frac{1}{k^{n-3}}\right) & \ldots & \mathcal{O}\left(\frac{1}{k^{2}}\right) & \mathcal{O}\left(\frac{1}{k}\right) & 1
\end{array}\right] .
$$

So, for $p \geq 1$, the $(j+p, j)$ entry of $\widetilde{L}_{k}$ is $\mathcal{O}\left(\left(\frac{1}{k}\right)^{p}\right)$ and thus

$$
\begin{equation*}
\widetilde{L}_{k} \rightarrow I+E_{k}, \quad E_{k} \rightarrow O \quad \text { as } \quad k \rightarrow \infty . \tag{4.24}
\end{equation*}
$$

But the convergence is very slow, governed by the terms of $\mathcal{O}\left(\frac{1}{k}\right)$.
Since $\widetilde{H}_{k}$ is symmetric we can write $L_{k} U_{k}=L_{k} D_{k} L_{k}^{T}$. The diagonal matrix $D_{k}$ is also a function of $k$ but it does not converge to a finite matrix. Nevertheless, for each $k$,
$(\lambda I+N)^{k} I$ has an LU factorization. We have

$$
\begin{aligned}
(\lambda I+N)^{k} I & =\lambda^{k-n+1} \Delta_{\lambda} \widetilde{H}_{k} \Delta_{\lambda} \\
& =\lambda^{k-n+1} \Delta_{\lambda} \widetilde{L}_{k} \widetilde{U}_{k} \Delta_{\lambda} \\
& =\Delta_{\lambda} \widetilde{L}_{k} \Delta_{\lambda}^{-1} \lambda^{k-n+1} \Delta_{\lambda} \widetilde{U}_{k} \Delta_{\lambda} \\
& =L_{k} U_{k}
\end{aligned}
$$

with $L_{k} \equiv \Delta_{\lambda} \widetilde{L}_{k} \Delta_{\lambda}^{-1}$ and $U_{k} \equiv \lambda^{k-n+1} \Delta_{\lambda} \widetilde{U}_{k} \Delta_{\lambda}$. Since (4.24) occur we will also have

$$
L_{k} \rightarrow I+G_{k}, \quad G_{k} \rightarrow O \quad \text { as } \quad k \rightarrow \infty
$$

with $G_{k}=\Delta_{\lambda} E_{k} \Delta_{\lambda}^{-1}$.
In what follows we will prove that the matrix $\widetilde{H}_{k}$ is strongly regular.

Lemma 4.3.2 $\widetilde{H}_{k}$ is strongly regular.

Proof: As noticed before, as a function of $k$, the $(i, j)$ entry of $\widetilde{H}_{k}, i+j \leq n+1$, is a polynomial of degree $n+1-(i+j)$. It is given by

$$
\begin{aligned}
\binom{k}{n+1-i-j} & =\frac{k(k-1) \cdots(k-(n-i-j))}{(n+1-i-j)!} \\
& =\frac{(k)_{n+1-i-j}}{(n+1-i-j)!}
\end{aligned}
$$

where $(k)_{m+1}$ is the Pochammer symbol defined as

$$
(k)_{0}:=1, \quad(k)_{m+1}:=k(k-1) \cdots(k-m), \quad m \in \mathbb{N}_{0}
$$

So, in each entry $(i, j), i+j \leq n+1$, the dominant term is

$$
\frac{k^{n+1-i-j}}{(n+1-i-j)!}
$$

and thus $\widetilde{H}_{k}=\widehat{H}_{k}$ plus lower order terms with

$$
\widehat{H}_{k}:=\left[\begin{array}{ccccccc}
\frac{k^{n-1}}{(n-1)!} & \frac{k^{n-2}}{(n-2)!} & \frac{k^{n-3}}{(n-3)!} & \ldots & \frac{k^{2}}{2!} & k & 1 \\
\frac{k^{n-2}}{(n-2)!} & \frac{k^{n-3}}{(n-3)!} & \ldots & \frac{k^{2}}{2!} & k & 1 & 0 \\
\frac{k^{n-3}}{(n-3)!} & \ldots & \frac{k^{2}}{2!} & k & 1 & 0 & 0 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
\frac{k^{2}}{2!} & k & 1 & 0 & 0 & \ldots & 0 \\
k & 1 & 0 & 0 & \ldots & 0 & 0 \\
1 & 0 & 0 & \ldots & 0 & 0 & 0
\end{array}\right] .
$$

In order to eliminate $k$ from $\widehat{H}_{k}$ we define

$$
\Delta_{k}:=\operatorname{diag}\left(1, k, k^{2}, \ldots, k^{n-1}\right)
$$

and we have

$$
I \Delta_{k} I=\operatorname{diag}\left(k^{n-1}, \ldots, k^{2}, k, 1\right) .
$$

Now it can be verified that

$$
\widehat{H}_{k}=\Delta_{k}^{-1} F_{n}^{(n-1)} I \Delta_{k} I
$$

with $F_{n}^{(n-1)}$ given by

$$
F_{n}^{(n-1)}=\left[\begin{array}{ccccccc}
\frac{1}{(n-1)!} & \frac{1}{(n-2)!} & \frac{1}{(n-3)!} & \ldots & \frac{1}{2!} & 1 & 1 \\
\frac{1}{(n-2)!} & \frac{1}{(n-3)!} & \ldots & \frac{1}{2!} & 1 & 1 & 0 \\
\frac{1}{(n-3)!} & \ldots & \frac{1}{2!} & 1 & 1 & 0 & 0 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
\frac{1}{2!} & 1 & 1 & 0 & 0 & \ldots & 0 \\
1 & 1 & 0 & 0 & \ldots & 0 & 0 \\
1 & 0 & 0 & \ldots & 0 & 0 & 0
\end{array}\right] .
$$

To prove that $\widetilde{H}_{k}$ is completely regular it suffices to prove that $F_{n}^{(n-1)}$ has this property (see lemma 4.2.4, page 93). First observe that $F_{n}^{(n-1)}$ is nonsingular since $\operatorname{det}\left(F_{n}^{(n-1)}\right)=(-1)^{n-1}$. Now we still have to show that all the $l^{\text {th }}$ leading principal
minors of $F_{n}^{(n-1)}$ are nonzero, $l=1, \ldots, n-1$. This follows immediately from the next lemma.

Recall that the factorial of a positive integer is defined as

$$
m!=m(m-1)(m-2) \cdots 2.1, \quad m \in \mathbb{N}
$$

By convention $0!=1$. We will say that for a negative integer $m, m!=0$. Also we define the double factorial symbol (!!) as follows

$$
\begin{aligned}
m!! & =m(m-1)!(m-2)!\cdots 2!1!, \quad m \in \mathbb{N} \\
0!! & =1 .
\end{aligned}
$$

This definition is not universal.
Define the $j \times j$ Hankel matrix

$$
F_{j}^{(n-1)}:=\left[\begin{array}{ccccc}
\frac{1}{(n-1)!} & \frac{1}{(n-2)!} & \frac{1}{(n-3)!} & \cdots & \frac{1}{(n-j)!} \\
\frac{1}{(n-2)!} & \frac{1}{(n-3)!} & \cdots & \cdots & \frac{1}{(n-j-1)!} \\
\frac{1}{(n-3)!} & \cdots & \cdots & \cdots & \frac{1}{(n-j-2)!} \\
\cdots & \cdots & \cdots & \cdots & \cdots \\
\frac{1}{(n-j)!} & \frac{1}{(n-j-1)!} & \frac{1}{(n-j-2)!} & \cdots & \frac{1}{(n-2 j+1)!}
\end{array}\right] .
$$

The following lemma gives us a useful result to calculate the determinant of $F_{j}^{(n-1)}$. It seems likely that this result may be already known but we couldn't find it in a search of the literature.

Lemma 4.3.3 Let $j<n$. Then

$$
\operatorname{det}\left(F_{j}^{(n-1)}\right)=(-1)^{j(j-1) / 2} \frac{(j-1)!!(n-j-1)!!}{(n-1)!!} .
$$

Proof. There are two cases to consider: first, $j \leq \frac{n+1}{2}$ and, second, $n>j>\frac{n+1}{2}$.
Case 1: $j \leq \frac{n+1}{2}$
Factor out $\frac{1}{(n-l)!}$ from column $l, l=1,2, \ldots, j$, to obtain the new integer matrix

$$
\widetilde{F}_{j}^{(n-1)}:=\left[\begin{array}{cccc}
1 & 1 & \cdots & 1 \\
n-1 & n-2 & \cdots & n-j \\
(n-1)(n-2) & (n-2)(n-3) & \cdots & (n-j)(n-j-1) \\
\vdots & \vdots & \vdots & \vdots \\
(n-1) \cdots(n-j+1) & (n-2) \cdots(n-j) & \cdots & (n-j) \cdots(n-2 j+2)
\end{array}\right]
$$

with

$$
\operatorname{det}\left(F_{j}^{(n-1)}\right)=\frac{1}{(n-1)!} \frac{1}{(n-2)!} \cdots \frac{1}{(n-j)!} \operatorname{det}\left(\widetilde{F}_{j}^{(n-1)}\right)
$$

By subtracting suitable multiples of higher rows from lower rows we are left with the transpose of the standard Vandermond Matrix (see (4.12), page 90)

$$
V_{\{n-1, n-2, \ldots, n-j\}}=: V_{j} .
$$

In fact, if we use the Pochammer symbol defined above, it is clear that $\widetilde{F}_{j}^{(n-1)}$ is a polynomial Vandermond matrix,

$$
\widetilde{F}_{j}^{(n-1)}=\left[\begin{array}{cccc}
1 & 1 & \ldots & 1  \tag{4.25}\\
(n-1)_{1} & (n-2)_{1} & \ldots & (n-j)_{1} \\
(n-1)_{2} & (n-2)_{2} & \ldots & (n-j)_{2} \\
\vdots & \vdots & \vdots & \vdots \\
(n-1)_{j-1} & (n-2)_{j-1} & \ldots & (n-j)_{j-1}
\end{array}\right]
$$

and the determinant of any polynomial Vandermond matrix is just the determinant of the standard Vandermond matrix (see lemma 4.2.2, page 91).

To be more precise, if we denote the $k^{\text {th }}$ row of $\widetilde{F}_{j}^{(n-1)}$ by $\operatorname{Row}_{l}, l=1, \ldots, j$, the row operations we have to perform are

$$
\operatorname{Row}_{k} \leftarrow \operatorname{Row}_{k}+(k-2) \operatorname{Row}_{k-1}+(k-3)(n-1) \operatorname{Row}_{k-2}+\ldots+1 .(n-1)^{k-3} \operatorname{Row}_{2},
$$

for $k=n, n-1, \ldots, 3$ (note that we start from the bottom). We will end with the leading (monic) terms

$$
\left[\begin{array}{cccc}
1 & 1 & \ldots & 1  \tag{4.26}\\
n-1 & n-2 & \ldots & n-j \\
(n-1)^{2} & (n-2)^{2} & \ldots & (n-j)^{2} \\
\vdots & \vdots & \vdots & \vdots \\
(n-1)^{j-1} & (n-2)^{j-1} & \ldots & (n-j)^{j-1}
\end{array}\right] .
$$

Hence,

$$
\operatorname{det}\left(F_{j}^{(n-1)}\right)=\frac{1}{(n-1)!} \frac{1}{(n-2)!} \cdots \frac{1}{(n-j)!} \operatorname{det}\left(V_{j}\right) .
$$

But, from (4.13), page 91, we know that

$$
\begin{aligned}
\operatorname{det}\left(V_{j}\right) & =\prod_{j \geq k>i \geq 1}((n-k)-(n-i)) \\
& =\prod_{j \geq k>i \geq 1}(i-k) \\
& =(-1)^{j(j-1) / 2} \prod_{j \geq k>i \geq 1}(k-i) \\
& =(-1)^{j(j-1) / 2}(j-1)!(j-2)!\cdots 2!1!.
\end{aligned}
$$

Notice that $j(j-1) / 2$ is the number of terms in the product $\prod_{j \geq k>i \geq 1}(i-k)$.
Finally, we can write

$$
\begin{align*}
\operatorname{det}\left(F_{j}^{(n-1)}\right) & =(-1)^{j(j-1) / 2} \frac{(j-1)!}{(n-1)!} \frac{(j-2)!}{(n-2)!} \cdots \frac{0!}{(n-j)!} \\
& =(-1)^{j(j-1) / 2} \frac{(j-1)!(n-j-1)!!}{(n-1)!!}, \quad n \geq 2 j-1 . \tag{4.27}
\end{align*}
$$

We have established the result for $n \geq 2 j-1$, that is, for $j \leq \frac{n+1}{2}$.

Case 2: Now consider the case $n>j>\frac{n+1}{2}$. In this case we will have

$$
F_{j}^{(n-1)}=\left[\begin{array}{cccccccccc}
\frac{1}{(n-1)!} & \frac{1}{(n-2)!} & \frac{1}{(n-3)!} & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \frac{1}{(n-j)!} \\
\frac{1}{(n-2)!} & \frac{1}{(n-3)!} & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \frac{1}{(n-j-1)!} \\
\frac{1}{(n-3)!} & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \frac{1}{(n-j-2)!} \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & 1 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & 1 & 1 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & 1 & 1 & 0 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
\ldots & \ldots & \ldots & \ldots & \ldots & 1 & 1 & 0 & \ldots & 0 \\
\frac{1}{(n-j)!} & \frac{1}{(n-j-1)!} & \frac{1}{(n-j-2)!} & \ldots & 1 & 1 & 0 & \ldots & 0 & 0
\end{array}\right] .
$$

Again, if we factor out $\frac{1}{(n-l)!}$ for column $l, l=1,2, \ldots, j$, we are left with the polynomial Vandermond matrix with the Pochammer symbol as in (4.25). Its just the case that some values of the Pochammer symbol are equal to zero. So, some zeros appear in the lower right corner of $\widetilde{F}_{j}^{(n-1)}$ but it is still valid that

$$
\operatorname{det}\left(\widetilde{F}_{j}^{(n-1)}\right)=\operatorname{det}\left(V_{j}\right)
$$

and, then, formula (4.27) also applies to the case $n>j>\frac{n+1}{2}$.
Observe that from the symmetry of the formula (4.27) it follows, up to the sign,

$$
\operatorname{det}\left(F_{j}^{(n-1)}\right)=\operatorname{det}\left(F_{n-j}^{(n-1)}\right) .
$$

## Chapter 5

## Triple dqds algorithm

In this chapter we describe the derivation of a first version of the triple dqds - an algorithm that performs implicitly three steps of simple dqds keeping real arithmetic in the presence of complex shifts. A preliminary version was developed by Z. Wu [64]. We start by describing the connection to the implicit double shifted LR algorithm and then go into the details of a practical implementation. The relation between dqds and the Gram-Schmidt orthogonalization process can be used to establish new results about triple dqds. We defer this study to the next chapter. For a discussion of the relation of LR to dqds see Fernando and Parlett [16].

### 5.1 Triple dqds algorithm

The essential point of triple dqds algorithm is the conversion of the implicit double shifted LR algorithm (see page 65) into a dqds format. There are two main reasons to pursue this goal - first, there is both theoretical and practical evidence that the eigenvalues are usually better defined by the entries of $L$ and $U$ rather than by the entries of the product $J=L U$, and, second, dqds produces with no cost the quantities $d_{i}, i=1,2, \ldots, n$, (see page 75 ) that may be used in a more efficient shift strategy than the one used by the LR algorithm.

Let $L_{1}$ and $U_{1}$ be real and have the form presented in section 3.5.1 (page 72),

$$
L_{1}=\left[\begin{array}{cccccc}
1 & & & & & \\
l_{1} & 1 & & & & \\
& l_{2} & 1 & & & \\
& & \ddots & \ddots & & \\
& & & l_{n-2} & 1 & \\
& & & & l_{n-1} & 1
\end{array}\right], \quad U_{1}=\left[\begin{array}{cccccc}
u_{1} & 1 & & & & \\
& u_{2} & 1 & & & \\
& & u_{3} & 1 & & \\
& & & \ddots & \ddots & \\
& & & & u_{n-1} & 1 \\
& & & & & u_{n}
\end{array}\right] .
$$

Consider that we apply three dqds steps to $L_{1}$ and $R_{1}$. The results of shifting by $\sigma_{1}, \sigma_{2}$ and $\sigma_{3}$ in succession are

$$
\begin{align*}
& U_{1} L_{1}-\sigma_{1} I=L_{2} U_{2}  \tag{5.1}\\
& U_{2} L_{2}-\sigma_{2} I=L_{3} U_{3}  \tag{5.2}\\
& U_{3} L_{3}-\sigma_{3} I=L_{4} U_{4} \tag{5.3}
\end{align*}
$$

The following lemma relates $L_{4}$ and $U_{4}$ to $L_{1}$ and $U_{1}$.
Lemma 5.1.1 Consider that matrices $L_{4}$ and $U_{4}$ result from the application of three dqds steps with successive shifts $\sigma_{1}, \sigma_{2}$ and $\sigma_{3}$. Then

$$
L_{4} U_{4}=\mathcal{L}^{-1} U_{1} L_{1} \mathcal{L}-\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right) I
$$

where $\mathcal{L}=L_{2} L_{3}$.
Proof. We have

$$
\begin{align*}
L_{4} U_{4} & =U_{3} L_{3}-\sigma_{3} I \\
& =L_{3}^{-1}\left(L_{3} U_{3}\right) L_{3}-\sigma_{3} I \\
& =L_{3}^{-1}\left(U_{2} L_{2}-\sigma_{2} I\right) L_{3}-\sigma_{3} I  \tag{5.4}\\
& =L_{3}^{-1}\left(L_{2}^{-1}\left(U_{1} L_{1}-\sigma_{1} I\right) L_{2}-\sigma_{2}\right) L_{3}-\sigma_{3} I  \tag{5.5}\\
& =\left(L_{2} L_{3}\right)^{-1} U_{1} L_{1}\left(L_{2} L_{3}\right)-\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right) I \\
& =\mathcal{L}^{-1} U_{1} L_{1} \mathcal{L}-\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right) I,
\end{align*}
$$

where $\mathcal{L} \equiv L_{2} L_{3}$.

Observe that

$$
\mathcal{L}^{-1} U_{1} L_{1} \mathcal{L}-\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right) I=\mathcal{L}^{-1}\left(U_{1} L_{1}-\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right) I\right) \mathcal{L}
$$

to realize that lemma 5.1 .1 just says that matrices $L_{4} U_{4}$ and $U_{1} L_{1}-\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right) I$ are similar.

Now consider the matrix $\mathcal{U} \equiv U_{3} U_{2}$. Next lemma tells us about the LU decomposition $\mathcal{L U}$.

Lemma 5.1.2 Let $\mathcal{L}=L_{2} L_{3}$ and $\mathcal{U}=U_{3} U_{2}$. Then

$$
\mathcal{L U}=\left(U_{1} L_{1}\right)^{2}-\left(2 \sigma_{1}+\sigma_{2}\right) U_{1} L_{1}+\left(\sigma_{1} \sigma_{2}+\sigma_{1}^{2}\right) I
$$

Proof. We have

$$
\begin{aligned}
\mathcal{L U} & =\left(L_{2} L_{3}\right)\left(U_{3} U_{2}\right) \\
& =L_{2}\left(L_{3} U_{3}\right) U_{2} \\
& =L_{2}\left(U_{2} L_{2}-\sigma_{2} I\right) U_{2} \\
& =\left(L_{2} U_{2}\right)^{2}-\sigma_{2} L_{2} U_{2} \\
& =\left(U_{1} L_{1}-\sigma_{1} I\right)^{2}-\sigma_{2}\left(U_{1} L_{1}-\sigma_{1} I\right) \\
& =\left(U_{1} L_{1}\right)^{2}-\left(2 \sigma_{1}+\sigma_{2}\right) U_{1} L_{1}+\left(\sigma_{1} \sigma_{2}+\sigma_{1}^{2}\right) I
\end{aligned}
$$

Next we will show how to do a triple shift in order to keep real arithmetic in the presence of complex eigenvalues. That is, we will see how to combine three consecutive dqds iterations with complex shifts $\sigma_{1}, \sigma_{2}$ and $\sigma_{3}$ such that the resulting matrices $L_{4}$ and $U_{4}$ after this triple shift will be again real.

Matrices $\mathcal{L}=L_{2} L_{3}$ and $\mathcal{U}=U_{3} U_{2}$ will have real entries if we choose

$$
\sigma_{2}=-2\left(\Im \sigma_{1}\right) \mathbf{i}
$$

where $\Im \sigma_{1}$ denotes the imaginary part of $\sigma_{1}$ and $\mathbf{i}$ is the imaginary unit. This is the only nontrivial complex solution to

$$
\left\{\begin{array}{l}
2 \sigma_{1}+\sigma_{2} \in \mathbb{R} \\
\sigma_{1} \sigma_{2}+\sigma_{1}^{2} \in \mathbb{R}
\end{array}\right.
$$

It is easy to verify that this choice for $\sigma_{2}$ is such that

$$
2 \sigma_{1}+\sigma_{2}=2\left(\Re \sigma_{1}\right) \quad \text { and } \quad \sigma_{1} \sigma_{2}+\sigma_{1}^{2}=\left|\sigma_{1}\right|^{2}
$$

where $\Re \sigma_{1}$ denotes the real part of $\sigma_{1}$. Thus, according to lemma 5.1.2 above,

$$
\mathcal{L U}=\left(U_{1} L_{1}\right)^{2}-2\left(\Re \sigma_{1}\right) U_{1} L_{1}+\left|\sigma_{1}\right|^{2} I,
$$

which is the same matrix that appears in the double shifted LR algorithm if we consider $A_{1} \equiv U_{1} L_{1}$ (see lemma 3.4.1, page 66 ).

By lemma 5.1.1, the accumulated shift is then

$$
\sigma_{1}+\sigma_{2}+\sigma_{3}=\bar{\sigma}_{1}+\sigma_{3}
$$

and if we choose

$$
\sigma_{3}=a+\left(\Im \sigma_{1}\right) \mathbf{i}, \quad a \in \mathbb{R},
$$

the arithmetic will be retained real.
Finally, notice that the choice

$$
\sigma_{3}=-\bar{\sigma}_{1}
$$

ensures that the transformation from $L_{1}$ and $U_{1}$ to $L_{4}$ and $U_{4}$ is a restoring shift transformation since $\sigma_{1}+\sigma_{2}+\sigma_{3}=0$.

What we have just showed is summarized in the following lemma.

Lemma 5.1.3 Performing three steps of dqds algorithm with successive shifts $\sigma_{1} \in \mathbb{C}$, $\sigma_{2}=-2\left(\Im \sigma_{1}\right) \mathbf{i}$ and $\sigma_{3}=\bar{\sigma}_{1}$ retains real arithmetic and the shifts are restored. We will have

$$
\begin{gathered}
\mathcal{L U}=\left(U_{1} L_{1}\right)^{2}-2\left(\Re \sigma_{1}\right) U_{1} L_{1}+\left|\sigma_{1}\right|^{2} I \\
L_{4} U_{4}=\mathcal{L}^{-1} U_{1} L_{1} \mathcal{L}
\end{gathered}
$$

where $\mathcal{L}=L_{2} L_{3}$ and $\mathcal{U}=U_{3} U_{2}$.

Now the three dqds steps (5.1), (5.2) and (5.3) become

$$
\begin{align*}
U_{1} L_{1}-\sigma_{1} I & =L_{2} U_{2}  \tag{5.6}\\
U_{2} L_{2}-\left(-2\left(\Im \sigma_{1}\right) \mathbf{i} I\right) & =L_{3} U_{3}  \tag{5.7}\\
U_{3} L_{3}-\left(-\bar{\sigma}_{1} I\right) & =L_{4} U_{4} \tag{5.8}
\end{align*}
$$

and can be rewritten in a more revealing way:

$$
\begin{aligned}
U_{1} L_{1}-\sigma_{1} I & =L_{2} U_{2} \\
\left(U_{2} L_{2}+\sigma_{1} I\right)-\bar{\sigma}_{1} I & =L_{3} U_{3} \\
\left(U_{3} L_{3}+\bar{\sigma}_{1} I\right)-0 I & =L_{4} U_{4}
\end{aligned}
$$

Notice that if we only applied the two dqds steps (5.6) and (5.7), the transformation from $L_{1}$ and $U_{1}$ to $L_{3}$ and $U_{3}$ would not be a restoring shift transformation, $L_{3}$ and $U_{3}$ would be complex as well as the accumulated shift $\sigma_{1}+\sigma_{2}=\sigma_{1}+\left(-2\left(\Im \sigma_{1}\right) \mathbf{i}\right)=\bar{\sigma}_{1}$. The third step (5.8) is therefore needed to come back to real factors $L_{4}$ and $U_{4}$, restoring the shift and getting $L_{4} U_{4}$ similar to $U_{1} L_{1}$.

Another way of analyzing the need of a triple shift is as follows. If we consider only the two steps (5.6) and (5.7) we will have

$$
\begin{gathered}
L_{3} U_{3}=L_{2}^{-1}\left(U_{1} L_{1}\right) L_{2}-\bar{\sigma}_{1} I \\
U_{3} L_{3}=L_{3}^{-1} L_{2}^{-1}\left(U_{1} L_{1}\right) L_{2} L_{3}-\bar{\sigma}_{1} I=\mathcal{L}^{-1}\left(U_{1} L_{1}\right) \mathcal{L}-\bar{\sigma}_{1} I \\
\mathcal{L U}=\left(U_{1} L_{1}\right)^{2}-2\left(\Re \sigma_{1}\right) U_{1} L_{1}+\left|\sigma_{1}\right|^{2} I
\end{gathered}
$$

where $\mathcal{L}=L_{2} L_{3}$ and $\mathcal{U}=U_{3} U_{2}$. Thus $L_{1}, U_{1}, \mathcal{L}$ and $\mathcal{U}$ are real and we can write

$$
\mathcal{L}^{-1}\left(U_{1} L_{1}\right) \mathcal{L}=U_{3} L_{3}+\bar{\sigma}_{1} I
$$

which means that $U_{3} L_{3}+\bar{\sigma}_{1}$ is also real. So, we must have to do the third step (5.8) to get the factors $L_{4}$ and $U_{4}$ of the LU factorization

$$
U_{3} L_{3}+\bar{\sigma}_{1}=L_{4} U_{4}
$$

The goal now will be to avoid complex factors $L_{2}, U_{2}, L_{3}$ and $U_{3}$ and go straight from $L_{1}$ and $U_{1}$ to $L_{4}$ and $U_{4}$, performing the three steps implicitly - implicit triple dqds algorithm.

### 5.2 Connection to the implicit double shifted LR algorithm

This section aims to relate the double shifted LR algorithm to the triple shift dqds algorithm and use the implicit version of the first to derive the implicit version of the second.

If we performed the products $U_{i} L_{i}, i=1,2,3$, and added back the shifts $\sigma=\sigma_{1}$ and $\bar{\sigma}=\bar{\sigma}_{1}$, the three dqds steps (5.6), (5.7) and (5.8) would correspond to a basic iteration (zero shift) followed by a double shifted iteration (shifts $\sigma$ and $\bar{\sigma}$ ) of the LR algorithm with an extra LU factorization in the end. In more detail,

> LR algorithm dqds
$\left\{\begin{array}{l}J_{1}=L_{1} U_{1} \\ J_{2}=U_{1} L_{1}\end{array}\right.$
$\left\{\begin{array}{rlrl}J_{2}-\sigma I & =L_{2} U_{2} & U_{1} L_{1}-\sigma I=L_{2} U_{2} \\ J_{3} & =U_{2} L_{2}+\sigma I\end{array}\right.$
$\left\{\begin{array}{rl}J_{3}-\bar{\sigma} I & =L_{3} U_{3} \\ J_{4} & =U_{3} L_{3}+\bar{\sigma} I\end{array} \quad\left(U_{2} L_{2}+\sigma I\right)-\bar{\sigma} I=L_{3} U_{3}\right.$
$\left\{\begin{array}{l}J_{4}=L_{4} U_{4} \\ \cdots\end{array} \quad\left(U_{3} L_{3}+\bar{\sigma} I\right)-0 I=L_{4} U_{4}\right.$

Figure 5.1 below also exhibits this relation - LR algorithm goes with matrices $J_{i}$ and triple dqds goes with the factors $L_{i}$ and $U_{i}$ - and shows clearly the need of three dqds steps in order to go from real $L_{1}$ and $U_{1}$ to real $L_{4}$ and $U_{4}$.


Figure 5.1: Double shift LR and three steps of dqds

Recall from lemma 3.4.1 (page 65) that after a double shifted LR iteration applied to $J_{2}$ with shifts $\sigma$ and $\bar{\sigma}$ we will have

$$
\mathbf{L U}=J_{2}^{2}-2(\Re \sigma) J_{2}+|\sigma|^{2} I \equiv M
$$

and

$$
J_{4}=\mathbf{L}^{-1} J_{2} \mathbf{L}
$$

where $\mathbf{L}=L_{2} L_{3}$ is unit lower triangular and $\mathbf{U}=U_{3} U_{2}$ is upper triangular. This is exactly what lemma 5.1.3 tells for the triple dqds algorithm but focusing on the factors of the LU decompositions $J_{2}=U_{1} L_{1}$ and $J_{4}=L_{4} U_{4}$,

$$
\begin{gathered}
\mathcal{L U}=\left(U_{1} L_{1}\right)^{2}-2(\Re \sigma) U_{1} L_{1}+|\sigma|^{2} I \\
L_{4} U_{4}=\mathcal{L}^{-1} U_{1} L_{1} \mathcal{U}
\end{gathered}
$$

with $\mathcal{L}=\mathbf{L}$ and $\mathcal{U}=\mathbf{U}$.
To obtain matrix $J_{4}$, the implicit double shifted LR algorithm uses the technique of bulge chasing described in section 3.4.1 (page 68) going directly from tridiagonal matrix $J_{2}$ to $J_{4}$. This is justified by the implicit L theorem that says that matrix $\mathbf{L}$ is uniquely determined by its first column which we can compute because it is proportional to the first column of $M$.

Using the same notation of example 3.4.1, $\mathbf{L}$ will be given as a product of elementary or

Gauss matrices $\mathbf{L}_{i}, i=1, \ldots, n-1$,

$$
\mathbf{L}=\mathbf{L}_{1} \mathbf{L}_{2} \cdots \mathbf{L}_{n-1}
$$

where $\mathbf{L}_{i}=I+\mathbf{l}_{i}^{\prime} e_{i}^{T}$ with $\mathbf{l}_{i}^{\prime}=\left[\begin{array}{llllllll}0 & \ldots & 0 & * & * & 0 & \ldots & 0\end{array}\right]^{T}$ (the non-zero elements $*$ appear in entries $i+1$ and $i+2$ ) and $J_{4}$ will be the result of successive similarity transformations given by $\mathbf{L}_{i}$,

$$
\begin{aligned}
J_{4} & =\mathbf{L}_{n-1}^{-1} \cdots \mathbf{L}_{1}^{-1} J_{2} \mathbf{L}_{1} \cdots \mathbf{L}_{n-1} \\
& =\left(\mathbf{L}_{1} \cdots \mathbf{L}_{n-1}\right)^{-1} J_{2} \mathbf{L}_{1} \cdots \mathbf{L}_{n-1} .
\end{aligned}
$$

So, $\mathbf{L}_{1}$ is computed and the first similarity transformation $\mathbf{L}_{1}^{-1} J_{2} \mathbf{L}_{1}$ spoils the tridiagonal form of $J_{2}$. Then the bulge is pushed down and to the right to restore the tridiagonal form using the transformations given by $\mathbf{L}_{i}, i=2, \ldots, n-1$. In the end we have $J_{4}$.

Analogously, to obtain $L_{4}$ and $U_{4}$, the implicit triple dqds algorithm will construct $\mathcal{L}$ as a product of elementary matrices $\mathcal{L}, i=1, \ldots, n-1$, obtained through a similar process of bulge chasing, such that the factors $L_{4}$ and $U_{4}$ will result from

$$
L_{4} U_{4}=\left(\mathcal{L}_{n-1}^{-1} \ldots \mathcal{L}_{2}^{-1} \mathcal{L}_{1}^{-1} U_{1}\right)\left(L_{1} \mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{n-1}\right)
$$

But things will be done in such a way that the products $U_{1} L_{1}$ and $L_{4} U_{4}$ are not computed explicitly. We will get $L_{4}$ and $U_{4}$ from $L_{1}$ and $U_{1}$ without computing any products of the form $L U$ or $U L$ explicitly.

We can now ask: "Do we have $\mathbf{L}_{i}=\mathcal{L}_{i}, i=1, \ldots, n-1$ ?" Implicit L theorem guarantees that $\mathbf{L}=\mathcal{L}$, because the first column of both matrices is the same and there is no more choice. And the expression for the product (see section 2.1, page 25)

$$
\mathbf{L}_{1} \mathbf{L}_{2} \cdots \mathbf{L}_{n-1}=\mathbf{L}=\mathcal{L}=\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{n-1}
$$

ensures that

$$
\mathbf{L}_{i}=\mathcal{L}_{i}, \quad i=1, \ldots, n-1
$$

In summary, starting with the factors $L_{1}$ and $U_{1}$ and the shift $\sigma$, if we

- normalize $1^{\text {st }}$ column of $\left(U_{1} L_{1}\right)^{2}-2\left(\Re \sigma_{1}\right) U_{1} L_{1}+\left|\sigma_{1}\right|^{2} I$ (equal to first column of $\mathcal{L}$ ) and compute $\mathcal{L}_{1}$
- spoil the bidiagonal form with $\underbrace{\mathcal{L}_{1}^{-1} U_{1}} \underbrace{L_{1} \mathcal{L}_{1}}$ and
- apply bulge chasing to get $L^{\prime}$ and $U^{\prime}$

$$
L^{\prime} U^{\prime}=\underbrace{\mathcal{L}_{n-1}^{-1} \ldots \mathcal{L}_{2}^{-1} \mathcal{L}_{1}^{-1} U_{1}} \underbrace{L_{1} \mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{n-1}},
$$

we will obtain implicitly the same result of performing the three dqds steps (5.6), (5.7) and (5.8), that is, we will get $L^{\prime}=L_{4}$ and $U^{\prime}=U_{4}$.

We could have started thinking only in performing the two dqds steps (5.6) and (5.7), that is a double dqds algorithm. If we decided to use bulge chasing, we would't get $L_{3}$ and $U_{3}$ but $L_{4}$ and $U_{4}$, ending up realizing that we were effectively performing a triple dqds iteration.

### 5.3 Derivation of triple dqds

The role of the implicit triple dqds algorithm will be to construct matrices $\mathcal{L}_{i}, i=1, \ldots, n$, and the unique matrix $X$ such that

$$
\begin{aligned}
& L_{4}=\mathcal{L}_{n}^{-1} \mathcal{L}_{n-1}^{-1} \ldots \mathcal{L}_{2}^{-1} \mathcal{L}_{1}^{-1} U_{1} X^{-1} \\
& U_{4}=X L_{1} \mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{n-1} \mathcal{L}_{n}
\end{aligned}
$$

(we will have $\mathcal{L}_{n}=I$ ). Matrix $X$ will be best written as

$$
X=X_{n} X_{2} \cdots X_{1}
$$

and each $X_{i}, i=1, \ldots, n$, will be the product of two matrices

$$
X_{i} \equiv Y_{i} Z_{i} .
$$

Conceptually, for each step of the implicit triple dqds algorithm we will start with $F:=U_{1}$ and $G:=L_{1}$ and, at each minor step $i, i=1, \ldots, n$, matrices $Z_{i}, \mathcal{L}_{i}$ and $Y_{i}$ are chosen to chase the bulges. After $n$ minor steps, $F$ will be transformed into $L_{4}$ and $G$ into $U_{4}$. So, we will be transforming an upper bidiagonal $U_{1}$ into a lower bidiagonal $L_{4}$ and a lower bidiagonal $L_{1}$ into an upper bidiagonal $U_{4}$. Details will be given in the next section.

Consider matrices $L_{1}$ and $U_{1}$ as follows

$$
L_{1}=\left[\begin{array}{cccccc}
1 & & & & &  \tag{5.9}\\
l_{1} & 1 & & & & \\
& l_{2} & 1 & & & \\
& & \ddots & \ddots & & \\
& & & l_{n-2} & 1 & \\
& & & & l_{n-1} & 1
\end{array}\right] \text { and } U_{1}=\left[\begin{array}{cccccc}
u_{1} & 1 & & & & \\
& u_{2} & 1 & & & \\
& & u_{3} & 1 & & \\
& & & \ddots & \ddots & \\
& & & & u_{n-1} & 1 \\
& & & & & u_{n}
\end{array}\right] .
$$

The product $U_{1} L_{1}$ is

$$
U_{1} L_{1}=\left[\begin{array}{cccccc}
u_{1}+l_{1} & 1 & & & & \\
u_{2} l_{1} & u_{2}+l_{2} & 1 & & & \\
& u_{3} l_{2} & u_{3}+l_{3} & 1 & & \\
& & \ddots & \ddots & \ddots & \\
& & & u_{n-1} l_{n-2} & u_{n-1}+l_{n-1} & 1 \\
& & & & u_{n} l_{n-1} & u_{n}
\end{array}\right]
$$

and the first column of $\left(U_{1} L_{1}\right)^{2}-2(\Re \sigma) U_{1} L_{1}+|\sigma|^{2} I$ is

$$
\left[\begin{array}{c}
\left(u_{1}+l_{1}\right)^{2}+u_{2} l_{1}-2(\Re \sigma)\left(u_{1}+l_{1}\right)+|\sigma|^{2}  \tag{5.10}\\
u_{2} l_{1}\left(u_{1}+l_{1}+u_{2}+l_{2}-2(\Re \sigma)\right) \\
u_{2} l_{1} u_{3} l_{2} \\
0 \\
\vdots \\
0
\end{array}\right] .
$$

We will start by presenting a $6 \times 6$ example to illustrate the general pattern of matrices $Z_{i}, \mathcal{L}_{i}$ and $Y_{i}$. If we work through the example in detail, we will understand the algorithm.

## Example 5.3.1

For a $6 \times 6$ example consider

$$
L_{1}=\left[\begin{array}{cccccc}
1 & & & & & \\
l_{1} & 1 & & & & \\
& l_{2} & 1 & & & \\
& & l_{3} & 1 & & \\
& & & l_{4} & 1 & \\
& & & & l_{5} & 1
\end{array}\right] \quad \text { and } \quad U_{1}=\left[\begin{array}{cccccc}
u_{1} & 1 & & & & \\
& u_{2} & 1 & & & \\
& & u_{3} & 1 & & \\
& & & u_{4} & 1 & \\
& & & & u_{5} & 1 \\
& & & & & u_{6}
\end{array}\right] .
$$

We seek

$$
L_{4}:=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & \\
& \hat{l}_{2} & 1 & & & \\
& & \hat{l}_{3} & 1 & & \\
& & & \hat{l}_{4} & 1 & \\
& & & & \hat{l}_{5} & 1
\end{array}\right] \quad \text { and } \quad U_{4}:=\left[\begin{array}{ccccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& & \hat{u}_{3} & 1 & & \\
& & & \hat{u}_{4} & 1 & \\
& & & & \hat{u}_{5} & 1 \\
& & & & & & \hat{u}_{6}
\end{array}\right] .
$$

Initially, we define

$$
F:=U_{1} \quad \text { and } \quad G:=L_{1}
$$

and

$$
F^{(1)}:=F \quad \text { and } \quad G^{(1)}:=G \text {. }
$$

## Step 1

1.a) Define $Z_{1}^{-1}$ and obtain

$$
\underbrace{F Z_{1}^{-1}} \underbrace{Z_{1} G}=\underbrace{U_{1} Z_{1}^{-1}} \underbrace{Z_{1} L_{1}} .
$$

Matrix $Z_{1}^{-1}$ is chosen to zero out entry $(1,2)$ of $F$ and put 1 into $(1,1)$ entry.
Observe that

$$
\left[\begin{array}{cc}
u_{1} & 1 \\
0 & u_{2}
\end{array}\right]=\left[\begin{array}{ll}
1 & 0 \\
0 & u_{2}
\end{array}\right]\left[\begin{array}{ll}
u_{1} & 1 \\
0 & 1
\end{array}\right]
$$

Then we have

Thus

$$
F Z_{1}^{-1}=\left[\begin{array}{cccccc}
1 & 0 & & & & \\
& u_{2} & 1 & & & \\
& & u_{3} & 1 & & \\
& & & u_{4} & 1 & \\
& & & & u_{5} & 1 \\
& & & & & \\
& & & & & u_{6}
\end{array}\right] \text { and } Z_{1} G=\left[\begin{array}{ccccc}
l_{1}+u_{1} & 1 & & & \\
l_{1} & 1 & & & \\
& & l_{2} & 1 & \\
& & l_{3} & 1 & \\
& & & l_{4} & 1 \\
& & & & l_{5}
\end{array}\right]
$$

Observe that entries $(1,1)$ of $F$ and $(1,2)$ of $G$ became 1 .
Let

$$
F \longleftarrow F Z_{1}^{-1} \quad \text { and } \quad G \longleftarrow Z_{1} G
$$

1.b) Define elementary matrix $\mathcal{L}_{1}$ to get

$$
\underbrace{\mathcal{L}_{1}^{-1} F} \underbrace{G \mathcal{L}_{1}}=\underbrace{\mathcal{L}_{1}^{-1} U_{1} Z_{1}^{-1}} \underbrace{Z_{1} L_{1} \mathcal{L}_{1}} .
$$

Let $\boldsymbol{l}_{1}=\left[\begin{array}{llllll}1 & * & * & 0 & 0 & 0\end{array}\right]^{T}$ be the first column of $\mathcal{L}_{1}$, obtained by normalizing (5.10) for $n=6$. We will have $\mathcal{L}_{1}=I+\boldsymbol{l}_{1}^{\prime} \boldsymbol{e}_{1}^{T}$, where $\boldsymbol{l}_{1}^{\prime}=\left[\begin{array}{llllll}0 & * & * & 0 & 0 & 0\end{array}\right]^{T}$, and $\mathcal{L}_{1}^{-1}=I-\boldsymbol{l}_{1}^{\prime} \boldsymbol{e}_{1}^{T}$. Thus,
$\mathcal{L}_{1}^{-1} F=\left[\begin{array}{ccccccc}1 & 0 & & & & \\ * & u_{2} & 1 & & & \\ + & & u_{3} & 1 & & & \\ & & & u_{4} & 1 & \\ & & & & u_{5} & 1 \\ & & & & & u_{6}\end{array}\right] \quad$ and $\quad G \mathcal{L}_{1}=\left[\begin{array}{llllll}\hat{u}_{1} & 1 & & & & \\ + & 1 & & & & \\ + & l_{2} & 1 & & & \\ + & & l_{3} & 1 & & \\ & & & l_{4} & 1 & \\ & & & & l_{5} & 1\end{array}\right]$.

Let

$$
F \longleftarrow \mathcal{L}_{1}^{-1} F \quad \text { and } \quad G \longleftarrow G \mathcal{L}_{1}
$$

Note that the preliminary transform $Z_{1}$ ensured that, at this point, column 2 of $F$ is unchanged. The first row of $F$ and first row of $G$ are in final form, but first column of $F$ and first column of $G$ have bulges that are indicated by plus signs. Next, the $3 \times 1$ bulge in $G$ will be chased.
1.c) Define elementary matrix $Y_{1}$ to obtain

$$
\underbrace{F Y_{1}^{-1}} \underbrace{Y_{1} G}=\underbrace{\mathcal{L}_{1}^{-1} U_{1} Z_{1}^{-1} Y_{1}^{-1}} \underbrace{Y_{1} Z_{1} L_{1} \mathcal{L}_{1}} .
$$

Matrix $Y_{1}$ affects only rows 2,3 and 4 of $Y_{1} G$, zeroing out entries $(2,1),(3,1)$ and $(4,1)$ of $G$. So, $Y_{1}^{-1}=I+\boldsymbol{y}_{1}^{\prime} \boldsymbol{e}_{1}^{T}$ with $\boldsymbol{y}_{1}^{\prime}=\left[\begin{array}{llllll}0 & * & * & * & 0 & 0\end{array}\right]^{T}, Y_{1}=I-\boldsymbol{y}_{1}^{\prime} \boldsymbol{e}_{1}^{T}$ and

$$
F Y_{1}^{-1}=\left[\begin{array}{cccccc}
1 & 0 & & & & \\
\hat{l}_{1} & u_{2} & 1 & & & \\
+ & & u_{3} & 1 & & \\
+ & & & u_{4} & 1 & \\
& & & & u_{5} & 1 \\
& & & & & u_{6}
\end{array}\right] \quad \text { and } \quad Y_{1} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& * & & & & \\
& + & 1 & & & \\
& + & l_{3} & 1 & & \\
& & & l_{4} & 1 & \\
& & & & l_{5} & 1
\end{array}\right]
$$

Let

$$
F \longleftarrow F Y_{1}^{-1} \quad \text { and } \quad G \longleftarrow Y_{1} G
$$

and define

$$
\begin{aligned}
F^{(2)} & :=F=\mathcal{L}_{1}^{-1} U_{1} Z_{1}^{-1} Y_{1}^{-1} \\
G^{(2)} & :=G=Y_{1} Z_{1} L_{1} \mathcal{L}_{1}
\end{aligned}
$$

This completes the first step. The first column of $G$ is in final form but the first column of $F$ is not yet in final form - there exists a $2 \times 1$ bulge.

The pattern shown above is carried down the matrix by later transformations $Z_{i}$, $\mathcal{L}_{i}$ and $Y_{i}, i=2, \ldots, 6$, forcing the bulges down by one row and one column at each step. This way in the end $F$ will be transformed into the lower bidiagonal matrix $L_{4}$ and $G$ into the upper bidiagonal matrix $U_{4}$.

## Step 2

2.a) Define $Z_{2}^{-1}$ and get

$$
\underbrace{F Z_{2}^{-1}} \underbrace{Z_{2} G}=\underbrace{F^{(2)} Z_{2}^{-1}} \underbrace{Z_{2} G^{(2)}}
$$

Matrix $Z_{2}^{-1}$ will zero out entry $(2,3)$ of $F$ and place 1 into entry $(2,2)$. Matrix $Z_{2}$ will turn entry $(2,3)$ of $G$ into 1 . We have

$$
Z_{2}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
& \frac{1}{u_{2}} & -\frac{1}{u_{2}} & & & \\
& 0 & 1 & & & \\
& & & 1 & & \\
& & & & 1 & \\
& & & & & \\
& & & & & 1
\end{array}\right], \quad Z_{2}=\left[\begin{array}{llllll}
1 & & & & & \\
& u_{2} & 1 & & & \\
& 0 & 1 & & \\
& & & 1 & \\
& & & & 1 & \\
& & & & & 1
\end{array}\right]
$$

and

$$
F Z_{2}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & 0 & & & \\
+ & & u_{3} & 1 & & \\
+ & & & u_{4} & 1 & \\
& & & & u_{5} & 1 \\
& & & & & u_{6}
\end{array}\right], \quad Z_{2} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & \\
& * & 1 & & & \\
& + & 1 & & & \\
& & l_{3} & 1 & & \\
& & & l_{4} & 1 & \\
& & & & l_{5} & 1
\end{array}\right]
$$

Let

$$
F \longleftarrow F Z_{2}^{-1} \quad \text { and } \quad G \longleftarrow Z_{2} G
$$

2.b) Define elementary matrix $\mathcal{L}_{2}$ to obtain

$$
\underbrace{\mathcal{L}_{2}^{-1} F} \underbrace{G \mathcal{L}_{2}}=\underbrace{\mathcal{L}_{2}^{-1} F^{(2)} Z_{2}^{-1}} \underbrace{Z_{2} G^{(2)} \mathcal{L}_{2}} .
$$

Matrix $\mathcal{L}_{2}$ will affect only rows 3 and 4 of $\mathcal{L}_{2}^{-1} F$, zeroing out entries $(3,1)$ and $(4,1)$ of $F$. So, $\mathcal{L}_{2}=I+l_{2}^{\prime} e_{2}^{T}$, where $\boldsymbol{l}_{2}^{\prime}=\left[\begin{array}{llllll}0 & 0 & * & * & 0 & 0\end{array}\right]^{T}$, and $\mathcal{L}_{2}^{-1}=I-l_{2}^{\prime} \boldsymbol{e}_{2}^{T}$. Thus,

$$
\mathcal{L}_{2}^{-1} F=\left[\begin{array}{ccccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & 0 & & & \\
& * & u_{3} & 1 & & \\
& + & & u_{4} & 1 & \\
& & & & u_{5} & 1 \\
& & & & & u_{6}
\end{array}\right] \quad \text { and } \quad G \mathcal{L}_{2}=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& + & 1 & & & \\
& + & l_{3} & 1 & & \\
& + & & l_{4} & 1 & \\
& & & & l_{5} & 1
\end{array}\right]
$$

Let

$$
F \longleftarrow \mathcal{L}_{2}^{-1} F \quad \text { and } \quad G \longleftarrow G \mathcal{L}_{2}
$$

Bulge $2 \times 1$ in $F$ was chased and second row of $F$ and second row of $G$ are now in final form.
2.c) Define matrix $Y_{2}$ to get

$$
\underbrace{F Y_{2}^{-1}} \underbrace{Y_{2} G}=\underbrace{\mathcal{L}_{2}^{-1} F^{(2)} Z_{2}^{-1} Y_{2}^{-1}} \underbrace{Y_{2} Z_{2} G^{(2)} \mathcal{L}_{2}} .
$$

Matrix $Y_{2}$ affects only rows 3,4 and 5 of $Y_{2} G$, zeroing out entries $(3,2),(4,2)$ and $(5,2)$ of $G$. So, $Y_{2}^{-1}=I+\boldsymbol{y}_{2}^{\prime} \boldsymbol{e}_{2}^{T}$ with $\boldsymbol{y}_{2}^{\prime}=\left[\begin{array}{llllll}0 & 0 & * & * & * & 0\end{array}\right]^{T}, Y_{2}=I-\boldsymbol{y}_{2}^{\prime} \boldsymbol{e}_{2}^{T}$ and

$$
F Y_{2}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & 0 & & & \\
& \hat{l}_{2} & u_{3} & 1 & & \\
& + & & u_{4} & 1 & \\
& + & & & u_{5} & 1 \\
& & & & & u_{6}
\end{array}\right] \quad \text { and } \quad Y_{2} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & \\
& \hat{u}_{2} & 1 & & & \\
& & * & & \\
& & & + & 1 & \\
& & & & l_{4} & 1 \\
& & & & l_{5} & 1
\end{array}\right] .
$$

Let

$$
F \longleftarrow F Y_{2}^{-1} \quad \text { and } \quad G \longleftarrow Y_{2} G
$$

and define

$$
\begin{aligned}
F^{(3)} & :=F=\mathcal{L}_{2}^{-1} F^{(2)} Z_{2}^{-1} Y_{2}^{-1} \\
G^{(3)} & :=G=Y_{2} Z_{2} G^{(2)} \mathcal{L}_{2} .
\end{aligned}
$$

The second step is complete. The second column of $G$ is in final form but the second column of $F$ is not yet in final form.

## Step 3

3.a) Define $Z_{3}^{-1}$ and get

$$
\underbrace{F Z_{3}^{-1}} \underbrace{Z_{3} G}=\underbrace{F^{(3)} Z_{3}^{-1}} \underbrace{Z_{3} G^{(3)}} .
$$

Matrix $Z_{3}^{-1}$ will zero out entry $(3,4)$ of $F$ and place 1 into entry $(3,3) . Z_{3}$ will turn into 1 entry $(3,4)$ of $G$. We have

$$
Z_{3}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
& 1 & & & & \\
& & \frac{1}{u_{3}} & -\frac{1}{u_{3}} & & \\
& & 0 & 1 & & \\
& & & & 1 & \\
& & & & & 1
\end{array}\right], \quad Z_{3}=\left[\begin{array}{llllll}
1 & & & & & \\
& 1 & & & & \\
& & u_{3} & 1 & \\
& & 0 & 1 & \\
& & & & 1 & \\
& & & & & 1
\end{array}\right]
$$

and

$$
F Z_{3}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & \\
& \hat{l}_{2} & 1 & 0 & & \\
& + & & u_{4} & 1 & \\
& + & & & u_{5} & 1 \\
& & & & & u_{6}
\end{array}\right], \quad Z_{3} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& & * & 1 & & \\
& & & & 1 & \\
& & & & l_{4} & 1 \\
& & & & l_{5} & 1
\end{array}\right] .
$$

Let

$$
F \longleftarrow F Z_{3}^{-1} \quad \text { and } \quad G \longleftarrow Z_{3} G
$$

3.b) Define elementary matrix $\mathcal{L}_{3}$ to obtain

$$
\underbrace{\mathcal{L}_{3}^{-1} F} \underbrace{G \mathcal{L}_{3}}=\underbrace{\mathcal{L}_{3}^{-1} F^{(3)} Z_{3}^{-1}} \underbrace{Z_{3} G^{(3)} \mathcal{L}_{3}} .
$$

Matrix $\mathcal{L}_{3}$ will affect only rows 4 and 5 of $\mathcal{L}_{3}^{-1} F$, zeroing out entries $(4,2)$ and $(5,2)$ of $F$. So, $\mathcal{L}_{3}=I+l_{3}^{\prime} e_{3}^{T}$, where $l_{3}^{\prime}=\left[\begin{array}{llllll}0 & 0 & 0 & * & * & 0\end{array}\right]^{T}$, and $\mathcal{L}_{3}^{-1}=I-l_{3}^{\prime} \boldsymbol{e}_{3}^{T}$. Thus,
$\mathcal{L}_{3}^{-1} F=\left[\begin{array}{lllllll}1 & & & & & \\ \hat{l}_{1} & 1 & & & & & \\ & \hat{l}_{2} & 1 & 0 & & & \\ & & * & u_{4} & 1 & \\ & & + & & u_{5} & 1 \\ & & & & & u_{6}\end{array}\right] \quad$ and $\quad G \mathcal{L}_{3}=\left[\begin{array}{cccccc}\hat{u}_{1} & 1 & & & & \\ & \hat{u}_{2} & 1 & & & \\ & & \hat{u}_{3} & 1 & & \\ & & & 1 & & \\ & & & & l_{4} & 1 \\ & & & & l_{5} & 1\end{array}\right]$.

Let

$$
F \longleftarrow \mathcal{L}_{3}^{-1} F \quad \text { and } \quad G \longleftarrow G \mathcal{L}_{3} .
$$

Third row of $F$ and third row of $G$ are in final form.
3.c) Define elementary matrix $Y_{3}$ to produce

$$
\underbrace{F Y_{3}^{-1}} \underbrace{Y_{3} G}=\underbrace{\mathcal{L}_{3}^{-1} F^{(3)} Z_{3}^{-1} Y_{3}^{-1}} \underbrace{Y_{3} Z_{3} G^{(3)} \mathcal{L}_{3}} .
$$

Matrix $Y_{3}$ affects only rows 4,5 and 6 of $Y_{3} G$, zeroing out entries $(4,3),(5,3)$ and $(6,3)$ of $G$. So, $Y_{3}^{-1}=I+\boldsymbol{y}_{3}^{\prime} e_{3}^{T}$ with $\boldsymbol{y}_{3}^{\prime}=\left[\begin{array}{llllll}0 & 0 & 0 & * & * & *\end{array}\right]^{T}, Y_{3}=I-\boldsymbol{y}_{3}^{\prime} \boldsymbol{e}_{3}^{T}$ and

$$
F Y_{3}^{-1}=\left[\begin{array}{ccccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & & \\
& \hat{l}_{2} & 1 & 0 & & \\
& & \hat{l}_{3} & u_{4} & 1 & \\
& & + & & u_{5} & 1 \\
& & + & & & u_{6}
\end{array}\right] \quad \text { and } \quad Y_{3} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& & \hat{u}_{3} & 1 & & \\
& & & * & & \\
& & & & + & 1 \\
& & & & & \\
& & & & l_{5} & 1
\end{array}\right] .
$$

Let

$$
F \longleftarrow F Y_{3}^{-1} \quad \text { and } \quad G \longleftarrow Y_{3} G
$$

and define

$$
\begin{aligned}
F^{(4)} & :=F=\mathcal{L}_{3}^{-1} F^{(3)} Z_{3}^{-1} Y_{3}^{-1} \\
G^{(4)} & :=G=Y_{3} Z_{3} G^{(3)} \mathcal{L}_{3} .
\end{aligned}
$$

The third step is complete. The third column of $G$ is in final form but the third column of $F$ is not yet in final form.

## Step 4

4.a) Define $Z_{4}^{-1}$ and get

$$
\underbrace{F Z_{4}^{-1}} \underbrace{Z_{4} G}=\underbrace{F^{(4)} Z_{4}^{-1}} \underbrace{Z_{4} G^{(4)}} .
$$

Matrix $Z_{4}^{-1}$ will zero out entry $(4,5)$ of $F$ and place 1 into entry $(4,4)$. $Z_{4}$ will turn into 1 entry $(4,5)$ of $G$. We have

$$
Z_{4}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
& 1 & & & & \\
& & 1 & & & \\
& & & \frac{1}{u_{4}} & -\frac{1}{u_{4}} & \\
& & & 0 & 1 & \\
& & & & & 1
\end{array}\right], \quad Z_{4}=\left[\begin{array}{cccccc}
1 & & & & & \\
& 1 & & & & \\
& & 1 & & & \\
& & & u_{4} & 1 & \\
& & & 0 & 1 & \\
& & & & & 1
\end{array}\right]
$$

and

$$
F Z_{4}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & \\
& \hat{l}_{2} & 1 & & & \\
& & \hat{l}_{3} & 1 & 0 & \\
& & + & & u_{5} & 1 \\
& & + & & & u_{6}
\end{array}\right], \quad Z_{4} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& & \hat{u}_{3} & 1 & & \\
& & & * & 1 & \\
& & & & + & 1 \\
& & & & & \\
& & & & l_{5} & 1
\end{array}\right] .
$$

Let

$$
F \longleftarrow F Z_{4}^{-1} \quad \text { and } \quad G \longleftarrow Z_{4} G .
$$

4.b) Apply elementary matrices $\mathcal{L}_{4}$ and $\mathcal{L}_{4}^{-1}$,

$$
\underbrace{\mathcal{L}_{4}^{-1} F} \underbrace{G \mathcal{L}_{4}}=\underbrace{\mathcal{L}_{4}^{-1} F^{(4)} Z_{4}^{-1}} \underbrace{Z_{4} G^{(4)} \mathcal{L}_{4}} .
$$

Matrix $\mathcal{L}_{4}$ will affect only rows 5 and 6 of $\mathcal{L}_{4}^{-1} F$, zeroing out entries $(5,3)$ and $(6,3)$ of $F$. So, $\mathcal{L}_{4}=I+\boldsymbol{l}_{4}^{\prime} \boldsymbol{e}_{4}^{T}$, where $\boldsymbol{l}_{4}^{\prime}=\left[\begin{array}{lllll}0 & 0 & 0 & 0 & *\end{array}\right]^{T}$, and $\mathcal{L}_{4}^{-1}=I-l_{4}^{\prime} e_{4}^{T}$. Thus,
$\mathcal{L}_{4}^{-1} F=\left[\begin{array}{cccccc}1 & & & & & \\ \hat{l}_{1} & 1 & & & & \\ & \hat{l}_{2} & 1 & & & \\ & & \hat{l}_{3} & 1 & 0 & \\ & & & * & u_{5} & 1 \\ & & & + & & u_{6}\end{array}\right] \quad$ and $\quad G \mathcal{L}_{4}=\left[\begin{array}{cccccc}\hat{u}_{1} & 1 & & & & \\ & \hat{u}_{2} & 1 & & & \\ & & \hat{u}_{3} & 1 & & \\ & & & & \hat{u}_{4} & 1 \\ & & & & & \\ & & & & 1 & \\ & & & & & l_{5} \\ & & & & \end{array}\right]$.

Let

$$
F \longleftarrow \mathcal{L}_{4}^{-1} F \quad \text { and } \quad G \longleftarrow G \mathcal{L}_{4} .
$$

Fourth row of $F$ and fourth row of $G$ are in final form.
4.c) Apply $Y_{4}$ and $Y_{4}^{-1}$,

$$
\underbrace{F Y_{4}^{-1}} \underbrace{Y_{4} G}=\underbrace{\mathcal{L}_{4}^{-1} F^{(4)} Z_{4}^{-1} Y_{4}^{-1}} \underbrace{Y_{4} Z_{4} G^{(4)} \mathcal{L}_{4}} .
$$

Matrix $Y_{4}$ affects only rows 5 and 6 of $Y_{4} G$, zeroing out entries $(5,4)$ and $(6,4)$ of $G$. So, $Y_{4}^{-1}=I+\boldsymbol{y}_{4}^{\prime} e_{4}^{T}$ with $\boldsymbol{y}_{4}^{\prime}=\left[\begin{array}{llllll}0 & 0 & 0 & 0 & * & *\end{array}\right]^{T}, Y_{4}=I-\boldsymbol{y}_{4}^{\prime} \boldsymbol{e}_{4}^{T}$ and

Let

$$
F \longleftarrow F Y_{4}^{-1} \quad \text { and } \quad G \longleftarrow Y_{4} G
$$

and define

$$
\begin{aligned}
F^{(5)} & :=F=\mathcal{L}_{4}^{-1} F^{(4)} Z_{4}^{-1} Y_{4}^{-1} \\
G^{(5)} & :=G=Y_{4} Z_{4} G^{(4)} \mathcal{L}_{4}
\end{aligned}
$$

The fourth step is complete.

## Step 5

5.a) Define $Z_{5}^{-1}$ and get

$$
\underbrace{F Z_{5}^{-1}} \underbrace{Z_{5} G}=\underbrace{F^{(5)} Z_{5}^{-1}} \underbrace{Z_{5} G^{(5)}} .
$$

Matrix $Z_{5}^{-1}$ will zero out entry $(5,6)$ of $F$ and place 1 into entry $(5,5) . Z_{5}$ will turn into 1 entry $(5,6)$ of $G$. We have

$$
Z_{5}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
& 1 & & & & \\
& & 1 & & & \\
& & & 1 & & \\
& & & & \frac{1}{u_{5}} & -\frac{1}{u_{5}} \\
& & & & 0 & 1
\end{array}\right], \quad Z_{5}=\left[\begin{array}{llllll}
1 & & & & & \\
& 1 & & & & \\
& & 1 & & & \\
& & & 1 & & \\
& & & & u_{5} & 1 \\
& & & & 0 & 1
\end{array}\right]
$$

and

$$
F Z_{5}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & \\
& \hat{l}_{2} & 1 & & & \\
& & \hat{l}_{3} & 1 & & \\
& & & \hat{l}_{4} & 1 & 0 \\
& & & + & & u_{6}
\end{array}\right], \quad Z_{5} G=\left[\begin{array}{ccccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & & \\
& & \hat{u}_{3} & 1 & & \\
& & & \hat{u}_{4} & 1 & \\
& & & & & * & 1 \\
& & & & & & \\
& & & & & & 1
\end{array}\right] .
$$

Let

$$
F \longleftarrow F Z_{5}^{-1} \quad \text { and } \quad G \longleftarrow Z_{5} G
$$

5.b) Apply matrices $\mathcal{L}_{5}$ and $\mathcal{L}_{5}^{-1}$ to obtain

$$
\underbrace{\mathcal{L}_{5}^{-1} F} \underbrace{G \mathcal{L}_{5}}=\underbrace{\mathcal{L}_{5}^{-1} F^{(5)} Z_{5}^{-1}} \underbrace{Z_{5} G^{(5)} \mathcal{L}_{5}}
$$

Matrix $\mathcal{L}_{5}$ will affect only rows 6 of $\mathcal{L}_{5}^{-1} F$, zeroing out entries $(6,4)$ of $F$. So, $\mathcal{L}_{5}=I+\boldsymbol{l}_{5}^{\prime} \boldsymbol{e}_{5}^{T}$, where $\boldsymbol{l}_{5}^{\prime}=\left[\begin{array}{cccccc}0 & 0 & 0 & 0 & 0 & *\end{array}\right]^{T}$, and $\mathcal{L}_{5}^{-1}=I-\boldsymbol{l}_{5}^{\prime} \boldsymbol{e}_{5}^{T}$. Thus,


Let

$$
F \longleftarrow \mathcal{L}_{5}^{-1} F \quad \text { and } \quad G \longleftarrow G \mathcal{L}_{5}
$$

Fifth row of $F$ and fifth row of $G$ are in final form.
5.c) Apply matrices $Y_{5}$ and $Y_{5}^{-1}$ to obtain

$$
\underbrace{F Y_{5}^{-1}} \underbrace{Y_{5} G}=\underbrace{\mathcal{L}_{5}^{-1} F^{(5)} Z_{5}^{-1} Y_{5}^{-1}} \underbrace{Y_{5} Z_{5} G^{(5)} \mathcal{L}_{5}}
$$

Matrix $Y_{5}$ affects only rows 6 of $Y_{5} G$, zeroing out entries $(6,5)$ of $G$. So, $Y_{5}^{-1}=I+\boldsymbol{y}_{5}^{\prime} \boldsymbol{e}_{5}^{T}$ with $\boldsymbol{y}_{5}^{\prime}=\left[\begin{array}{llllll}0 & 0 & 0 & 0 & 0 & *\end{array}\right]^{T}, Y_{5}=I-\boldsymbol{y}_{5}^{\prime} \boldsymbol{e}_{5}^{T}$ and

$$
F Y_{5}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & \\
& \hat{l}_{2} & 1 & & & \\
& & \hat{l}_{3} & 1 & & \\
& & & \hat{l}_{4} & 1 & 0 \\
& & & & \hat{l}_{5} & u_{6}
\end{array}\right] \quad \text { and } \quad Y_{5} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& & \hat{u}_{3} & 1 & & \\
& & & \hat{u}_{4} & 1 & \\
& & & & & \hat{u}_{5} \\
& & & & & \\
& & & & & *
\end{array}\right] .
$$

Let

$$
F \longleftarrow F Y_{5}^{-1} \quad \text { and } \quad G \longleftarrow Y_{5} G
$$

and define

$$
\begin{aligned}
F^{(6)} & :=F=\mathcal{L}_{5}^{-1} F^{(5)} Z_{5}^{-1} Y_{5}^{-1} \\
G^{(6)} & :=G=Y_{5} Z_{5} G^{(5)} \mathcal{L}_{5}
\end{aligned}
$$

The fifth step is complete and the two bidiagonals are found. There is only need to place 1 in entry $(6,6)$ of $F$.

## Step 6

6.a) Apply $Z_{6}$ and $Z_{6}^{-1}$,

$$
\underbrace{F Z_{6}^{-1}} \underbrace{Z_{6} G}=\underbrace{F^{(6)} Z_{6}^{-1}} \underbrace{Z_{6} G^{(6)}} .
$$

Matrix $Z_{6}^{-1}$ place 1 in entry $(6,6)$ of $F$. We have
and

$$
F Z_{6}^{-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & \\
& \hat{l}_{2} & 1 & & & \\
& & \hat{l}_{3} & 1 & & \\
& & & \hat{l}_{4} & 1 & \\
& & & & \hat{l}_{5} & 1
\end{array}\right], \quad Z_{6} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& & \hat{u}_{3} & 1 & & \\
& & & \hat{u}_{4} & 1 & \\
& & & & \hat{u}_{5} & 1 \\
& & & & & \\
& & & & & *
\end{array}\right] .
$$

Let

$$
F \longleftarrow F Z_{6}^{-1} \quad \text { and } \quad G \longleftarrow Z_{6} G .
$$

Matrices $F$ and $G$ are already in final form.
6.b) Let $\mathcal{L}_{6}=I$ and

$$
\underbrace{\mathcal{L}_{6}^{-1} F} \underbrace{G \mathcal{L}_{6}}=\underbrace{\mathcal{L}_{6}^{-1} F^{(6)} Z_{6}^{-1}} \underbrace{Z_{6} G^{(6)} \mathcal{L}_{6}} .
$$

Thus,

$$
\mathcal{L}_{6}^{-1} F=\left[\begin{array}{cccccc}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & \\
& \hat{l}_{2} & 1 & & & \\
& & \hat{l}_{3} & 1 & & \\
& & & \hat{l}_{4} & 1 & \\
& & & & \hat{l}_{5} & 1
\end{array}\right] \quad \text { and } \quad G \mathcal{L}_{6}=\left[\begin{array}{ccccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& & \hat{u}_{3} & 1 & & \\
& & & \hat{u}_{4} & 1 & \\
& & & & \hat{u}_{5} & 1 \\
& & & & & \hat{u}_{6}
\end{array}\right] \text {. }
$$

Let

$$
F \longleftarrow \mathcal{L}_{6}^{-1} F \quad \text { and } \quad G \longleftarrow G \mathcal{L}_{6}
$$

6.c) Let $Y_{6}=I$ and finally

$$
\underbrace{F Y_{6}^{-1}} \underbrace{Y_{6} G}=\underbrace{\mathcal{L}_{6}^{-1} F^{(6)} Z_{6}^{-1} Y_{6}^{-1}} \underbrace{Y_{6} Z_{6} G^{(6)} \mathcal{L}_{6}}
$$

and

$$
F \longleftarrow F Y_{6}^{-1} \quad \text { and } \quad G \longleftarrow Y_{6} G
$$

In the end we have

$$
L_{4}=F \text { and } U_{4}=G . \diamond
$$

For general matrices $L_{1}$ and $U_{1}$ as in (5.9), we will have

$$
\begin{aligned}
L_{4} U_{4} & =\underbrace{\mathcal{L}_{n}^{-1} \cdots \mathcal{L}_{1}^{-1} U_{1} Z_{1}^{-1} Y_{1}^{-1} \cdots Z_{n}^{-1} Y_{n}^{-1}} \underbrace{Y_{n} Z_{n} \cdots Y_{1} Z_{1} L_{1} \mathcal{L}_{1} \cdots \mathcal{L}_{n}} \\
& =\underbrace{\mathcal{L}_{n}^{-1} \cdots \mathcal{L}_{1}^{-1} U_{1} X_{1}^{-1} \cdots X_{n}^{-1}} \underbrace{X_{n} \cdots X_{1} L_{1} \mathcal{L}_{1} \cdots \mathcal{L}_{n}} \\
& =\underbrace{\mathcal{L}^{-1} U_{1} X^{-1}} \underbrace{X L_{1} \mathcal{L}}
\end{aligned}
$$

where

$$
\mathcal{L} \equiv \mathcal{L}_{1} \cdots \mathcal{L}_{n} \quad \text { and } \quad X \equiv X_{n} \cdots X_{1}
$$

with $X_{i}=Y_{i} Z_{i}, i=1, \ldots, n$.

And for a complex shift $\sigma$, the triple dqds algorithm will be called tridqds and is given by:
tridqds( $\sigma$ ) :

$$
\begin{aligned}
& F=U_{1} ; G=L_{1} \\
& F=F Z_{1}^{-1} ; G=Z_{1} G \\
& F=\mathcal{L}_{1}^{-1} F ; G=G \mathcal{L}_{1} \quad\left[\text { form } \mathcal{L}_{1} \operatorname{using}(5.10)\right] \\
& F=F Y_{1}^{-1} ; G=Y_{1} G \\
& \text { for } i=2, \ldots, n-3 \\
& F=F Z_{i}^{-1} ; G=Z_{i} G \\
& F=\mathcal{L}_{i}^{-1} F ; G=G \mathcal{L}_{i} \\
& F=F Y_{i}^{-1} ; G=Y_{i} G \quad\left[Z_{i} \text { with one, } \mathcal{L}_{i} \text { with two and } Y_{i}\right. \text { with three } \\
& \text { end for } \\
& \text { off-diagonal entries] } \\
& \text { \% step n-2 } \\
& F=F Z_{n-2}^{-1} ; G=Z_{n-2} G \\
& F=\mathcal{L}_{n-2}^{-1} F ; G=G \mathcal{L}_{n-2} \\
& F=F Y_{n-2}^{-1} ; G=Y_{n-2} G \quad\left[Y_{n-2}\right. \text { with two off-diagonal entries] } \\
& \text { \% step n-1 } \\
& F=F Z_{n-1}^{-1} ; G=Z_{n-1} G \\
& F=\mathcal{L}_{n-1}^{-1} F ; G=G \mathcal{L}_{n-1} \\
& F=F Y_{n-1}^{-1} ; G=Y_{n-1} G \quad\left[Y_{n-1} \text { and } \mathcal{L}_{n-1} \text { with one off-diagonal entry }\right] \\
& \text { \% step n } \\
& \mathcal{L}_{n}=I ; Y_{n}=I \\
& F=F Z_{n}^{-1} ; G=Z_{n} G \quad\left[Z_{n}\right. \text { diagonal] } \\
& L_{4}=F ; F_{4}=G
\end{aligned}
$$

### 5.3.1 Details of tridqds

In this section we will go into the details of the tridqds algorithm described in the previous section. Consider matrices $L_{1}$ and $U_{1}$ as defined in (5.9) and let us say that matrices $L_{4}$ and $U_{4}$ have the same shape with subdiagonal entries $\hat{l}_{1}, \ldots, \hat{l}_{n-1}$ and diagonal entries $\hat{u}_{1}, \ldots, \hat{u}_{n}$, respectively.

For each iteration of tridqds, in the beginning of a minor step $i, i=2, \ldots, n-2$, the active parts of $F$ and $G$ are

$$
F=\left[\begin{array}{ccccccc}
\ddots & & & & &  \tag{5.11}\\
\ddots & 1 & & & & \\
& \hat{l}_{i-1} & u_{i} & 1 & & & \\
& + & u_{i+1} & 1 & \\
& + & & & u_{i+2} & \ddots \\
& & & & & \ddots
\end{array}\right] \text { and } G=\left[\begin{array}{ccccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i-1} & 1 & & & \\
& & * & & & \\
& & & & 1 & & \\
& & & + & l_{i+1} & 1 & \\
& & & & \ddots & \ddots
\end{array}\right] .
$$

So, $F$ starts to be $U_{1}$ and is transformed into $L_{4}$ and $G$ is given initially as $L_{1}$ and is transformed into $U_{4}$.

For $i=2, \ldots, n-3$, a general minor step $i$ consists in what follows.

## Typical minor step i

i.a) Define $Z_{i}$ and let

$$
F \longleftarrow F Z_{i}^{-1} \quad \text { and } \quad G \longleftarrow Z_{i} G .
$$

Matrix $Z_{i}^{-1}$ will zero out entry $(i, i+1)$ of $F$ and place 1 into entry $(i, i)$. $Z_{i}$ turns into 1 entry $(i, i+1)$ of $G$.

$$
Z_{i}^{-1}=\left[\begin{array}{cccccc}
\ddots & & & & & \\
& 1 & & & & \\
& & \frac{1}{u_{i}} & -\frac{1}{u_{i}} & & \\
& & 0 & 1 & & \\
& & & & 1 & \\
& & & & & \ddots
\end{array}\right], \quad Z_{i}=\left[\begin{array}{llllll}
\ddots & & & & & \\
& 1 & & & & \\
& & u_{i} & 1 & & \\
& & 0 & 1 & & \\
& & & & & 1
\end{array}\right]
$$

and

$$
F Z_{i}^{-1}=\left[\begin{array}{ccccccc}
\ddots & & & & & & \\
\ddots & 1 & & & & & \\
& \hat{l}_{i-1} & 1 & 0 & & & \\
& + & u_{i+1} & 1 & & \\
& + & & & u_{i+2} & \ddots & \\
& & & & & \ddots
\end{array}\right], \quad Z_{i} G=\left[\begin{array}{ccccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i-1} & 1 & & & \\
& & * & 1 & & \\
& & & & & & \\
& & & & & & \\
& & l_{i+1} & 1 & \\
& & & & & \ddots & \ddots
\end{array}\right] .
$$

i.b) Define elementary matrix $\mathcal{L}_{i}$ and let

$$
F \longleftarrow \mathcal{L}_{i}^{-1} F \quad \text { and } \quad G \longleftarrow G \mathcal{L}_{i} .
$$

$\mathcal{L}_{i}$ will affect only rows $i+1$ and $i+2$ of $\mathcal{L}_{i}^{-1} F$, zeroing out entries $(i+1, i-1)$ and $(i+2, i-1)$ of $F$;
$\mathcal{L}_{i}=I+\boldsymbol{l}_{i}^{\prime} \boldsymbol{e}_{i}^{T}$, where $\boldsymbol{l}_{i}^{\prime}=\left[\begin{array}{lllll}0 \ldots 0 & * & * & 0 \ldots 0\end{array}\right]^{T}, \mathcal{L}_{i}^{-1}=I-\boldsymbol{l}_{i}^{\prime} \boldsymbol{e}_{i}^{T}$ and

$$
\mathcal{L}_{i}^{-1} F=\left[\begin{array}{ccccccc}
\ddots & & & & & & \\
\ddots & 1 & & & & & \\
& \hat{l}_{i-1} & 1 & & & & \\
& & * & u_{i+1} & 1 & \\
& & & & u_{i+2} & 1 \\
& & & & & \ddots
\end{array}\right], \quad G \mathcal{L}_{i}=\left[\begin{array}{ccccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i} & 1 & & & \\
& + & 1 & & & \\
& + & l_{i+1} & 1 & & \\
& + & & l_{i+2} & 1 & \\
& & & & & \ddots & \ddots
\end{array}\right] .
$$

i.c) Define elementary matrix $Y_{i}$ and let

$$
F \longleftarrow F Y_{i}^{-1} \quad \text { and } \quad G \longleftarrow Y_{i} G .
$$

$Y_{i}$ affects only rows $i+1, i+2$ and $i+3$ of $Y_{i} G$, zeroing out entries $(i+1, i),(i+2, i)$ and $(i+3, i)$ of $G$;
$Y_{i}^{-1}=I+\boldsymbol{y}_{i}^{\prime} \boldsymbol{e}_{i}^{T}$ with $\boldsymbol{y}_{i}^{\prime}=\left[\begin{array}{llllll}0 \ldots 0 & * & * & * & 0 \ldots 0\end{array}\right]^{T}, Y_{i}=I-\boldsymbol{y}_{i}^{\prime} \boldsymbol{e}_{i}^{T}$ and

Compare with (5.11) to see that we are back to the same pattern from which we began: the bulges have been chased one column to the right and one row to the bottom. Turning to a practical implementation, next we will see how the calculations involved in this process can be organized.

For each minor step $i, i=2, \ldots, n-3$, consider $F$ and $G$ as in (5.11). Denote the $2 \times 1$ bulge in $F$, indicated with plus signs, by $\left[\begin{array}{ll}x_{l} & y_{l}\end{array}\right]^{T}$. And denote the entries $(i, i),(i+1, i)$ and $(i+2, i)$ in $G$, indicated with $*,+,+$, by $\left[\begin{array}{lll}x_{r} & y_{r} & z_{r}\end{array}\right]^{T}$. Subscripts $l$ and $r$ derive from "left" and "right", respectively. This way we have

$$
F=\left[\begin{array}{cccccc}
\ddots & & & & &  \tag{5.12}\\
\ddots & 1 & & & & \\
& \hat{l}_{i-1} & u_{i} & 1 & & \\
& x_{l} & & u_{i+1} & 1 & \\
& y_{l} & & & u_{i+2} & \ddots \\
& & & & & \ddots
\end{array}\right] \text { and } G=\left[\begin{array}{ccccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i-1} & 1 & & & \\
& & x_{r} & & & \\
& & y_{r} & 1 & & \\
& & z_{r} & l_{i+1} & 1 & \\
& & & & \ddots & \ddots
\end{array}\right]
$$

and the typical minor step $i$ can be accomplished using only these auxiliary variables. The details will be shown in the following pages.

## Typical minor step i

i.a) - The inverse $Z_{i}^{-1}$ and $Z_{i}$

$$
Z_{i}^{-1}=\left[\begin{array}{cccccc}
\ddots & & & & & \\
& 1 & & & & \\
& & \frac{1}{u_{i}} & -\frac{1}{u_{i}} & & \\
& & 0 & 1 & & \\
& & & & 1 & \\
& & & & & \ddots
\end{array}\right], \quad Z_{i}^{-1}=\left[\begin{array}{cccccc}
\ddots & & & & & \\
& 1 & & & & \\
& & u_{i} & 1 & & \\
& & 0 & 1 & & \\
& & & & 1 & \\
& & & & & \ddots
\end{array}\right]
$$

- The effect of $Z_{i}^{-1}$

$$
F Z_{i}^{-1}=\left[\begin{array}{cccccc}
\ddots & & & & & \\
\ddots & 1 & & & & \\
& \hat{l}_{i-1} & 1 & 0 & & \\
& x_{l} & & u_{i+1} & 1 & \\
& y_{l} & & & u_{i+2} & \ddots \\
& & & & & \ddots
\end{array}\right]
$$

- The effect of $Z_{i}$

$$
Z_{i} G=\left[\begin{array}{cccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i-1} & 1 & & & \\
& & x_{r} & 1 & & \\
& & y_{r} & 1 & & \\
& & z_{r} & l_{i+1} & 1 & \\
& & & & \ddots & \ddots
\end{array}\right]
$$

where

$$
x_{r} \longleftarrow x_{r} * u_{i}+y_{r}
$$

i.b) - The inverse $\mathcal{L}_{i}^{-1}$

$$
\mathcal{L}_{i}^{-1}=\left[\begin{array}{llllll}
\ddots & & & & & \\
& \ddots & & & & \\
& & 1 & & & \\
& & x_{l} & 1 & & \\
& & y_{l} & & 1 & \\
& & & & & \ddots
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{l} \longleftarrow-x_{l} / \hat{l}_{i-1} \\
& y_{l} \longleftarrow-y_{l} / \hat{l}_{i-1}
\end{aligned}
$$

- The effect of $\mathcal{L}_{i}^{-1}$

$$
\mathcal{L}_{i}^{-1} F=\left[\begin{array}{cccccc}
\ddots & & & & & \\
\ddots & 1 & & & & \\
& \hat{l}_{i-1} & 1 & & & \\
& & x_{l} & u_{i+1} & 1 & \\
& & y_{l} & & u_{i+2} & \ddots \\
& & & & & \ddots
\end{array}\right]
$$

- The effect of $\mathcal{L}_{i}$

$$
G \mathcal{L}_{i}=\left[\begin{array}{cccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i} & 1 & & & \\
& x_{r} & 1 & & & \\
& y_{r} & l_{i+1} & 1 & & \\
& z_{r} & & l_{i+2} & 1 & \\
& & & & \ddots & \ddots
\end{array}\right]
$$

where

$$
\begin{aligned}
& \hat{u}_{i} \longleftarrow x_{r}-x_{l} \\
& x_{r} \longleftarrow y_{r}-x_{l} \\
& y_{r} \longleftarrow z_{r}-y_{l}-x_{l} * l_{i+1} \\
& z_{r} \longleftarrow-y_{l} * l_{i+2}
\end{aligned}
$$

i.c) - The inverse $Y_{i}^{-1}$

$$
Y_{i}^{-1}=\left[\begin{array}{llllll}
\ddots & & & & & \\
& 1 & & & & \\
& x_{r} & 1 & & & \\
& y_{r} & & 1 & & \\
& z_{r} & & & 1 & \\
& & & & & \ddots
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{r} \longleftarrow x_{r} / \hat{u}_{i} \\
& y_{r} \longleftarrow y_{r} / \hat{u}_{i} \\
& z_{r} \longleftarrow z_{r} / \hat{u}_{i}
\end{aligned}
$$

- The effect of $Y_{i}^{-1}$

$$
F Y_{i}^{-1}=\left[\begin{array}{cccccc}
\ddots & & & & & \\
\ddots & 1 & & & & \\
& \hat{l}_{i} & u_{i+1} & 1 & & \\
& x_{l} & & u_{i+2} & 1 & \\
& y_{l} & & & u_{i+3} & \ddots \\
& & & & & \ddots
\end{array}\right]
$$

where

$$
\begin{aligned}
& \hat{l}_{i} \longleftarrow x_{l}+y_{r}+x_{r} * u_{i+1} \\
& x_{l} \longleftarrow y_{l}+z_{r}+y_{r} * u_{i+2} \\
& y_{l} \longleftarrow z_{r} * u_{i+3}
\end{aligned}
$$

- The effect of $Y_{i}$

$$
Y_{i} G=\left[\begin{array}{llllll}
\ddots & \ddots & & & & \\
& \hat{u}_{i} & 1 & & & \\
& & x_{r} & & & \\
& & y_{r} & 1 & & \\
& & z_{r} & l_{i+2} & 1 & \\
& & & & \ddots & \ddots
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{r} \longleftarrow 1-x_{r} \\
& y_{r} \longleftarrow l_{i+1}-y_{r} \\
& z_{r} \longleftarrow-z_{r}
\end{aligned}
$$

Now the details of the last three minor steps.

## Minor step n-2

[n-2].a) - The inverse $Z_{n-2}^{-1}$ and $Z_{n-2}$

$$
Z_{n-2}^{-1}=\left[\begin{array}{ccccc}
\ddots & & & & \\
& 1 & & & \\
& & \frac{1}{u_{n-2}} & -\frac{1}{u_{n-2}} & \\
& & 0 & 1 & \\
& & & & 1
\end{array}\right], \quad Z_{n-2}=\left[\begin{array}{ccccc}
\ddots & & & & \\
& 1 & & & \\
& & u_{n-2} & 1 & \\
& & 0 & 1 & \\
& & & & 1
\end{array}\right]
$$

- The effect of $Z_{n-2}^{-1}$

$$
F Z_{n-2}^{-1}=\left[\begin{array}{ccccc}
\ddots & & & & \\
\ddots & 1 & & & \\
& \hat{l}_{n-3} & 1 & 0 & \\
& x_{l} & & u_{n-1} & 1 \\
& y_{l} & & & u_{n}
\end{array}\right]
$$

- The effect of $Z_{n-2}$

$$
Z_{n-2} G=\left[\begin{array}{ccccc}
\ddots & \ddots & & & \\
& \hat{u}_{n-3} & 1 & & \\
& & x_{r} & 1 & \\
& & y_{r} & 1 & \\
& & z_{r} & l_{n-1} & 1
\end{array}\right]
$$

where

$$
x_{r} \longleftarrow x_{r} * u_{n-2}+y_{r}
$$

[n-2].b) - The inverse $\mathcal{L}_{n-2}^{-1}$

$$
\mathcal{L}_{n-2}^{-1}=\left[\begin{array}{lllll}
\ddots & & & & \\
& 1 & & & \\
& & 1 & & \\
& & x_{l} & 1 & \\
& & y_{l} & & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{l} \longleftarrow-x_{l} / \hat{l}_{n-3} \\
& y_{l} \longleftarrow-y_{l} / \hat{l}_{n-3}
\end{aligned}
$$

- The effect of $\mathcal{L}_{n-2}^{-1}$

$$
\mathcal{L}_{n-2}^{-1} F=\left[\begin{array}{ccccc}
\ddots & & & & \\
\ddots & 1 & & & \\
& \hat{l}_{n-3} & 1 & & \\
& & x_{l} & u_{n-1} & 1 \\
& & y_{l} & & u_{n}
\end{array}\right]
$$

- The effect of $\mathcal{L}_{n-2}$

$$
G \mathcal{L}_{n-2}=\left[\begin{array}{ccccc}
\ddots & \ddots & & & \\
& \hat{u}_{n-3} & 1 & & \\
& & \hat{u}_{n-2} & 1 & \\
& & x_{r} & 1 & \\
& & y_{r} & l_{n-1} & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
\hat{u}_{n-2} & \longleftarrow x_{r}-x_{l} \\
x_{r} & \longleftarrow y_{r}-x_{l} \\
y_{r} & \longleftarrow z_{r}-y_{l}-x_{l} * l_{n-1}
\end{aligned}
$$

[n-2].c) • The inverse $Y_{n-2}^{-1}$

$$
Y_{n-2}^{-1}=\left[\begin{array}{lllll}
\ddots & & & & \\
& 1 & & & \\
& & 1 & & \\
& & x_{r} & 1 & \\
& & y_{r} & & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{r} \longleftarrow x_{r} / \hat{u}_{n-2} \\
& y_{r} \longleftarrow y_{r} / \hat{u}_{n-2}
\end{aligned}
$$

- The effect of $Y_{n-2}^{-1}$

$$
F Y_{n-2}^{-1}=\left[\begin{array}{ccccc}
\ddots & \ddots & & & \\
& \hat{l}_{n-3} & 1 & & \\
& & \hat{l}_{n-2} & u_{n-1} & 1 \\
& & x_{l} & & u_{n}
\end{array}\right]
$$

where

$$
\begin{gathered}
\hat{l}_{n-2} \longleftarrow x_{l}+y_{r}+x_{r} * u_{n-1} \\
x_{l} \longleftarrow y_{l}+y_{r} * u_{n}
\end{gathered}
$$

- The effect of $Y_{n-2}$

$$
Y_{n-2} G=\left[\begin{array}{cccccc}
\ddots & \ddots & & & \\
& \hat{u}_{n-3} & 1 & & \\
& & \hat{u}_{n-2} & 1 & \\
& & & x_{r} & \\
& & & y_{r} & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{r} \longleftarrow 1-x_{r} \\
& y_{r} \longleftarrow l_{n-1}-y_{r}
\end{aligned}
$$

## Minor step n-1

[n-1].a) - The inverse $Z_{n-1}^{-1}$ and $Z_{n-1}$

$$
Z_{n-1}^{-1}=\left[\begin{array}{cccc}
\ddots & & & \\
& 1 & & \\
& & \frac{1}{u_{n-1}} & -\frac{1}{u_{n-1}} \\
& 0 & 1
\end{array}\right], \quad Z_{n-1}=\left[\begin{array}{llll}
\ddots & & & \\
& 1 & & \\
& & u_{n-1} & 1 \\
& & 0 & 1
\end{array}\right]
$$

- The effect of $Z_{n-1}^{-1}$

$$
F Z_{n-1}^{-1}=\left[\begin{array}{cccc}
\ddots & & & \\
\ddots & 1 & & \\
& \hat{l}_{n-2} & 1 & 0 \\
& x_{l} & & u_{n}
\end{array}\right]
$$

- The effect of $Z_{n-1}$

$$
Z_{n-1} G=\left[\begin{array}{cccc}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & x_{r} & 1 \\
& & y_{r} & 1
\end{array}\right]
$$

where

$$
x_{r} \longleftarrow x_{r} * u_{n-1}+y_{r}
$$

[n-1].b) - The inverse $\mathcal{L}_{n-1}^{-1}$

$$
\mathcal{L}_{n-1}^{-1}=\left[\begin{array}{llll}
\ddots & & & \\
& 1 & & \\
& & 1 & \\
& & x_{l} & 1
\end{array}\right]
$$

where

$$
x_{l} \longleftarrow-x_{l} / \hat{l}_{n-2}
$$

- The effect of $\mathcal{L}_{n-1}^{-1}$

$$
\mathcal{L}_{n-1}^{-1} F=\left[\begin{array}{cccc}
\ddots & & & \\
\ddots & 1 & & \\
& \hat{l}_{n-2} & 1 & \\
& & x_{l} & u_{n}
\end{array}\right]
$$

- The effect of $\mathcal{L}_{n-1}$

$$
G \mathcal{L}_{n-1}=\left[\begin{array}{cccc}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & \hat{u}_{n-1} & 1 \\
& & x_{r} & 1
\end{array}\right]
$$

where

$$
\begin{gathered}
\hat{u}_{n-1} \longleftarrow x_{r}-x_{l} \\
x_{r} \longleftarrow y_{r}-x_{l}
\end{gathered}
$$

[n-1].c) - The inverse $Y_{n-1}^{-1}$

$$
Y_{n-1}^{-1}=\left[\begin{array}{llll}
\ddots & & & \\
& 1 & & \\
& & 1 & \\
& & x_{r} & 1
\end{array}\right]
$$

where

$$
x_{r} \longleftarrow x_{r} / \hat{u}_{n-1}
$$

- The effect of $Y_{n-1}^{-1}$

$$
F Y_{n-1}^{-1}=\left[\begin{array}{cccc}
\ddots & & & \\
\ddots & 1 & & \\
& \hat{l}_{n-2} & 1 & \\
& & \hat{l}_{n-1} & u_{n}
\end{array}\right]
$$

where

$$
\hat{l}_{n-1} \longleftarrow x_{l}+x_{r} * u_{n}
$$

- The effect of $Y_{n-1}$

$$
Y_{n-1} G=\left[\begin{array}{cccc}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & \hat{u}_{n-1} & 1 \\
& & & x_{r}
\end{array}\right]
$$

where

$$
x_{r} \longleftarrow 1-x_{r}
$$

## $\underline{\text { Minor step n }}$

n.a) - The inverse $Z_{n}^{-1}$ and $Z_{n}$

$$
Z_{n}^{-1}=\left[\begin{array}{llll}
\ddots & & & \\
& \ddots & & \\
& & 1 & \\
& & & \frac{1}{u_{n}}
\end{array}\right], \quad Z_{n}=\left[\begin{array}{llll}
\ddots & & & \\
& \ddots & & \\
& & 1 & \\
& & & u_{n}
\end{array}\right]
$$

- The effect of $Z_{n}^{-1}$

$$
F Z_{n}^{-1}=\left[\begin{array}{ccccc}
\ddots & & & \\
\ddots & 1 & & \\
& \hat{l}_{n-2} & 1 & \\
& & \hat{l}_{n-1} & 1
\end{array}\right]
$$

- The effect of $Z_{n}$

$$
Z_{n} G=\left[\begin{array}{cccc}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & \hat{u}_{n-1} & 1 \\
& & & x_{r}
\end{array}\right]
$$

where

$$
x_{r} \longleftarrow x_{r} * u_{n}
$$

n.b) - The inverse $\mathcal{L}_{n}^{-1}=I$

- The effect of $\mathcal{L}_{n}$

$$
G \mathcal{L}_{n}=\left[\begin{array}{cccc}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & \hat{u}_{n-1} & 1 \\
& & & \hat{u}_{n}
\end{array}\right]
$$

where

$$
\hat{u}_{n} \longleftarrow x_{r}
$$

n.c) - The inverse $Y_{n}^{-1}=I$

Now the last detail that is missing is how to initialize auxiliary variables $x_{l}, y_{l}$ and $x_{r}$, $y_{r}$ and $z_{r}$, that is, how to organize the first step of the tridqds algorithm.

## Minor step 1

Initially we let

$$
F=\left[\begin{array}{ccccc}
u_{1} & 1 & & & \\
& u_{2} & 1 & & \\
& & u_{3} & \ddots & \\
& & & \ddots & 1 \\
& & & & u_{n}
\end{array}\right], \quad G=\left[\begin{array}{ccccc}
x_{r} & & & & \\
y_{r} & 1 & & & \\
z_{r} & l_{2} & 1 & & \\
& & \ddots & \ddots & \\
& & & l_{n-1} & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{r} \longleftarrow 1 \\
& y_{r} \longleftarrow l_{1} \\
& z_{r} \longleftarrow 0
\end{aligned}
$$

1.a) - The inverse $Z_{1}^{-1}$ and $Z_{1}$

$$
Z_{1}^{-1}=\left[\begin{array}{ccccc}
\frac{1}{u_{1}} & -\frac{1}{u_{1}} & & & \\
0 & 1 & & & \\
& & 1 & & \\
& & & \ddots & \\
& & & & 1
\end{array}\right], \quad Z_{1}=\left[\begin{array}{ccccc}
u_{1} & 1 & & & \\
0 & 1 & & \\
& & 1 & \\
& & & \ddots & \\
& & & & \\
& & & & 1
\end{array}\right]
$$

- The effect of $Z_{1}^{-1}$

$$
F Z_{1}^{-1}=\left[\begin{array}{ccccc}
1 & 0 & & & \\
& u_{2} & 1 & & \\
& & u_{3} & \ddots & \\
& & & \ddots & 1 \\
& & & & u_{n}
\end{array}\right]
$$

- The effect of $Z_{1}$

$$
Z_{1} G=\left[\begin{array}{ccccc}
x_{r} & 1 & & & \\
y_{r} & 1 & & & \\
z_{r} & l_{2} & 1 & & \\
& & \ddots & \ddots & \\
& & & l_{n-1} & 1
\end{array}\right]
$$

where

$$
x_{r} \longleftarrow x_{r} * u_{1}+y_{r}
$$

1.b) - The inverse $\mathcal{L}_{1}^{-1}$

$$
\mathcal{L}_{1}^{-1}=\left[\begin{array}{ccccc}
1 & & & & \\
x_{l} & 1 & & & \\
y_{l} & & 1 & & \\
& & & \ddots & \\
& & & & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{l} \longleftarrow\left(u_{1}+l_{1}\right)^{2}+u_{2} l_{1}-2(\Re \sigma)\left(u_{1}+l_{1}\right)+|\sigma|^{2} \\
& y_{l} \longleftarrow-u_{2} l_{1} u_{3} l_{2} / x_{l} \\
& x_{l} \longleftarrow-u_{2} l_{1}\left(u_{1}+l_{1}+u_{2}+l_{2}-2(\Re \sigma)\right) / x_{l}
\end{aligned}
$$

- The effect of $\mathcal{L}_{1}^{-1}$

$$
\mathcal{L}_{1}^{-1} F=\left[\begin{array}{ccccc}
1 & 0 & & & \\
x_{l} & u_{2} & 1 & & \\
y_{l} & & u_{3} & \ddots & \\
& & & \ddots & 1 \\
& & & & u_{n}
\end{array}\right]
$$

- The effect of $\mathcal{L}_{1}$

$$
G \mathcal{L}_{1}=\left[\begin{array}{ccccc}
\hat{u}_{1} & 1 & & & \\
x_{r} & 1 & & & \\
y_{r} & l_{2} & 1 & & \\
z_{r} & & \ddots & \ddots & \\
& & & l_{n-1} & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
& \hat{u}_{1} \longleftarrow x_{r}-x_{l} \\
& x_{r} \longleftarrow y_{r}-x_{l} \\
& y_{r} \longleftarrow z_{r}-y_{l}-x_{l} * l_{2} \\
& z_{r} \longleftarrow-y_{l} * l_{3}
\end{aligned}
$$

1.c) - The inverse $Y_{1}^{-1}$

$$
Y_{1}^{-1}=\left[\begin{array}{ccccc}
1 & & & & \\
x_{r} & 1 & & & \\
y_{r} & & 1 & & \\
z_{r} & & & \ddots & \\
& & & & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{r} \longleftarrow x_{r} / \hat{u}_{1} \\
& y_{r} \longleftarrow y_{r} / \hat{u}_{1} \\
& z_{r} \longleftarrow z_{r} / \hat{u}_{1}
\end{aligned}
$$

- The effect of $Y_{1}^{-1}$

$$
F Y_{1}^{-1}=\left[\begin{array}{ccccc}
1 & 0 & & & \\
\hat{l}_{1} & u_{2} & 1 & & \\
x_{l} & & u_{3} & \ddots & \\
y_{l} & & & \ddots & 1 \\
& & & & u_{n}
\end{array}\right]
$$

where

$$
\begin{aligned}
& \hat{l}_{1} \longleftarrow x_{l}+y_{r}+x_{r} * u_{2} \\
& x_{l} \longleftarrow y_{l}+z_{r}+y_{r} * u_{3} \\
& y_{l} \longleftarrow z_{r} * u_{4}
\end{aligned}
$$

- The effect of $Y_{1}$

$$
Y_{1} G=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& x_{r} & & & & \\
& y_{r} & 1 & & & \\
& z_{r} & l_{3} & 1 & & \\
& & & \ddots & \ddots & \\
& & & & l_{n-1} & 1
\end{array}\right]
$$

where

$$
\begin{aligned}
& x_{r} \longleftarrow 1-x_{r} \\
& y_{r} \longleftarrow l_{2}-y_{r} \\
& z_{r} \longleftarrow-z_{r}
\end{aligned}
$$

It can be verified that minor step 1 could be incorporated in the inner loop, $i=2, \ldots, n-3$, but that would demand an extra auxiliary variable and it is not worthwhile. The final version of tridqds algorithm is presented in the following pages (comments are included to make it easier to relate to the matrix formulation in page 140).

## tridqds( $\sigma$ ) :

$$
\begin{aligned}
& x_{r}=1 ; y_{r}=l_{1} ; z_{r}=0 \\
& \% \text { the effect of } Z_{1} \\
& x_{r}=x_{r} * u_{1}+y_{r} \\
& \% \text { the inverse } \mathcal{L}_{1}^{-1} \\
& x_{l}=\left(u_{1}+l_{1}\right)^{2}+u_{2} l_{1}-2(\Re \sigma)\left(u_{1}+l_{1}\right)+|\sigma|^{2} \\
& y_{l}=-u_{2} l_{1} u_{3} l_{2} / x_{l} \\
& x_{l}=-u_{2} l_{1}\left(u_{1}+l_{1}+u_{2}+l_{2}-2(\Re \sigma)\right) / x_{l} \\
& \% \text { the effect of } \mathcal{L}_{1} \\
& \hat{u}_{1}=x_{r}-x_{l} ; \\
& x_{r}=y_{r}-x_{l} ; y_{r}=z_{r}-y_{l}-x_{l} * l_{2} ; \quad z_{r}=-y_{l} * l_{3} \\
& \% \text { the inverse } Y_{1}^{-1} \\
& x_{r}=x_{r} / \hat{u}_{1} ; y_{r}=y_{r} / \hat{u}_{1} ; z_{r}=z_{r} / \hat{u}_{1} \\
& \% \text { the effect of } Y_{1}^{-1} \\
& \hat{l}_{1}=x_{l}+y_{r}+x_{r} * u_{2} \\
& x_{l}=y_{l}+z_{r}+y_{r} * u_{3} ; y_{l}=z_{r} * u_{4} \\
& \% \text { the effect of } Y_{1} \\
& x_{r}=1-x_{r} ; y_{r}=l_{2}-y_{r} ; z_{r}=-z_{r} \\
& \text { for } i=2, \ldots, n-3 \\
& \% \text { the effect of } Z_{i} \\
& x_{r}=x_{r} * u_{i}+y_{r} \\
& \% \text { the inverse } \mathcal{L}_{i}^{-1} \\
& x_{l}=-x_{l} / \hat{l}_{i-1} ; y_{l}=-y_{l} / \hat{l_{i-1}} ;
\end{aligned}
$$

$\%$ the effect of $\mathcal{L}_{i}$

$$
\begin{aligned}
& \hat{u}_{i}=x_{r}-x_{l} \\
& x_{r}=y_{r}-x_{l} ; \quad y_{r}=z_{r}-y_{l}-x_{l} * l_{i+1} ; \quad z_{r}=-y_{l} * l_{i+2}
\end{aligned}
$$

$\%$ the inverse $Y_{i}^{-1}$
$x_{r}=x_{r} / \hat{u}_{i} ; \quad y_{r}=y_{r} / \hat{u}_{i} ; \quad z_{r}=z_{r} / \hat{u}_{i}$
$\%$ the effect of $Y_{i}^{-1}$
$\hat{l}_{i}=x_{l}+y_{r}+x_{r} * u_{i+1}$
$x_{l}=y_{l}+z_{r}+y_{r} * u_{i+2} ; \quad y_{l}=z_{r} * u_{i+3}$
$\%$ the effect of $Y_{i}$

$$
x_{r}=1-x_{r} ; \quad y_{r}=l_{i+1}-y_{r} ; \quad z_{r}=-z_{r}
$$

end for
\% step n-2
\% the effect of $Z_{n-2}$
$x_{r}=x_{r} * u_{n-2}+y_{r}$
$\%$ the inverse $\mathcal{L}_{n-2}^{-1}$
$x_{l}=-x_{l} / \hat{l}_{n-3} ; \quad y_{l}=-y_{l} / \hat{l}_{n-3} ;$
$\%$ the effect of $\mathcal{L}_{n-2}$
$\hat{u}_{n-2}=x_{r}-x_{l} ;$
$x_{r}=y_{r}-x_{l} ; \quad y_{r}=z_{r}-y_{l}-x_{l} * l_{n-1}$
$\%$ the inverse $Y_{n-2}^{-1}$
$x_{r}=x_{r} / \hat{u}_{n-2} ; \quad y_{r}=y_{r} / \hat{u}_{n-2}$
$\%$ the effect of $Y_{n-2}^{-1}$
$\hat{l}_{n-2}=x_{l}+y_{r}+x_{r} * u_{n-1}$
$x_{l}=y_{l}+y_{r} * u_{n}$
$\%$ the effect of $Y_{n-2}$
$x_{r}=1-x_{r} ; \quad y_{r}=l_{n-1}-y_{r}$

$$
\begin{aligned}
& \% \text { step n-1 } \\
& \% \text { the effect of } Z_{n-1} \\
& x_{r}=x_{r} * u_{n-1}+y_{r} \\
& \% \text { the inverse } \mathcal{L}_{n-1}^{-1} \\
& x_{l}=-x_{l} / \hat{l}_{n-2} \\
& \% \text { the effect of } \mathcal{L}_{n-1} \\
& \hat{u}_{n-1}=x_{r}-x_{l} ; \\
& x_{r}=y_{r}-x_{l} \\
& \% \text { the inverse } Y_{n-1}^{-1} \\
& x_{r}=x_{r} / \hat{u}_{n-1} \\
& \% \text { the effect of } Y_{n-1}^{-1} \\
& \hat{l}_{n-1}=x_{l}+x_{r} * u_{n} \\
& \% \text { the effect of } Y_{n-1} \\
& x_{r}=1-x_{r} \\
& \% \text { step } \mathrm{n} \\
& \% \text { the effect of } Z_{n} \\
& x_{r}=x_{r} * u_{n} \\
& \% \text { the inverse } \mathcal{L}_{n}^{-1}=I \\
& \% \text { the effect of } \mathcal{L}_{n} \\
& \hat{u}_{n}=x_{r} ; \\
& \% \text { the inverse } Y_{n}^{-1}=I
\end{aligned}
$$

### 5.3.2 Operation count for tridqds

In this section we will see how three steps of simple dqds algorithm compares with one step of tridqds in what respects to the number of floating point operations required. In order to make easier this comparison we first remember dqds algorithm from section 3.5.2 (page 75).

$$
\begin{array}{cl}
\operatorname{dqds}(\sigma): & d_{1}=u_{1}-\sigma \\
& \text { for } i=1, \ldots, n-1 \\
\hat{u}_{i}=d_{i}+l_{i} \\
\hat{l}_{i}=l_{i}\left(u_{i+1} / \hat{u}_{i}\right)  \tag{5.13}\\
& d_{i+1}=d_{i}\left(u_{i+1} / \hat{u}_{i}\right)-\sigma
\end{array}
$$

end for

$$
\hat{u}_{n}=d_{n} .
$$

In practice, each $d_{i+1}$ may be written over its predecessor in a single variable $d$ and, if the common subexpression $u_{i+1} / \hat{u}_{i}$ is recognized, then only one division is needed if we use an auxiliary variable.

Table 5.1 below shows that the operation count of one step of tridqds is comparable to three steps of dqds (table expresses only the number of floating point operations in the inner loops).

|  | tridqds | 3 dqds steps |
| :---: | :---: | :---: |
| Divisions | 5 | 3 |
| Multiplications | 6 | 6 |
| Additions | 5 | 3 |
| Subtractions | 6 | 3 |
| Assignments | 16 | 12 |
| Auxiliary variables | 5 | 2 |

Table 5.1: Operation count of tridqds and 3 dqds steps

But to make three steps of dqds equivalent to tridqds we have to consider dqds in complex arithmetic and the total cost is raised by a factor of about 4 . Thus, in complex arithmetic, three steps of dqds are much more expensive than one step of tridqds.

## Chapter 6

## New version of triple dqds

In this chapter we will first describe the connection of a generalized Gram-Schmidt process with the LU factorization and explain how dqds may be derived from this relation independently of qd. Then, we show that quantities $d_{i}$ in dqds provide useful information about the diagonal of $(U L)^{-1}$ and can be incorporated in a shift strategy. For the case of complex eigenvalues we try to generalize these results and exhibit similar ones for tridqds.

A new version of the triple dqds algorithm follows from several results that we establish for the quantities involved in tridqds. This version will be called 3dqds and is more elegant and more efficient than tridqds. Numerical results of an implementation of 3dqds will be shown in the next chapter.

### 6.1 Gram-Schmidt factors

Finding the $L D U$ factorization of any product $B C$ is equivalent to applying a generalized Gram-Schmidt process to the rows of $B$ and to the columns of $C$ so that $B=L P^{*}, C=Q U$, and $P^{*} Q$ is diagonal. See Parlett [40]. When this Gram-Schmidt process is applied to $D U L$ in an efficient manner one obtains the variant dqd of qd algorithm that, as we have discussed earlier, requires a little more arithmetic effort than qd itself. So, in this section we show that dqd could have been discovered independently of qd.

The Gram-Schmidt process is the standard way of producing an orthonormal set of vectors $\left\{\boldsymbol{q}_{1}, \boldsymbol{q}_{2}, \ldots, \boldsymbol{q}_{k}\right\}$ from a linearly independent set $\left\{\boldsymbol{f}_{1}, \boldsymbol{f}_{2}, \ldots, \boldsymbol{f}_{k}\right\}$. The defining property is that $\operatorname{span}\left(\boldsymbol{q}_{1}, \boldsymbol{q}_{2}, \ldots, \boldsymbol{q}_{j}\right)=\operatorname{span}\left(\boldsymbol{f}_{1}, \boldsymbol{f}_{2}, \ldots, \boldsymbol{f}_{j}\right)$, for each $j=1, \ldots, k$. The matrix formulation of this process is the QR factorization

$$
F=Q R
$$

where $F=\left[\begin{array}{llll}\boldsymbol{f}_{1} & \boldsymbol{f}_{2} & \ldots & \boldsymbol{f}_{k}\end{array}\right], Q=\left[\begin{array}{llll}\boldsymbol{q}_{1} & \boldsymbol{q}_{2} & \ldots & \boldsymbol{q}_{k}\end{array}\right]$ and $R$ is a $k \times k$ upper triangular matrix.

The generalization of this process to a pair of vector sets $\left\{\boldsymbol{f}_{1}, \boldsymbol{f}_{2}, \ldots, \boldsymbol{f}_{k}\right\}$ and $\left\{\boldsymbol{g}_{1}, \boldsymbol{g}_{2}, \ldots, \boldsymbol{g}_{k}\right\}$ is so natural that there can be little objection to keeping the name Gram-Schmidt. The context determines immediately whether one or two sets of vectors are involved. Recall that $F^{*}$ denotes the conjugate transpose of $F$.

Theorem 6.1.1 Let $F$ and $G$ be complex $n \times k$ matrices, $n \geq k$, such that $G^{*} F$ permits triangular factorization

$$
G^{*} F=\tilde{L} \tilde{D} \tilde{R},
$$

where $\tilde{L}$ and $\tilde{R}$ are unit triangular, lower and upper, respectively, and $\tilde{D}$ is diagonal. Then there exist unique $n \times k$ matrices $\tilde{Q}$ and $\tilde{P}$ such that

$$
F=\tilde{Q} \tilde{R}, \quad G=\tilde{P} \tilde{L}^{*}, \quad \tilde{P}^{*} \tilde{Q}=\tilde{D} .
$$

This result can be proved by construction.

We will use $X_{i j}$ to denote the $(i, j)$ element of matrix $X$.
Note 6.1.1 When $G$ and $F$ are real and $G=F$, the traditional QR factorization is recovered with an unconventional normalization. Suppose that $F^{T} F$ permits an LU factorization. We must have

$$
F^{T} F=\tilde{L} \tilde{D} \tilde{L}^{T}
$$

and then, from

$$
F=\tilde{Q} \tilde{L}^{T}=\tilde{P} \tilde{L}^{T},
$$

we get $\tilde{Q}=\tilde{P}$, since $\tilde{L}^{T}$ is invertible. Thus

$$
\tilde{P}^{T} \tilde{Q}=\tilde{Q}^{T} \tilde{Q}=\tilde{D}
$$

Now,

$$
\begin{aligned}
\tilde{Q}^{T} \tilde{Q}=\tilde{D} & \Leftrightarrow\left(\tilde{D}^{-\frac{1}{2}}\right)^{T} \tilde{Q}^{T} \tilde{Q} \tilde{D}^{-\frac{1}{2}}=\left(\tilde{D}^{-\frac{1}{2}}\right)^{T} \tilde{D} \tilde{D}^{-\frac{1}{2}} \\
& \Leftrightarrow\left[\tilde{Q} \tilde{D}^{-\frac{1}{2}}\right]^{T} \tilde{Q} \tilde{D}^{-\frac{1}{2}}=\left(\tilde{D}^{-\frac{1}{2}}\right)^{T} \tilde{D}^{\frac{1}{2}} \\
& \Leftrightarrow Q^{T} Q=I
\end{aligned}
$$

where $Q \equiv \tilde{Q} \tilde{D}^{-\frac{1}{2}}$. We have

$$
F=\tilde{Q} \tilde{R}=\tilde{Q} \tilde{D}^{-\frac{1}{2}} \tilde{D}^{\frac{1}{2}} \tilde{R}=Q R
$$

$R \equiv \tilde{D}^{\frac{1}{2}} \tilde{R}=\tilde{D}^{\frac{1}{2}} \tilde{L}^{T}$.
Observe that we have assumed that $\tilde{D}$ is invertible.

Note 6.1.2 In practice, when $n=k$ and $\tilde{D}$ is invertible one can omit $\tilde{Q}$ and write

$$
F=\left(\tilde{P}^{*}\right)^{-1}(\tilde{D} \tilde{R}), \quad G=\tilde{P} \tilde{L}^{*}
$$

and still call it the Gram-Schmidt factorization. The important feature is the uniqueness of $\tilde{Q}$ and $\tilde{P}$.

The Gram-Schmidt factorization leads directly to the differential qd algorithms. Let us show how.

### 6.2 Derivation of dqds from Gram-Schmidt

Consider

$$
\begin{align*}
& L=\left[\begin{array}{llllll}
1 & & & & & \\
l_{1} & 1 & & & & \\
& l_{2} & 1 & & & \\
& & \ddots & \ddots & & \\
& & & l_{n-2} & 1 & \\
& & & & l_{n-1} & 1
\end{array}\right], \quad U=\left[\begin{array}{ccccccc}
u_{1} & 1 & & & & \\
& u_{2} & 1 & & & \\
& & u_{3} & 1 & & \\
& & & \ddots & \ddots & \\
& & & & u_{n-1} & 1 \\
& & & & & u_{n}
\end{array}\right],  \tag{6.1}\\
& \hat{L}=\left[\begin{array}{llllll}
1 & & & & & \\
\hat{l}_{1} & 1 & & & & \\
& \hat{l}_{2} & 1 & & \\
& & \ddots & \ddots & & \\
& & & \hat{l}_{n-2} & 1 & \\
& & & & \hat{l}_{n-1} & 1
\end{array}\right], \quad \hat{U}=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
& \hat{u}_{2} & 1 & & & \\
& & \hat{u}_{3} & 1 & & \\
& & & \ddots & \ddots & \\
& & & & \hat{u}_{n-1} & 1 \\
& & & & & \hat{u}_{n}
\end{array}\right] \tag{6.2}
\end{align*}
$$

and dqds algorithm shown in (5.13), page 161, with $\sigma=0$, that is, dqd algorithm.

Corollary 6.2.1 Let $L$ and $U$ be $n \times n$ bidiagonal matrices as given above. Suppose $U L$ permits factorization

$$
\begin{equation*}
U L=\hat{L} \hat{D} \hat{R}=\hat{L} \hat{U}, \tag{6.3}
\end{equation*}
$$

where $\hat{L}$ and $\hat{R}$ are unit bidiagonal, lower and upper, respectively, and $\hat{D}$ is diagonal. Then there exist unique matrices $\tilde{P}$ and $\tilde{Q}$ such that

$$
U=\hat{L} \tilde{P}^{T}, \quad L=\tilde{Q} \hat{R}, \quad \tilde{P}^{T} \tilde{Q}=\hat{D}
$$

Proof. In previous theorem just take $F=L$ and $G^{*}=U$ and use the fact that $U$ and $L$ are real.

In words, apply Gram-Schmidt to the columns of $L$ and to the rows of $U$ (or columns of $U^{T}$ ), in natural order, to obtain $\hat{L}$ and $\hat{R}$. Then note that $\hat{U}=\hat{D} \hat{R}$. In fact, after Gram-Schmidt, we have

$$
U L=\hat{L} \tilde{P}^{T} \tilde{Q} \hat{R}=\hat{L} \hat{D} \hat{R}=\hat{L} \hat{U} .
$$

We certainly have $\tilde{P}^{T} \tilde{Q}=\hat{D}$ because $L$ and $U$ are invertible and the two QR factorizations are unique.

Also notice that if $\hat{U}_{i i}=0, i<n$, then $U L$ does not permit triangular factorization. However, the theorem allows $\hat{U}_{n n}=0$. When $\hat{U}_{n n} \neq 0$ then $\hat{U}$ is invertible and so is $\hat{D}$. In this case matrices $\tilde{Q}$ and $\tilde{P}$ are also invertible and, since $\tilde{Q}=\tilde{P}^{-T} \hat{D}$ and $\tilde{Q}^{-1} L=\hat{D}^{-1} \tilde{P}^{T} L=\hat{R}$, we can rewrite the factorization as

$$
U L=U \tilde{P}^{-T} \tilde{P}^{T} L=\left(U \tilde{P}^{-T}\right)\left[\hat{D}\left(\hat{D}^{-1} \tilde{P}^{T} L\right)\right]=\hat{L} \hat{D} \hat{R}=\hat{L} \hat{U} .
$$

If we define $K:=\tilde{P}^{T}$ we can write

$$
\begin{equation*}
\underbrace{U K^{-1}} \underbrace{K L}=\hat{L} \hat{U} . \tag{6.4}
\end{equation*}
$$

The matrix $K$ is hidden when we just write $U L=\hat{L} \hat{U}$. However the identification of $K$ with the Gram-Schmidt process goes only half way in the derivation of the dqd algorithm. The nature of Gram-Schmidt process shows that $\tilde{P}$ and $\tilde{Q}$ are upper Hessenberg matrices. We are going to show that they may be written as the product of $n$ simple matrices that are non-orthogonal analogues of plane rotations. That means that $L$ may be changed into $\hat{U}$ and $U$ into $\hat{L}$ by a sequence of simple transformations and neither $K, \tilde{Q}$ nor $\tilde{P}$ need appear explicitly.

A plane transformer in plane $(i, j), i \neq j$, is an identity matrix except for the entries $(i, i),(i, j),(j, i)$ and $(j, j)$. The $2 \times 2$ submatrix they define must be invertible.

Let us describe the first minor step in mapping $U \longrightarrow \hat{L}, L \longrightarrow \hat{U}$. We seek an invertible matrix

$$
X=\left[\begin{array}{cccccc}
x & z & & & & \\
-y & w & & & & \\
& & 1 & & & \\
& & & 1 & & \\
& & & & \ddots & \\
& & & & & 1
\end{array}\right]
$$

such that

$$
U X^{-1}=\left[\begin{array}{cccccc}
1 & 0 & & & & \\
\hat{l}_{1} & * & 1 & & & \\
& & u_{3} & 1 & & \\
& & & \ddots & \ddots & \\
& & & & u_{n-1} & 1 \\
& & & & & u_{n}
\end{array}\right] \text { and } X L=\left[\begin{array}{cccccc}
\hat{u}_{1} & 1 & & & & \\
0 & 1 & & & & \\
& l_{2} & 1 & & & \\
& & \ddots & \ddots & & \\
& & & l_{n-2} & 1 & \\
& & & & l_{n-1} & 1
\end{array}\right]
$$

where * may be anything. Note that

$$
X^{-1}=\left[\begin{array}{ccccc}
\frac{w}{d e t} & -\frac{z}{d e t} & & & \\
\frac{y}{d e t} & \frac{x}{d e t} & & & \\
& & 1 & & \\
& & & \ddots & \\
& & & & 1
\end{array}\right]
$$

where det $=x w+y z$. Thus, we are left with the following two equations

$$
\begin{aligned}
{\left[\begin{array}{cc}
x & z \\
-y & w
\end{array}\right]\left[\begin{array}{ll}
1 & 0 \\
l_{1} & 1
\end{array}\right] } & =\left[\begin{array}{ll}
\hat{u}_{1} & 1 \\
0 & 1
\end{array}\right] \\
{\left[\begin{array}{cc}
u_{1} & 1 \\
0 & u_{2}
\end{array}\right]\left[\begin{array}{cc}
w & -z \\
y & x
\end{array}\right] } & =\left[\begin{array}{ll}
1 & 0 \\
\hat{l}_{1} & *
\end{array}\right] \cdot \operatorname{det}
\end{aligned}
$$

We must have

$$
\begin{align*}
x+z l_{1} & =\hat{u}_{1}  \tag{6.5}\\
z & =1  \tag{6.6}\\
-y+w l_{1} & =0  \tag{6.7}\\
w & =1  \tag{6.8}\\
u_{1} w+y & =\operatorname{det}  \tag{6.9}\\
-u_{1} z+x & =0  \tag{6.10}\\
u_{2} y & =\hat{l}_{1} . \operatorname{det}  \tag{6.11}\\
u_{2} x & =* . d e t . \tag{6.12}
\end{align*}
$$

So, $z=w=1$, equation (6.7) shows that $y=l_{1}$ and equation (6.10) gives $x=u_{1}$. Thus det $=u_{1}+l_{1}$ which verifies (6.9). From (6.5),

$$
\hat{u}_{1}=u_{1}+l_{1}
$$

and from (6.11) we learn that

$$
u_{2} y=u_{2} l_{1}=\hat{l}_{1} \operatorname{det}=\hat{l}_{1} \hat{u}_{1} .
$$

Finally, and of most interest, from (6.12),

$$
*=u_{2} x / \operatorname{det}=u_{2} u_{1} / \operatorname{det}=u_{2} u_{1} / \hat{u}_{1} .
$$

This gives the intermediate quantity $u_{2}$ in dqd and we see it here as something that gets carried to the next minor step.

If we write $d_{1}=u_{1}$ we obtain the start of the inner loop of dqd:

$$
\begin{aligned}
\hat{u}_{1} & =d_{1}+l_{1} \\
\hat{l}_{1} & =l_{1}\left(u_{2} / \hat{u}_{1}\right) \\
d_{2} & =d_{1}\left(u_{2} / \hat{u}_{1}\right) .
\end{aligned}
$$

The typical minor step is similar. Below we show matrices part way through the transformation $U \rightarrow \hat{L}, L \rightarrow \hat{U}$. At minor step $i, i=1, \ldots, n-1$, we have

$$
X_{i}=\left[\begin{array}{llllll}
\ddots & & & & &  \tag{6.13}\\
& 1 & & & & \\
& & d_{i} & 1 & & \\
& & -l_{i} & 1 & & \\
& & & & 1 & \\
& & & & & \ddots
\end{array}\right], \quad X_{i}^{-1}=\left[\begin{array}{llllll}
\ddots & & & & & \\
& 1 & & & \\
& & 1 / \hat{u}_{i} & -1 / \hat{u}_{i} & \\
& & l_{i} / \hat{u}_{i} & d_{i} / \hat{u}_{i} & & \\
& & & & 1 & \\
& & & & & \ddots
\end{array}\right] .
$$

Consider that after completion of step $i$ we let

$$
U \longleftarrow U X_{i}^{-1} \text { and } L \longleftarrow X_{i} L
$$

So, at the end of minor step $i$ we have

$$
U X_{i}^{-1}=\left[\begin{array}{ccccccc}
\ddots & & & & & & \\
\ddots & 1 & 0 & & & & \\
& \hat{l}_{i} & d_{i+1} & 1 & & & \\
& & & u_{i+2} & 1 & \\
& & & & u_{i+3} & \ddots \\
& & & & & \ddots
\end{array}\right], \quad X_{i} L=\left[\begin{array}{ccccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i} & 1 & & & \\
& 0 & 1 & & & \\
& & & l_{i+1} & 1 & & \\
& & & l_{i+2} & 1 & \\
& & & & & \ddots & \ddots
\end{array}\right] .
$$

At the end of minor step $n-1$,

$$
U X_{n-1}^{-1}=\left[\begin{array}{cccc}
\ddots & & & \\
\ddots & 1 & & \\
& \hat{l}_{n-2} & 1 & 0 \\
& & \hat{l}_{n-1} & d_{n}
\end{array}\right] \text { and } X_{n-1} L=\left[\begin{array}{ccccc}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & \hat{u}_{n-1} & 1 \\
& & 0 & 1
\end{array}\right] .
$$

If $d_{n} \neq 0$ we simply multiply row $n$ on the right by $d_{n}$, letting $\hat{u}_{n}=d_{n}$, and divide column $n$ on the left by $d_{n}$ leaving 1 in its place. This is achieved by defining

$$
X_{n}=\left[\begin{array}{llll}
\ddots & & & \\
& \ddots & & \\
& & 1 & \\
& & & d_{n}
\end{array}\right]
$$

and finally

$$
\hat{L}:=\left(U \longleftarrow U X_{n}^{-1}\right) \text { and } \hat{U}:=\left(L \longleftarrow X_{n} L\right)
$$

That is,

$$
\begin{aligned}
& \hat{L}=U X_{1}^{-1} \ldots X_{n}^{-1} \\
& \hat{U}=X_{n} \ldots X_{1} L
\end{aligned}
$$

So, recalling matrix $K$ in (6.4) we can write

$$
\underbrace{U K^{-1}} \underbrace{K L}=\hat{L} \hat{U}
$$

where $K \equiv X_{n} \ldots X_{1}$.
If $d_{n}=0$ we can not define $X_{n}$. Observe that

$$
\begin{aligned}
& \left.\left[\begin{array}{lllll}
\ddots & & & \\
\ddots & 1 & & \\
& \hat{l}_{n-2} & 1 & \\
& & \hat{l}_{n-1} & d_{n}
\end{array}\right]\left[\begin{array}{cccc}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & & \hat{u}_{n-1}
\end{array}\right] 1 \text { } \begin{array}{llll} 
& & & 1
\end{array}\right]= \\
& =\left[\begin{array}{ccccc}
\ddots & & & \\
\ddots & 1 & & \\
& \hat{l}_{n-2} & 1 & \\
& & \hat{l}_{n-1} & 1
\end{array}\right]\left[\begin{array}{lllll}
\ddots & & & \\
& 1 & & \\
& & & 1 & 0 \\
& & & d_{n}
\end{array}\right]\left[\begin{array}{lllll}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & & \hat{u}_{n-1} & 1 \\
& & & & 1
\end{array}\right] .
\end{aligned}
$$

So, $\hat{L}$ does not change and we let $\hat{u}_{n}=1$ and $d_{n}=0 . \hat{U}$ is not invertible but the factorization $\hat{L} \hat{D} \hat{R}$ in (6.3) with $\hat{U}=\hat{D} \hat{R}$ exists.

We have derived the dqd algorithm without reference to qd. Of more significance is the fact that the quantities $d_{i}, i=1, \ldots, n$, provide useful information about $U L$, that qd does not reveal, and so dqd facilitates the choice of a shift.

### 6.2.1 The meaning of $d_{i}$

The meaning of $d_{i}$ is revealed by the next theorem.

Theorem 6.2.2 Consider $L$ and $U$ as described in (6.1). If $U$ is invertible and $U L$ permits triangular factorization then the quantities $d_{i}, i=1, \ldots, n$, generated by the dqd algorithm applied to $L$ and $U$ satisfy

$$
d_{i}^{-1}=\left[(U L)^{-1}\right]_{i i}, \quad i=1, \ldots, n .
$$

Note that if $U$ is invertible and $U L=\hat{L} \hat{U}$ then $\hat{U}$ must also be invertible. This means that $u_{i}$ and $\hat{u}_{i}$ are all non-zero.

Proof. The algorithm may be considered as transforming $L$ to $\hat{U}$ by premultiplications and $U$ to $\hat{L}$ by inverse multiplications on the right, as described in the previous section. At the end of the $(i-1)^{\text {th }}$ plane transformer $X_{i-1}, i=2, \ldots, n$, the situation is as indicated below:

$$
U K_{i-1}^{-1}=\left[\begin{array}{lllllll}
\ddots & & & & & & \\
\ddots & 1 & 0 & & & \\
& \hat{l}_{i-1} & d_{i} & 1 & & & \\
& & & u_{i+1} & 1 & \\
& & & & u_{i+2} & \ddots & \\
& & & & & \ddots
\end{array}\right], \quad K_{i-1} L=\left[\begin{array}{ccccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i-1} & 1 & & & & \\
& 0 & 1 & & & \\
& & & l_{i} & 1 & & \\
& & & & l_{i+1} & 1 & \\
& & & & & \ddots & \ddots
\end{array}\right] .
$$

where

$$
K_{i-1} \equiv X_{i-1} X_{i-2} \ldots X_{1}
$$

and $X_{i}$ is the plane transform given in (6.13).
The striking fact is that row $i$ of $K_{i-1} L$ and column $i$ of $U K_{i-1}^{-1}$ are singletons. That is

$$
\boldsymbol{e}_{i}^{T} K_{i-1} L=\boldsymbol{e}_{i}^{T} \quad \text { and } \quad U K_{i-1}^{-1} \boldsymbol{e}_{i}=\boldsymbol{e}_{i} d_{i}, \quad i=2, \ldots, n
$$

Rearranging these equations yields, for $i=2, \ldots, n$,

$$
\boldsymbol{e}_{i}^{T} K_{i-1}=\boldsymbol{e}_{i}^{T} L^{-1}, \quad K_{i-1}^{-1} \boldsymbol{e}_{i}=U^{-1} \boldsymbol{e}_{i} d_{i}
$$

and

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{i i} } & =\left(\boldsymbol{e}_{i}^{T} L^{-1}\right)\left(U^{-1} \boldsymbol{e}_{i}\right) \\
& =\left(\boldsymbol{e}_{i}^{T} K_{i-1}\right)\left(K_{i-1}^{-1} \boldsymbol{e}_{i} d_{i}^{-1}\right) \\
& =d_{i}^{-1} .
\end{aligned}
$$

Now for case $i=1$. We also have

$$
\boldsymbol{e}_{1}^{T} L=\boldsymbol{e}_{1}^{T} \quad \text { and } \quad U \boldsymbol{e}_{1}=\boldsymbol{e}_{1} u_{1}
$$

and, analogously,

$$
\left[(U L)^{-1}\right]_{11}=\left(\boldsymbol{e}_{1}^{T} L^{-1}\right)\left(U^{-1} \boldsymbol{e}_{1}\right)=\boldsymbol{e}_{1}^{T} u_{1}^{-1} \boldsymbol{e}_{1}=d_{1}^{-1}
$$

In the positive case $\left(l_{i}>0, u_{i}>0\right), U L$ is diagonally similar to a symmetric positive definite matrix (see lemma 2.2.3) and the quantity $\min \left|d_{j}\right|$ gives useful information on the eigenvalue nearest 0 .

Corollary 6.2.3 In the positive case

$$
\frac{\min _{i} d_{i}}{n}<\lambda_{\min }(U L) \leq \min _{i} d_{i} .
$$

Before writing the proof of this corollary we need to remember that if $A$ is a symmetric positive definite matrix then

$$
\operatorname{trace}(A)>\lambda_{\max }(A) \text { and } \lambda_{\max }(A) \geq A_{i i}
$$

This is simple to prove: eigenvalues $\lambda_{1}, \ldots, \lambda_{n}$ of $A$ are all positive and the first inequality is immediate. Also, $A$ is orthogonally similar to $D=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right)$. We have $A=Q^{T} D Q$, for an orthogonal matrix $Q$ and then

$$
\begin{aligned}
A_{i i} & =\left[Q^{T} D Q\right]_{i i} \\
& =\boldsymbol{e}_{i}^{T} Q^{T} D Q \boldsymbol{e}_{i} \\
& =\boldsymbol{q}_{i}^{T} D \boldsymbol{q}_{i} \\
& =\lambda_{1} q_{1 i}^{2}+\ldots+\lambda_{n} q_{n i}^{2} \\
& \leq \lambda_{\max }(A)\left\|\boldsymbol{q}_{i}\right\|_{2}^{2}=\lambda_{\max }(A)
\end{aligned}
$$

where we let $\boldsymbol{q}_{i}:=\left[q_{1 i}, \ldots, q_{n i}\right]^{T}$.
Also recall that for a matrix $M$ diagonally similar to $A$ we have

$$
M_{i i}=A_{i i}
$$

because a diagonal similarity does not affect the diagonal entries of $A$.
Now we write the proof of corollary 6.2.3.
Proof. For any matrix $M=\left(m_{i j}\right)_{1 \leq i, j \leq n}$ that is diagonally similar to a positive definite symmetric matrix we have

$$
\max _{i} m_{i i} \leq \lambda_{\max }(M)<\operatorname{trace}(M)
$$

If we let $M=(U L)^{-1}$ we have $\left[(U L)^{-1}\right]_{i i}=d_{i}^{-1}$ and then

$$
\left(\min _{i} d_{i}\right)^{-1}=\max _{i} d_{i}^{-1} \leq \lambda_{\max }\left((U L)^{-1}\right)<\sum_{i=1}^{n} d_{i}^{-1}<n\left(\min _{i} d_{i}\right)^{-1}
$$

Since $\left(\lambda_{\max }(U L)^{-1}\right)=\left(\lambda_{\min }(U L)\right)^{-1}$ we can write

$$
\left(\min _{i} d_{i}\right)^{-1} \leq\left(\lambda_{\min }(U L)\right)^{-1}<n\left(\min _{i} d_{i}\right)^{-1}
$$

or, equivalently,

$$
\frac{\min _{i} d_{i}}{n}<\lambda_{\min }(U L) \leq \min _{i} d_{i}
$$

Even in the general case, as $u_{n} \rightarrow 0, \min _{i}\left|d_{i}\right|$ becomes an increasingly accurate approximation to $\left|\lambda_{\text {min }}\right|$.

LR, QR and qd algorithms are only as good as their shift strategies. In practice, one uses qds and dqds, the shifted versions of qd and dqd.

The derivation of $\operatorname{dqds}(\sigma)$ in terms of Gram-Schmidt process is not obvious. Formally if $\left(U-\sigma L^{-1}\right) L$ admits triangular factorization we write

$$
U L-\sigma I=\left(U-\sigma L^{-1}\right) L=\hat{L} \hat{D} \hat{R}=\hat{L} \hat{U}
$$

and apply Gram-Schmidt to the columns of $L$ and to the rows of $U-\sigma L^{-1}$ to obtain

$$
U-\sigma L^{-1}=\hat{L} \tilde{P}^{T}, \quad L=\tilde{Q} \hat{R}, \quad \tilde{P}^{T} \tilde{Q}=\hat{D}
$$

If $\hat{D}$ is invertible, we can eliminate $\tilde{Q}$ :

$$
L=\left(\tilde{Q} \hat{D}^{-1}\right)(\hat{D} \hat{R})=\tilde{P}^{-T} \hat{U}, \quad U-\sigma L^{-1}=\hat{L} \tilde{P}^{T}
$$

that is

$$
K L=\hat{U}, \quad\left(U-\sigma L^{-1}\right) K^{-1}=\hat{L},
$$

with $K:=\tilde{P}^{T}$. At first glance the new term $-\sigma L^{-1}$ appears to spoil the derivation of $K$ as a product of plane rotations. However, it is not necessary to know all the terms of $L^{-1}$ but only the $(i+1, i)$ entries immediately below the main diagonal. The change for the unshifted case is small.

If one looks at the two matrices part way through the transformations $L \longrightarrow \hat{U}$, $U-\sigma L^{-1} \longrightarrow \hat{L}$, the singleton column in the second matrix (from theorem 6.2.2) has disappeared and the relation of $d_{i}$ to $(U L)^{-1}$ will be more complicated.

### 6.2.2 Complex eigenvalues and dqd

Theorem 6.2.2 gives information about the diagonal entries of $(U L)^{-1}$ that can be used in a choice of a shift strategy. In the presence of a pair of complex conjugate eigenvalues it will
be useful to know not just the diagonal entries but also the entries $(i+1, i)$ and $(i, i+1)$, that is, the diagonal $2 \times 2$ blocks of $(U L)^{-1}$.

We will use the notation $A_{i: j, k: l}$ to represent the submatrix of $A$ lying in rows $i$ through $j$ and columns $k$ through $l$.

Next result gives the entries of the $2 \times 2$ submatrix

$$
\begin{equation*}
\left[(U L)^{-1}\right]_{i: i+1, i: i+1} \tag{6.14}
\end{equation*}
$$

and then we will analyze when we may have complex eigenvalues.
Theorem 6.2.4 Consider $L$ and $U$ as described in (6.1). If $U$ is invertible and $d_{i}$, $i=1, \ldots, n$, are the intermediate quantities generated by the dqd algorithm applied to $L$ and $U$, then

$$
\left[(U L)^{-1}\right]_{i: i+1, i: i+1}=\left[\begin{array}{cc}
\frac{1}{d_{i}} & -\frac{1}{d_{i} u_{i+1}} \\
-\frac{l_{i}}{d_{i}} & \frac{1}{d_{i+1}}
\end{array}\right], \quad i=1, \ldots, n-1 .
$$

Proof. From theorem 6.2.2

$$
\frac{1}{d_{i}}=\left[(U L)^{-1}\right]_{i i}, \quad i=1, \ldots, n .
$$

Also from the proof of theorem 6.2.2, we have

$$
\boldsymbol{e}_{i}^{T} K_{i-1}=\boldsymbol{e}_{i}^{T} L^{-1} \quad \text { and } \quad d_{i}^{-1} K_{i-1}^{-1} \boldsymbol{e}_{i}=U^{-1} \boldsymbol{e}_{i}, \quad i=2, \ldots, n,
$$

where

$$
K_{i-1}=X_{i-1} X_{i-2} \ldots X_{1}, \quad i=2, \ldots, n-1
$$

and $X_{i}$ is given by (6.13).
Thus,

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{i, i+1} } & =\left(e_{i}^{T} L^{-1}\right)\left(U^{-1} \boldsymbol{e}_{i+1}\right) \\
& =\left(\boldsymbol{e}_{i}^{T} K_{i-1}\right)\left(d_{i+1}^{-1} K_{i}^{-1} \boldsymbol{e}_{i+1}\right) \\
& =d_{i+1}^{-1} \boldsymbol{e}_{i}^{T}\left(K_{i-1} K_{i}^{-1}\right) \boldsymbol{e}_{i+1} \\
& =d_{i+1}^{-1}\left(\boldsymbol{e}_{i}^{T} X_{i}^{-1} \boldsymbol{e}_{i+1}\right) \\
& =-\frac{1}{d_{i+1}} \frac{1}{\hat{u}_{i}} .
\end{aligned}
$$

In dqd algorithm we have

$$
d_{i+1}=d_{i}\left(u_{i+1} / \hat{u}_{i}\right)
$$

and then

$$
\left[(U L)^{-1}\right]_{i, i+1}=-\frac{1}{d_{i} u_{i+1}}
$$

Analogously,

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{i+1, i} } & =\left(\boldsymbol{e}_{i+1}^{T} L^{-1}\right)\left(U^{-1} \boldsymbol{e}_{i}\right) \\
& =\left(\boldsymbol{e}_{i+1}^{T} K_{i}\right)\left(d_{i}^{-1} K_{i-1}^{-1} \boldsymbol{e}_{i}\right) \\
& =d_{i}^{-1} \boldsymbol{e}_{i+1}^{T}\left(K_{i} K_{i-1}^{-1}\right) \boldsymbol{e}_{i} \\
& =d_{i}^{-1}\left(\boldsymbol{e}_{i+1}^{T} X_{i} \boldsymbol{e}_{i}\right) \\
& =-\frac{l_{i}}{d_{i}}
\end{aligned}
$$

For $i=1$, notice that $d_{1}=u_{1}$ and we have

$$
\boldsymbol{e}_{1}^{T}=L^{-1} \boldsymbol{e}_{1}^{T} \quad \text { and } \quad u_{1}^{-1} \boldsymbol{e}_{1}=U^{-1} \boldsymbol{e}_{1}
$$

So,

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{1,2} } & =\left(\boldsymbol{e}_{1}^{T} L^{-1}\right)\left(U^{-1} \boldsymbol{e}_{2}\right) \\
& =\boldsymbol{e}_{1}^{T}\left(d_{2}^{-1} K_{1}^{-1} \boldsymbol{e}_{2}\right) \\
& =d_{2}^{-1}\left(\boldsymbol{e}_{1}^{T} K_{1}^{-1} \boldsymbol{e}_{2}\right) \\
& =d_{2}^{-1}\left(\boldsymbol{e}_{2}^{T} X_{1}^{-1} \boldsymbol{e}_{2}\right) \\
& =-\frac{1}{d_{2}} \frac{1}{\hat{u}_{1}}=-\frac{1}{d_{1} u_{2}}
\end{aligned}
$$

and, finally,

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{2,1} } & =\left(e_{2}^{T} L^{-1}\right)\left(U^{-1} e_{1}\right) \\
& =\left(e_{2}^{T} K_{1}\right)\left(d_{1}^{-1} e_{1}\right) \\
& =d_{1}^{-1}\left(\boldsymbol{e}_{2}^{T} K_{1} e_{1}\right) \\
& =d_{1}^{-1}\left(\boldsymbol{e}_{2}^{T} X_{1} e_{1}\right) \\
& =-\frac{l_{1}}{d_{1}} .
\end{aligned}
$$

Now we are interested in finding the eigenvalues of the $2 \times 2$ submatrix of theorem 6.2.4. It can be verified that the discriminant of the characteristic polynomial is

$$
\begin{equation*}
\Delta_{i} \equiv\left(\frac{d_{i+1}-d_{i}}{d_{i} d_{i+1}}\right)^{2}+\frac{4 l_{i}}{d_{i}^{2} u_{i+1}} . \tag{6.15}
\end{equation*}
$$

Lemma 6.2.1 If $l_{i} u_{i+1}>0$ then the $2 \times 2$ matrix in (6.14) has real eigenvalues.

Proof. According to (6.15), if $l_{i} u_{i+1} \geq 0$ then $\Delta_{i} \geq 0$.
If $l_{i} u_{i+1}<0$ we must evaluate $\Delta_{i}$ to know if we are in the presence of complex eigenvalues.

### 6.3 New version of triple dqds

Consider matrices $L$ and $U$ as given in (6.1) and the application of tridqds algorithm. Initially, we define

$$
F:=U \quad \text { and } \quad G:=L
$$

and

$$
F^{(1)}:=F \quad \text { and } \quad G^{(1)}:=G \text {. }
$$

At the beginning of each minor step $i, i=2, \ldots, n$, we have

$$
\begin{aligned}
F^{(i)} & :=\mathcal{L}_{i-1}^{-1} F^{(i-1)} Z_{i-1}^{-1} Y_{i-1}^{-1} \\
G^{(i)} & :=Y_{i-1} Z_{i-1} G^{(i-1)} \mathcal{L}_{i-1}
\end{aligned}
$$

Therefore

$$
\begin{aligned}
F^{(i)} & :=\left(\mathcal{L}_{i-1}^{-1} \ldots \mathcal{L}_{2}^{-1} \mathcal{L}_{1}^{-1}\right) U\left(Z_{1}^{-1} Y_{1}^{-1} Z_{2}^{-1} Y_{2}^{-1} \ldots Z_{i-1}^{-1} Y_{i-1}^{-1}\right) \\
G^{(i)} & :=\left(Y_{i-1} Z_{i-1} \ldots Y_{2} Z_{2} Y_{1} Z_{1}\right) L\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)
\end{aligned}
$$

and

$$
\begin{equation*}
F^{(i)} G^{(i)}=\left(\mathcal{L}_{i-1}^{-1} \ldots \mathcal{L}_{2}^{-1} \mathcal{L}_{1}^{-1}\right)(U L)\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right) \tag{6.16}
\end{equation*}
$$

As established in (5.12), page 143 , at the beginning of minor step $i=2, \ldots, n-2$, of tridqds algorithm we have

$$
F^{(i)}=\left[\begin{array}{llllll}
\ddots & & & & &  \tag{6.17}\\
\ddots & 1 & & & & \\
& \hat{l}_{i-1} & u_{i} & 1 & & \\
& f_{1}^{i-1} & & u_{i+1} & 1 & \\
& f_{2}^{i-1} & & & u_{i+2} & \ddots \\
& & & & \ddots
\end{array}\right], \quad G^{(i)}=\left[\begin{array}{cccccc}
\ddots & \ddots & & & & \\
& \hat{u}_{i-1} & 1 & & & \\
& & d_{i} & & & \\
& & g_{1}^{i} & 1 & & \\
& & & g_{2}^{i} & l_{i+1} & 1
\end{array}\right]
$$

where we defined

$$
\begin{array}{rlrl}
f_{1}^{i-1}:=x_{l} \quad \text { and } \quad d_{i} & :=x_{r} \\
f_{2}^{i-1}:=y_{l} & & g_{1}^{i} & :=y_{r} \\
& g_{2}^{i} & :=z_{r} .
\end{array}
$$

The superscript refers to the column at which the element is positioned.

At the beginning of minor step $n-1$,

$$
F^{n-1}=\left[\begin{array}{ccccc}
\ddots & \ddots & & & \\
& \hat{l}_{n-3} & 1 & & \\
& & \hat{l}_{n-2} & u_{n-1} & 1 \\
& & f_{1}^{n-2} & & u_{n}
\end{array}\right], \quad G^{n-1}=\left[\begin{array}{ccccc}
\ddots & \ddots & & & \\
& \hat{u}_{n-3} & 1 & & \\
& & \hat{u}_{n-2} & 1 & \\
& & & d_{n-1} & \\
& & & g_{1}^{n-1} & 1
\end{array}\right]
$$

and for minor step $n$

$$
F_{n}=\left[\begin{array}{cccc}
\ddots & & & \\
\ddots & 1 & & \\
& \hat{l}_{n-2} & 1 & \\
& & \hat{l}_{n-1} & u_{n}
\end{array}\right], \quad G_{n}=\left[\begin{array}{cccc}
\ddots & \ddots & & \\
& \hat{u}_{n-2} & 1 & \\
& & \hat{u}_{n-1} & 1 \\
& & & d_{n}
\end{array}\right] .
$$

Since initially we let

$$
F^{(1)}=U \quad \text { and } \quad G^{(1)}=L,
$$

then

$$
d_{1}=1, \quad g_{1}^{1}=l_{1} \quad \text { and } \quad g_{2}^{1}=0 .
$$

### 6.3.1 New notation for tridqds

Next we will rewrite tridqds algorithm with the notation introduced in (6.17). First column of $F^{(1)}$ is $u_{1} \boldsymbol{e}_{1}$ and quantities $f_{1}^{0}$ and $f_{2}^{0}$ presented at minor step 1 are just auxiliary variables used for clarity and do not belong to $F^{(1)}$.

The first minor step of tridqds.

Minor step 1 :

$$
\begin{aligned}
& d_{1}=1 ; g_{1}^{1}=l_{1} ; g_{2}^{1}=0 \\
& d_{2}=d_{1} * u_{1}+g_{1}^{1} \\
& f_{1}^{0}=\left(u_{1}+l_{1}\right)^{2}+u_{2} l_{1}-2(\Re \sigma)\left(u_{1}+l_{1}\right)+|\sigma|^{2} \\
& f_{2}^{0}=u_{2} l_{1} u_{3} l_{2} / f_{1}^{0} \\
& f_{1}^{0}=u_{2} l_{1}\left(u_{1}+l_{1}+u_{2}+l_{2}-2(\Re \sigma)\right) / f_{1}^{0} \\
& f_{1}^{1}=-f_{1}^{0} ; \quad f_{2}^{1}=-f_{2}^{0} \\
& \hat{u}_{1}=d_{2}-f_{1}^{1} ; \\
& d_{2}=g_{1}^{1}-f_{1}^{1} ; g_{1}^{2}=g_{2}^{1}-f_{2}^{1}-f_{1}^{1} * l_{2} ; g_{2}^{2}=-f_{2}^{1} * l_{3} \\
& d_{2}=d_{2} / \hat{u}_{1} ; g_{1}^{2}=g_{1}^{2} / \hat{u}_{1} ; \quad g_{2}^{2}=g_{2}^{2} / \hat{u}_{1} \\
& \hat{l}_{1}=f_{1}^{1}+g_{1}^{2}+d_{2} * u_{2} \\
& f_{1}^{1}=f_{2}^{1}+g_{2}^{2}+g_{1}^{2} * u_{3} ; \quad f_{2}^{1}=g_{2}^{2} * u_{4} \\
& d_{2}=1-d_{2} ; g_{1}^{2}=l_{2}-g_{1}^{2} ; g_{2}^{2}=-g_{2}^{2}
\end{aligned}
$$

The inner loop is changed to:

$$
\begin{aligned}
& \text { for } i=2, \ldots, n-3 \\
& \qquad \begin{array}{l}
d_{i+1}=d_{i} * u_{i}+g_{1}^{i} \\
f_{1}^{i}=-f_{1}^{i-1} / \hat{l}_{i-1} ; \quad f_{2}^{i}=-f_{2}^{i-1} / \hat{l}_{i-1} ; \\
\hat{u}_{i}=d_{i+1}-f_{1}^{i} ; \\
d_{i+1}=g_{1}^{i}-f_{1}^{i} ; \quad g_{1}^{i+1}=g_{2}^{i}-f_{2}^{i}-f_{1}^{i} * l_{i+1} ; \quad g_{2}^{i+1}=-f_{2}^{i} * l_{i+2} \\
d_{i+1}=d_{i+1} / \hat{u} i ; \quad g_{1}^{i+1}=g_{1}^{i+1} / \hat{u_{i}} ; \quad g_{2}^{i+1}=g_{2}^{i+1} / \hat{u}_{i} \\
\hat{l}_{i}=f_{1}^{i}+g_{1}^{i+1}+d_{i+1} * u_{i+1} \\
f_{1}^{i}=f_{2}^{i}+g_{2}^{i+1}+g_{1}^{i+1} * u_{i+2} ; \quad f_{2}^{i}=g_{2}^{i+1} * u_{i+3} \\
d_{i+1}=1-d_{i+1} ; \quad g_{1}^{i+1}=l_{i+1}-g_{1}^{i+1} ; g_{2}^{i+1}=-g_{2}^{i+1}
\end{array}
\end{aligned}
$$

## end for

We begin the loop with quantities $f_{1}^{i-1}, f_{2}^{i-1}, d_{i}, g_{1}^{i}$ and $g_{2}^{i}$ and end with quantities $f_{1}^{i}$, $f_{2}^{i}, d_{i+1}, g_{1}^{i+1}$ and $g_{2}^{i+1}$, and all values are preserved. For completeness we will also write down the minor steps $n-2, n-1$ and $n$ with this new notation.

## Minor step n-2 :

$$
\begin{aligned}
& d_{n-1}=d_{n-2} * u_{n-2}+g_{1}^{n-2} \\
& f_{1}^{n-2}=-f_{1}^{n-3} / \hat{l}_{n-3} ; \quad f_{2}^{n-2}=-f_{2}^{n-3} / \hat{l}_{n-3} ; \\
& \hat{u}_{n-2}=d_{n-1}-f_{1}^{n-2} ; \\
& d_{n-1}=g_{1}^{n-2}-f_{1}^{n-2} ; \quad g_{1}^{n-1}=g_{2}^{n-2}-f_{2}^{n-2}-f_{1}^{n-2} * l_{n-1} \\
& d_{n-1}=d_{n-1} / \hat{u}_{n-2} ; g_{1}^{n-1}=g_{1}^{n-1} / \hat{u}_{n-2} \\
& \hat{l}_{n-2}=f_{1}^{n-2}+g_{1}^{n-1}+d_{n-1} * u_{n-1} \\
& f_{1}^{n-2}=f_{2}^{n-2}+g_{1}^{n-1} * u_{n} \\
& d_{n-1}=1-d_{n-1} ; g_{1}^{n-1}=l_{n-1}-g_{1}^{n-1}
\end{aligned}
$$

Minor step n-1 :

$$
\begin{aligned}
& d_{n}=d_{n-1} * u_{n-1}+g_{1}^{n-1} \\
& f_{1}^{n-1}=-f_{1}^{n-2} / \hat{l}_{n-2} \\
& \hat{u}_{n-1}=d_{n}-f_{1}^{n-1} \\
& d_{n}=g_{1}^{n-1}-f_{1}^{n-1} \\
& d_{n}=d_{n} / \hat{u}_{n-1} \\
& \hat{l}_{n-1}=f_{1}^{n-1}+d_{n} * u_{n} \\
& d_{n}=1-d_{n}
\end{aligned}
$$

Minor step n :

$$
\begin{aligned}
& \hat{u}_{n}=d_{n} * u_{n} \\
& \hat{u}_{n}=\hat{u}_{n} ;
\end{aligned}
$$

To represent elementary matrix $\mathcal{L}_{i}$ at minor step $i, i=1, \ldots, n-2$, we will use the notation

$$
\mathcal{L}_{i}=\left[\begin{array}{cccccc}
1 & & & & &  \tag{6.18}\\
& \ddots & & & & \\
& & 1 & & & \\
& & h_{1}^{i} & 1 & & \\
& & h_{2}^{i} & & 1 & \\
& & & & & \ddots \\
& & & & & \\
& & & & & \\
&
\end{array}\right]
$$

Elements $h_{1}^{i}$ and $h_{2}^{i}$ are at positions $(i+1, i)$ and $(i+2, i)$, respectively. For $i=n-1$ and $i=n$,

$$
\mathcal{L}_{n-1}=\left[\begin{array}{cccccc}
1 & & & & & \\
& \ddots & & & & \\
& & 1 & & & \\
& & & \ddots & & \\
& & & & 1 & \\
& & & & h_{1}^{n-1} & 1
\end{array}\right] \quad \text { and } \quad \mathcal{L}_{n}=I .
$$

We will write

$$
\mathcal{L}_{i}=I+\boldsymbol{h}_{i} \boldsymbol{e}_{i}^{T}, \quad i=1, \ldots, n-1
$$

where

It is well known that

$$
\mathcal{L}_{i}^{-1}=I-\boldsymbol{h}_{i} \boldsymbol{e}_{i}^{T}
$$

and

$$
\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i}=\left[\begin{array}{cccccccc}
1 & & & & & & & \\
h_{1}^{1} & 1 & & & & & & \\
h_{2}^{1} & h_{1}^{2} & \ddots & & & & & \\
& h_{2}^{2} & \ddots & 1 & & & & \\
& & \ddots & h_{1}^{i} & 1 & & & \\
& & & h_{2}^{i} & & 1 & & \\
& & & & & \ddots & \\
& & & & & & & 1
\end{array}\right] .
$$

In what follows we will need to refer to the columns and to the rows of $\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i}$. For $i=1, \ldots, n$, the columns are

$$
\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i}\right) \boldsymbol{e}_{k}=\left\{\begin{array}{lllllllll} 
& \overbrace{}^{k^{\text {th }}}  \tag{6.19}\\
{\left[\begin{array}{llllllll}
0 & \ldots & 0 & 1 & h_{1}^{k} & h_{2}^{k} & 0 & \ldots
\end{array}\right.} & 0
\end{array}\right]^{T}, \quad k \leq i .
$$

Note that $\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{n-1}\right) \boldsymbol{e}_{n-1}=\left[\begin{array}{lllll}0 & \ldots & 0 & 1 & h_{1}^{n-1}\end{array}\right]^{T}$ and $\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{n}\right) \boldsymbol{e}_{n}=\boldsymbol{e}_{n}$. And for $i=1, \ldots, n$, the rows are given by

$$
\boldsymbol{e}_{k}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i}\right)=\left\{\begin{array}{lllllllll} 
& \overbrace{0}^{k^{\text {th }}}  \tag{6.20}\\
0 & \ldots & 0 & h_{2}^{k-2} & h_{1}^{k-1} & 1 & 0 & \ldots & 0
\end{array}\right], \quad k \leq i+1 . \quad k=i+2
$$

Note that $\boldsymbol{e}_{1}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i}\right)=\boldsymbol{e}_{1}^{T}$ and $\boldsymbol{e}_{2}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i}\right)=\left[\begin{array}{lllll}h_{1}^{1} & 1 & 0 & \ldots & 0\end{array}\right]$.

### 6.3.2 From tridqds to $3 d q d s$

The next sequence of lemmas will permit us to change the tridqds algorithm into an equivalent but more efficient format. This is a central contribution of the thesis and lead us to the new version, 3dqds, that will be presented in page 201.

Lemma 6.3.1 Consider tridqds algorithm applied to matrices $L$ and $U$ as given in (6.1). For $\mathcal{L}_{1}$ we have

$$
h_{1}^{1}=f_{1}^{0} \quad \text { and } \quad h_{2}^{1}=f_{2}^{0}
$$

For $\mathcal{L}_{i}$ at minor step $i, i=2, \ldots, n-2$, we have

$$
h_{1}^{i}=f_{1}^{i-1} / \hat{l}_{i-1} \quad \text { and } \quad h_{2}^{i}=f_{2}^{i-1} / \hat{l}_{i-1} .
$$

Finally, for $\mathcal{L}_{n-1}$

$$
h_{1}^{n-1}=f_{1}^{n-2} / \hat{l}_{n-2}
$$

Proof. In tridqds algorithm the inverse $\mathcal{L}_{i}^{-1}, i=2, \ldots, n-2$, is chosen to zero out entries $(i+1, i-1)$ and $(i+2, i-1)$ of $F^{(i)}$ using entry $(i, i-1)$ as the pivot. Thus

$$
\mathcal{L}_{i}^{-1}=\left[\begin{array}{cccccc}
\ddots & & & & &  \tag{6.21}\\
& \ddots & & & & \\
& & 1 & & \\
& & -f_{1}^{i-1} / \hat{l}_{i-1} & 1 \\
& -f_{2}^{i-1} / \hat{l}_{i-1} & & & \\
& & & & & \ddots
\end{array}\right]
$$

So, entries $h_{1}^{i}$ and $h_{2}^{i}$ in $\mathcal{L}_{i}$ are the symmetric of the corresponding entries in $\mathcal{L}_{i}^{-1} . \mathcal{L}_{n-1}$ is chosen to zero out entry $(n, n-2)$ of $F^{n-1}$. The case $i=1$ is immediate (see minor step 1 to observe the initialization of $f_{1}^{0}$ and $f_{2}^{0}$ ).

In all the results that follow we will consider $n \geq 4$.

Lemma 6.3.2 Consider tridqds algorithm and matrix $F^{(i)}$ at the beginning of minor step $i$, $i=2, \ldots, n-1$. Then, for $i=2$,

$$
f_{1}^{1}=\frac{u_{3} l_{2}}{\hat{u}_{1}} f_{1}^{0}+\frac{l_{3}+u_{3}-\hat{u}_{1}}{\hat{u}_{1}} f_{2}^{0}+\frac{u_{3}}{\hat{u}_{1}} g_{2}^{1} \quad \text { and } \quad f_{2}^{1}=\frac{u_{4} l_{3}}{\hat{u}_{1}} f_{2}^{0}
$$

For $i=3, \ldots, n-2$,

$$
f_{1}^{i-1}=\frac{u_{i+1} l_{i}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{1}^{i-2}+\frac{l_{i+1}+u_{i+1}-\hat{u}_{i-1}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{2}^{i-2}+\frac{u_{i+1}}{\hat{u}_{i-1}} g_{2}^{i-1}
$$

and

$$
f_{2}^{i-1}=\frac{u_{i+2} l_{i+1}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{2}^{i-2}
$$

At the beginning of minor step $n-1$,

$$
f_{1}^{n-2}=\frac{u_{n} l_{n-1}}{\hat{l}_{n-3} \hat{u}_{n-2}} f_{1}^{n-3}+\frac{u_{n}-\hat{u}_{n-2}}{\hat{l}_{n-3} \hat{u}_{n-2}} f_{2}^{n-3}+\frac{u_{n}}{\hat{u}_{n-2}} g_{2}^{n-2}
$$

Proof. At the end of minor step 1 we obtain

$$
\begin{aligned}
f_{1}^{1} & =-f_{2}^{0}+\left[\left(f_{2}^{0} * l_{3}\right) / \hat{u}_{1}\right]+\left[\left(g_{2}^{1}+f_{2}^{0}+f_{1}^{0} * l_{2}\right) / \hat{u}_{1}\right] * u_{3} \\
& =\frac{u_{3} l_{2}}{\hat{u}_{1}} f_{1}^{0}+\frac{l_{3}+u_{3}-\hat{u}_{1}}{\hat{u}_{1}} f_{2}^{0}+\frac{u_{3}}{\hat{u}_{1}} g_{2}^{1}
\end{aligned}
$$

and

$$
f_{2}^{1}=\left[\left(f_{2}^{0} * l_{3}\right) / \hat{u}_{1}\right] * u_{4}=\frac{u_{4} l_{3}}{\hat{u}_{1}} f_{2}^{0}
$$

The sequence of assignments to compute $f_{2}^{i}, i=2, \ldots, n-3$, is

$$
\begin{gathered}
f_{2}^{i} \longleftarrow-f_{2}^{i-1} / \hat{l}_{i-1} \\
g_{2}^{i+1} \longleftarrow-f_{2}^{i} * l_{i+2} \\
g_{2}^{i+1} \longleftarrow g_{2}^{i+1} / \hat{u}_{i} \\
f_{2}^{i} \longleftarrow g_{2}^{i+1} * u_{i+3} .
\end{gathered}
$$

Doing forward substitution we get

$$
\begin{aligned}
f_{2}^{i} & =\left\{\left[-\left(-f_{2}^{i-1} / \hat{l}_{i-1}\right) * l_{i+2}\right] / \hat{u}_{i}\right\} * u_{i+3} \\
& =\frac{f_{2}^{i-1} l_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} * u_{i+3} \\
& =\frac{u_{i+3} l_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} f_{2}^{i-1}
\end{aligned}
$$

So, in the beginning of minor step $i, i=3, \ldots, n-2$, we have

$$
f_{2}^{i-1}=\frac{u_{i+2} l_{i+1}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{2}^{i-2} .
$$

Analogously, to obtain $f_{1}^{i}, i=2, \ldots, n-3$, the sequence is

$$
\begin{aligned}
f_{1}^{i} \longleftarrow-f_{1}^{i-1} / \hat{l}_{i-1} \\
f_{2}^{i} \longleftarrow-f_{2}^{i-1} / \hat{l}_{i-1} \\
g_{1}^{i+1} \longleftarrow g_{2}^{i}-f_{2}^{i}-f_{1}^{i} * l_{i+1} ; \\
g_{2}^{i+1} \longleftarrow-f_{2}^{i} * l_{i+2} \\
g_{1}^{i+1} \longleftarrow g_{1}^{i+1} / \hat{u}_{i} ; \\
g_{2}^{i+1} \longleftarrow g_{2}^{i+1} / \hat{u}_{i} \\
f_{1}^{i} \longleftarrow f_{2}^{i}+g_{2}^{i+1}+g_{1}^{i+1} * u_{i+2} .
\end{aligned}
$$

Thus, after minor step $i$,

$$
\begin{aligned}
f_{1}^{i}= & \left(-f_{2}^{i-1} / \hat{l}_{i-1}\right)+\left\{\left[-\left(-f_{2}^{i-1} / \hat{l}_{i-1}\right) * l_{i+2}\right] / \hat{u}_{i}\right\}+ \\
& +\left\{\left[g_{2}^{i}-\left(-f_{2}^{i-1} / \hat{l}_{i-1}\right)-\left(-f_{1}^{i-1} / \hat{l}_{i-1}\right) * l_{i+1}\right] / \hat{u}_{i}\right\} * u_{i+2} \\
= & -\frac{f_{2}^{i-1}}{\hat{l}_{i-1}}+\frac{f_{2}^{i-1} l_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}}+\left(\frac{g_{2}^{i}}{\hat{u}_{i}}+\frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}+\frac{f_{1}^{i-1} l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}}\right) u_{i+2} \\
= & -\frac{1}{\hat{l}_{i-1}} f_{2}^{i-1}+\frac{l_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} f_{2}^{i-1}+\frac{u_{i+2}}{\hat{u}_{i}} g_{2}^{i}+\frac{u_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} f_{2}^{i-1}+\frac{u_{i+2} l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1} \\
= & \frac{u_{i+2} l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{l_{i+2}+u_{i+2}-\hat{u}_{i}}{\hat{l}_{i-1} \hat{u}_{i}} f_{2}^{i-1}+\frac{u_{i+2}}{\hat{u}_{i}} g_{2}^{i} .
\end{aligned}
$$

In the beginning of minor step $i, i=3, \ldots, n-2$, we have

$$
f_{1}^{i-1}=\frac{u_{i+1} l_{i}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{1}^{i-2}+\frac{l_{i+1}+u_{i+1}-\hat{u}_{i-1}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{2}^{i-2}+\frac{u_{i+1}}{\hat{u}_{i-1}} g_{2}^{i-1} .
$$

At the end of minor step $i=n-2$, or at the beginning of minor step $n-1$,

$$
\begin{aligned}
f_{1}^{n-2} & =\left(-f_{2}^{n-3} / \hat{l}_{n-3}\right)+\left\{\left[g_{2}^{n-2}-\left(-f_{2}^{n-3} / \hat{l}_{n-3}\right)-\left(-f_{1}^{n-3} / \hat{l}_{n-3}\right) * l_{n-1}\right] / \hat{u}_{n-2}\right\} * u_{n} \\
& =\frac{u_{n} l_{n-1}}{\hat{l}_{n-3} \hat{u}_{n-2}} f_{1}^{n-3}+\frac{u_{n}-\hat{u}_{n-2}}{\hat{l}_{n-3} \hat{u}_{n-2}} f_{2}^{n-3}+\frac{u_{n}}{\hat{u}_{n-2}} g_{2}^{n-2} . \square
\end{aligned}
$$

Lemma 6.3.3 Consider matrix $G^{(i)}$ at the beginning of minor step $i, i=2, \ldots, n-1$, of tridqds algorithm. Then, for $i=2$,

$$
g_{1}^{2}=l_{2}-\left(\frac{l_{2} f_{1}^{0}}{\hat{u}_{1}}+\frac{f_{2}^{0}}{\hat{u}_{1}}+\frac{g_{2}^{1}}{\hat{u}_{1}}\right) \quad \text { and } \quad g_{2}^{2}=-\frac{l_{3}}{\hat{u}_{1}} f_{2}^{0} .
$$

For $i=3, \ldots, n-1$,

$$
g_{1}^{i}=l_{i}-\left(\frac{l_{i}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{1}^{i-2}+\frac{f_{2}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}+\frac{g_{2}^{i-1}}{\hat{u}_{i-1}}\right)
$$

and for $i=3, \ldots, n-2$,

$$
g_{2}^{i}=-\frac{l_{i+1}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{2}^{i-2}
$$

Proof. At the beginning of minor step $i=2$, from minor step 1, we have

$$
g_{2}^{2}=-\left[\left(f_{2}^{0} * l_{3}\right) / \hat{u}_{1}\right]=-\frac{l_{3}}{\hat{u}_{1}} f_{2}^{0}
$$

and

$$
g_{1}^{2}=l_{2}-\left[\left(g_{2}^{1}+f_{2}^{0}+f_{1}^{0} * l_{2}\right) / \hat{u}_{1}\right]=l_{2}-\left(\frac{l_{2} f_{1}^{0}}{\hat{u}_{1}}+\frac{f_{2}^{0}}{\hat{u}_{1}}+\frac{g_{2}^{1}}{\hat{u}_{1}}\right) .
$$

The sequence of assignments to compute $g_{2}^{i+1}, i=2, \ldots, n-3$, is

$$
\begin{aligned}
& f_{2}^{i} \longleftarrow-f_{2}^{i-1} / \hat{l}_{i-1} \\
& g_{2}^{i+1} \longleftarrow-f_{2}^{i} * l_{i+2} \\
& g_{2}^{i+1} \longleftarrow g_{2}^{i+1} / \hat{u}_{i} \\
& g_{2}^{i+1} \longleftarrow-g_{2}^{i+1} .
\end{aligned}
$$

Again, substituting forwards, we get

$$
g_{2}^{i+1}=-\left\{\left[-\left(-f_{2}^{i-1} / \hat{l}_{i-1}\right) * l_{i+2}\right] / \hat{u}_{i}\right\}=-\frac{l_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} f_{2}^{i-1}
$$

Thus in the beginning of minor step $i=3, \ldots, n-2$ we have

$$
g_{2}^{i}=-\frac{l_{i+1}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{2}^{i-2} .
$$

Now the sequence of assignments to compute $g_{1}^{i+1}, i=2, \ldots, n-3$, is

$$
\begin{gathered}
f_{1}^{i} \longleftarrow-f_{1}^{i-1} / \hat{l}_{i-1} \\
f_{2}^{i} \longleftarrow-f_{2}^{i-1} / \hat{l}_{i-1} \\
g_{1}^{i+1} \longleftarrow g_{2}^{i}-f_{2}^{i}-f_{1}^{i} * l_{i+1} \\
g_{1}^{i+1} \longleftarrow g_{1}^{i+1} / \hat{u}_{i} \\
g_{1}^{i+1} \longleftarrow l_{i+1}-g_{1}^{i+1}
\end{gathered}
$$

and we obtain

$$
\begin{aligned}
g_{1}^{i+1} & =l_{i+1}-\left\{\left[g_{2}^{i}-\left(-f_{2}^{i-1} / \hat{l}_{i-1}\right)-\left(-f_{1}^{i-1} / \hat{l}_{i-1}\right) * l_{i+1}\right] / \hat{u}_{i}\right\} \\
& =l_{i+1}-\frac{g_{2}^{i}}{\hat{u}_{i}}-\frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}-\frac{f_{1}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}} l_{i+1} \\
& =l_{i+1}-\left(\frac{l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}+\frac{g_{2}^{i}}{\hat{u}_{i}}\right) .
\end{aligned}
$$

Thus, in the beginning of minor step $i, i=3, \ldots, n-2$, we have

$$
g_{1}^{i}=l_{i}-\left(\frac{l_{i}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{1}^{i-2}+\frac{f_{2}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}+\frac{g_{2}^{i-1}}{\hat{u}_{i-1}}\right) .
$$

Finally, at the beginning of minor step $n-1$,

$$
\begin{aligned}
g_{1}^{n-1} & =l_{n-1}-\left\{\left[g_{2}^{n-2}-\left(-f_{2}^{n-3} / \hat{l}_{n-3}\right)-\left(-f_{1}^{n-3} / \hat{l}_{n-3}\right) * l_{n-1}\right] / \hat{u}_{n-2}\right\} \\
& =l_{n-1}-\left(\frac{l_{n-1}}{\hat{l}_{n-3} \hat{u}_{n-2}} f_{1}^{n-3}+\frac{f_{2}^{n-3}}{\hat{l}_{n-3} \hat{u}_{n-2}}+\frac{g_{2}^{n-2}}{\hat{u}_{n-2}}\right) .
\end{aligned}
$$

Lemma 6.3.4 Let $F^{(i)}$ and $G^{(i)}$ be the matrices obtained at the beginning of minor step $i$, $i=2, \ldots, n-1$, when we apply tridqds algorithm to $L$ and $U$. Then

$$
f_{2}^{i-1}=-u_{i+2} g_{2}^{i},
$$

for $i=2, \ldots, n-2$, and

$$
f_{2}^{n-2}=-f_{2}^{n-3} / \hat{l}_{n-3} .
$$

Proof. From lemmas 6.3 .2 and 6.3 .3 we know that, for $i=3, \ldots, n-2$,

$$
f_{2}^{i-1}=\frac{u_{i+2} l_{i+1}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{2}^{i-2} \quad \text { and } \quad g_{2}^{i}=-\frac{l_{i+1}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{2}^{i-2} .
$$

The result follows immediately.
For $i=2$, that is at the end of minor step 1 , we have in sequence

$$
\begin{aligned}
& f_{2}^{1} \longleftarrow g_{2}^{2} * u_{4} \\
& g_{2}^{2} \longleftarrow-g_{2}^{2} .
\end{aligned}
$$

Thus,

$$
g_{2}^{2}=-\frac{f_{2}^{1}}{u_{4}} \Leftrightarrow f_{2}^{1}=-u_{4} g_{2}^{2} .
$$

For $i=n-1, f_{2}^{n-2}=-f_{2}^{n-3} / \hat{l}_{n-3}$ comes from the last assignment for $f_{2}^{n-2}$ in step $n-2$.

Lemma 6.3.5 Let $F^{(i)}$ and $G^{(i)}$ be the matrices obtained at the beginning of each minor step $i, i=2, \ldots, n$, when we apply tridqds algorithm to $L$ and $U$. Then

$$
f_{1}^{1}=u_{3}\left(l_{2}-g_{1}^{2}\right)-g_{2}^{2}-f_{2}^{0},
$$

for $i=2, \ldots, n-3$,

$$
f_{1}^{i}=u_{i+2}\left(l_{i+1}-g_{1}^{i+1}\right)-g_{2}^{i+1}-\frac{f_{2}^{i-1}}{\hat{l}_{i-1}}
$$

and

$$
f_{1}^{n-2}=u_{n}\left(l_{n-1}-g_{1}^{n-1}\right)-\frac{f_{2}^{n-3}}{\hat{l}_{n-3}}, \quad f_{1}^{n-1}=-\frac{f_{1}^{n-2}}{\hat{l}_{n-2}} .
$$

Proof. From lemmas 6.3.2 and 6.3.3 we know that, for $i=2, \ldots, n-3$,

$$
f_{1}^{i}=\frac{u_{i+2} l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{l_{i+2}+u_{i+2}-\hat{u}_{i}}{\hat{l}_{i-1} \hat{u}_{i}} f_{2}^{i-1}+\frac{u_{i+2}}{\hat{u}_{i}} g_{2}^{i}
$$

and

$$
g_{1}^{i+1}=l_{i+1}-\left(\frac{l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}+\frac{g_{2}^{i}}{\hat{u}_{i}}\right) .
$$

Also for $i=3, \ldots, n-2$, from lemma 6.3.4,

$$
g_{2}^{i+1}=-\frac{l_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} f_{2}^{i-1}
$$

So,

$$
\begin{aligned}
f_{1}^{i} & =u_{i+2}\left(\frac{l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}+\frac{g_{2}^{i}}{\hat{u}_{i}}\right)+\left(l_{i+2}-\hat{u}_{i}\right) \frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}} \\
& =u_{i+2}\left(l_{i+1}-g_{1}^{i+1}\right)+l_{i+2} \frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}-\frac{f_{2}^{i-1}}{\hat{l}_{i-1}} \\
& =u_{i+1}\left(l_{i}-g_{1}^{i}\right)-g_{2}^{i+1}-\frac{f_{2}^{i-1}}{\hat{l}_{i-1}}
\end{aligned}
$$

The cases $i=2$ and $i=n-2$ also derive from lemmas 6.3.2 and 6.3.3. The case $i=n-1$ follows directly from step $n-1$.

Lemma 6.3.6 Consider $G^{(i)}$ at the beginning of minor step $i, i=1, \ldots, n$, of tridqds algorithm. Then quantities $d_{i}$ satisfy

$$
d_{1}=1, \quad d_{2}=1-\left(\frac{g_{1}^{1}}{\hat{u}_{1}}+\frac{f_{1}^{0}}{\hat{u}_{1}}\right)
$$

and

$$
d_{i}=1-\left(\frac{g_{1}^{i-1}}{\hat{u}_{i-1}}+\frac{f_{1}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}\right), \quad i=3, \ldots, n
$$

Proof. From minor step 1 the result for $d_{2}$ is straightforward,

$$
d_{2}=1-\left[\left(g_{1}^{1}+f_{1}^{0}\right) / \hat{u}_{1}\right]
$$

In minor step $i, i=2, \ldots, n-1$, we obtain $d_{i+1}$ through the sequence

$$
\begin{gathered}
f_{1}^{i} \longleftarrow-f_{1}^{i-1} / \hat{l}_{i-1} \\
d_{i+1} \longleftarrow g_{1}^{i}-f_{1}^{i} \\
d_{i+1} \longleftarrow d_{i+1} / \hat{u}_{i} \\
d_{i+1} \longleftarrow 1-d_{i+1} .
\end{gathered}
$$

Then,

$$
d_{i+1}=1-\left\{\left[g_{1}^{i}-\left(-f_{1}^{i-1} / \hat{l}_{i-1}\right)\right] / \hat{u}_{i}\right\}=1-\left(\frac{g_{1}^{i}}{\hat{u}_{i}}+\frac{f_{1}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}\right) .
$$

For $i=3, \ldots, n$ we have

$$
d_{i}=1-\left(\frac{g_{1}^{i-1}}{\hat{u}_{i-1}}+\frac{f_{1}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}\right) .
$$

Lemma 6.3.7 Consider $\hat{u}_{i}, i=1, \ldots, n$, obtained when tridqds algorithm is applied to $L$ and $U$. Then

$$
\begin{aligned}
& \hat{u}_{1}=d_{1} u_{1}+g_{1}^{1}+f_{1}^{0}, \\
& \hat{u}_{i}=d_{i} u_{i}+g_{1}^{i}+\frac{f_{1}^{i-1}}{\hat{l}_{i-1}}, \quad i=2, \ldots, n-1, \\
& \hat{u}_{n}=d_{n} u_{n} .
\end{aligned}
$$

Proof. For $i=1$ we have

$$
\hat{u}_{1}=\left(d_{1} * u_{1}+g_{1}^{1}\right)+f_{1}^{0}=d_{1} u_{1}+g_{1}^{1}+f_{1}^{0} .
$$

To obtain $\hat{u}_{i}, i=2, \ldots, n-1$, we need to compute

$$
\begin{array}{r}
d_{i+1} \longleftarrow d_{i} * u_{i}+g_{1}^{i} \\
f_{1}^{i} \longleftarrow-f_{1}^{i-1} / \hat{l}_{i-1} \\
\hat{u}_{i} \longleftarrow d_{i+1}-f_{1}^{i} .
\end{array}
$$

Thus,

$$
\hat{u}_{i}=\left(d_{i} * u_{i}+g_{1}^{i}\right)-\left(-f_{1}^{i-1} / \hat{l}_{i-1}\right)=d_{i} u_{i}+g_{1}^{i}+\frac{f_{1}^{i-1}}{\hat{l}_{i-1}} .
$$

In the end

$$
\hat{u}_{n}=d_{n} * u_{n} .
$$

The result of the next theorem is quite surprising. It gives for the quantities $d_{i}$ generated by the triple dqds a relation to the quantities $u_{i}$ and $\hat{u}_{i}$ similar to the one that exists in dqd. Recall from (5.13), page 161, that in dqd algorithm we have

$$
\begin{aligned}
d_{1} & =u_{1} \\
d_{i+1} & =d_{i}\left(\frac{u_{i+1}}{\widehat{u}_{i}}\right), \quad i=1, \ldots, n-1,
\end{aligned}
$$

and for tridqds
Theorem 6.3.1 For quantities $d_{i}, i=1, \ldots, n$, generated by tridqds algorithm applied to $L$ and $U$ we have $d_{1}=1$ and

$$
\begin{aligned}
d_{i+1} & =d_{i}\left(\frac{u_{i}}{\hat{u}_{i}}\right), \quad i=1, \ldots, n-2 \\
d_{n} & =\frac{u_{n}}{\hat{u}_{n}} .
\end{aligned}
$$

Proof. From lemma 6.3.6 and lemma 6.3.7, for $i=3, \ldots, n-1$, we have

$$
d_{i}=1-\frac{\hat{l}_{i-2} g_{1}^{i-1}+f_{1}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}} \quad \text { and } \quad \hat{u_{i}}=d_{i} u_{i}+\frac{\hat{l}_{i-1} g_{1}^{i}+f_{1}^{i-1}}{\hat{l}_{i-1}} .
$$

Thus

$$
\hat{u}_{i}=d_{i} u_{i}+\hat{u}_{i} \frac{\hat{l}_{i-1} g_{1}^{i}+f_{1}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}=d_{i} u_{i}+\hat{u}_{i}\left(1-d_{i+1}\right), \quad i=2, \ldots, n-2 .
$$

Solving this equation yields

$$
\hat{u}_{i}=d_{i} u_{i}+\hat{u}_{i}-\hat{u}_{i} d_{i+1} \Leftrightarrow d_{i+1}=d_{i}\left(\frac{u_{i}}{\hat{u}_{i}}\right) .
$$

For $i=1$ this result also holds. Using the same lemmas, we have

$$
d_{2}=1-\frac{l_{1}+f_{1}^{0}}{\hat{u}_{1}} \quad \text { and } \quad \hat{u_{1}}=u_{1}+l_{1}+f_{1}^{0}
$$

since $d_{1}=1$ and $g_{1}^{1}=l_{1}$. Then

$$
d_{2}=\frac{\hat{u}_{1}-\left(\hat{u}_{1}-u_{1}\right)}{\hat{u}_{1}}=\frac{u_{1}}{\hat{u}_{1}}=d_{1}\left(\frac{u_{1}}{\hat{u}_{1}}\right) .
$$

For $d_{n}$ the result is immediate.
Next lemma show us the expressions for $\hat{l}_{i}, i=1, \ldots, n-1$.

Lemma 6.3.8 Consider $\hat{l}_{i}, i=1, \ldots, n-1$, obtained when tridqds algorithm is applied to $L$ and $U$. Then

$$
\begin{aligned}
\hat{l}_{1} & =\frac{l_{2}+u_{2}-\hat{u}_{1}}{\hat{u}_{1}} f_{1}^{0}+\frac{u_{2}}{\hat{u}_{1}} g_{1}^{1}+\frac{g_{2}^{1}+f_{2}^{0}}{\hat{u}_{1}}, \\
\hat{l}_{i} & =\frac{u_{i+1}+l_{i+1}-\hat{u}_{i}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{u_{i+1}}{\hat{u}_{i}} g_{1}^{i}+\frac{\hat{l}_{i-1}-u_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} g_{2}^{i}, \quad i=2, \ldots, n-2, \\
\hat{l}_{n-1} & =\frac{u_{n}-\hat{u}_{n-1}}{\hat{l}_{n-2} \hat{u}_{n-1}} f_{1}^{n-2}+\frac{u_{n}}{\hat{u}_{n-1}} g_{1}^{n-1} .
\end{aligned}
$$

Proof. For $\hat{l}_{1}$ we have

$$
\begin{aligned}
\hat{l}_{1} & =-f_{1}^{0}+\left[\left(g_{2}^{1}+f_{2}^{0}+f_{1}^{0} * l_{2}\right) / \hat{u}_{1}\right]+\left[\left(g_{1}^{1}+f_{1}^{0}\right) / \hat{u}_{1}\right] * u_{2} \\
& =\frac{l_{2}+u_{2}-\hat{u}_{1}}{\hat{u}_{1}} f_{1}^{0}+\frac{u_{2}}{\hat{u}_{1}} g_{1}^{1}+\frac{g_{2}^{1}+f_{2}^{0}}{\hat{u}_{1}} .
\end{aligned}
$$

To calculate $\hat{l}_{i}, i=2, \ldots, n-2$ we need to follow

$$
\begin{aligned}
& f_{1}^{i} \longleftarrow-f_{1}^{i-1} / \hat{l}_{i-1} \\
& f_{2}^{i} \longleftarrow-f_{2}^{i-1} / \hat{l}_{i-1} \\
& d_{i+1} \longleftarrow g_{1}^{i}-f_{1}^{i} \\
& g_{1}^{i+1} \longleftarrow g_{2}^{i}-f_{2}^{i}-f_{1}^{i} * l_{i+1} \\
& d_{i+1} \longleftarrow d_{i+1} / \hat{u}_{i} \\
& g_{1}^{i+1} \longleftarrow g_{1}^{i+1} / \hat{u}_{i} \\
& \hat{l}_{i} \longleftarrow f_{1}^{i}+g_{1}^{i+1}+d_{i+1} * u_{i+1} .
\end{aligned}
$$

This way,

$$
\begin{aligned}
\hat{l}_{i}= & \left(-f_{1}^{i-1} / \hat{l}_{i-1}\right)+\left\{\left[g_{2}^{i}-\left(-f_{2}^{i-1} / \hat{l}_{i-1}\right)-\left(-f_{1}^{i-1} / \hat{l}_{i-1}\right) * l_{i+1}\right] / \hat{u}_{i}\right\}+ \\
& +\left\{\left[g_{1}^{i}-\left(-f_{1}^{i-1} / \hat{l}_{i-1}\right)\right] / \hat{u}_{i}\right\} * u_{i+1} \\
= & -\frac{\hat{u}_{i} f_{1}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}+\frac{\hat{l}_{i-1} g_{2}^{i}+f_{2}^{i-1}+l_{i+1} f_{1}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}+\frac{\hat{l}_{i-1} g_{1}^{i}+f_{1}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}} u_{i+1} \\
= & \frac{u_{i+1}+l_{i+1}-\hat{u}_{i}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{u_{i+1}}{\hat{u}_{i}} g_{1}^{i}+\frac{f_{2}^{i-1}+\hat{l}_{i-1} g_{2}^{i}}{\hat{l}_{i-1} \hat{u}_{i}} .
\end{aligned}
$$

From lemma 6.3 .4 we know that, for $i=2, \ldots, n-2$,

$$
f_{2}^{i-1}=-u_{i+2} g_{2}^{i} .
$$

So,

$$
\hat{l}_{i}=\frac{u_{i+1}+l_{i+1}-\hat{u}_{i}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{u_{i+1}}{\hat{u}_{i}} g_{1}^{i}+\frac{\hat{l}_{i-1}-u_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} g_{2}^{i} .
$$

To compute $\hat{l}_{n-1}$ we have

$$
\begin{aligned}
& f_{1}^{n-1} \longleftarrow-f_{1}^{n-2} / \hat{l}_{n-2} \\
& d_{n} \longleftarrow g_{1}^{n-1}-f_{1}^{n-1} \\
& d_{n} \longleftarrow d_{n} / \hat{u}_{n-1} \\
& \hat{l}_{n-1} \longleftarrow f_{1}^{n-1}+d_{n} * u_{n} .
\end{aligned}
$$

Thus,

$$
\begin{aligned}
\hat{l}_{n-1} & =\left(-f_{1}^{n-2} / \hat{l}_{n-2}\right)+\left\{\left[g_{1}^{n-1}-\left(-f_{1}^{n-2} / \hat{l}_{n-2}\right)\right] / \hat{u}_{n-1}\right\} * u_{n} \\
& =\frac{u_{n}-\hat{u}_{n-1}}{\hat{l}_{n-2} \hat{u}_{n-1}} f_{1}^{n-2}+\frac{u_{n}}{\hat{u}_{n-1}} g_{1}^{n-1} .
\end{aligned}
$$

But these expressions for $\hat{l}_{i}, i=1, \ldots, n-1$, can still be simplified.
Lemma 6.3.9 Consider $\hat{l}_{i}, i=1, \ldots, n-1$, obtained when tridqds algorithm is applied to $L$ and $U$. Then

$$
\begin{aligned}
\hat{l}_{1} & =u_{2}\left(1-d_{2}\right)+\left(l_{2}-g_{1}^{2}\right)-f_{1}^{0}, \\
\hat{l}_{i} & =u_{i+1}\left(1-d_{i+1}\right)+\left(l_{i+1}-g_{1}^{i+1}\right)-\frac{f_{1}^{i-1}}{\hat{l}_{i-1}}, \quad i=2, \ldots, n-2, \\
\hat{l}_{n-1} & =u_{n}\left(1-d_{n}\right)-\frac{f_{1}^{n-2}}{\hat{l}_{n-2}} .
\end{aligned}
$$

Proof. By previous lemma, for $i=2, \ldots, n-2$,

$$
\begin{align*}
\hat{l}_{i} & =\frac{u_{i+1}+l_{i+1}-\hat{u}_{i}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{u_{i+1}}{\hat{u}_{i}} g_{1}^{i}+\frac{\hat{l}_{i-1}-u_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} g_{2}^{i} \\
& =\frac{u_{i+1}}{\hat{u}_{i}}\left(\frac{f_{1}^{i-1}}{\hat{l}_{i-1}}+g_{1}^{i}\right)+\frac{l_{i+1}-\hat{u}_{i}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{\hat{l}_{i-1}-u_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} g_{2}^{i} . \tag{6.22}
\end{align*}
$$

From lemma 6.3.7 and theorem 6.3.1, we can write

$$
\begin{equation*}
\frac{u_{i+1}}{\hat{u}_{i}}\left(\hat{u}_{i}-d_{i} u_{i}\right)=u_{i+1}\left(1-d_{i} \frac{u_{i}}{\hat{u}_{i}}\right)=u_{i+1}\left(1-d_{i+1}\right) . \tag{6.23}
\end{equation*}
$$

Now the last term in (6.22). We have

$$
\frac{\hat{l}_{i-1}-u_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} g_{2}^{i}=\frac{g_{2}^{i}}{\hat{u}_{i}}+\frac{-u_{i+2} g_{2}^{i}}{\hat{l}_{i-1} \hat{u}_{i}}=\frac{g_{2}^{i}}{\hat{u}_{i}}+\frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}},
$$

since, from lemma 6.3.4, it holds

$$
f_{2}^{i-1}=-u_{i+2} g_{2}^{i} .
$$

Then joining the last two terms in (6.22) gives

$$
\begin{align*}
\frac{l_{i+1}-\hat{u}_{i}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{\hat{l}_{i-1}-u_{i+2}}{\hat{l}_{i-1} \hat{u}_{i}} g_{2}^{i} & =\frac{l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}-\frac{f_{1}^{i-1}}{\hat{l}_{i-1}}+\frac{g_{2}^{i}}{\hat{u}_{i}}+\frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}} \\
& =\left(\frac{l_{i+1}}{\hat{l}_{i-1} \hat{u}_{i}} f_{1}^{i-1}+\frac{f_{2}^{i-1}}{\hat{l}_{i-1} \hat{u}_{i}}+\frac{g_{2}^{i}}{\hat{u}_{i}}\right)-\frac{f_{1}^{i-1}}{\hat{l}_{i-1}} \\
& =\left(l_{i+1}-g_{1}^{i+1}\right)-\frac{f_{1}^{i-1}}{\hat{l}_{i-1}}, \tag{6.24}
\end{align*}
$$

since, from lemma 6.3.3,

$$
g_{1}^{i}=l_{i}-\left(\frac{l_{i}}{\hat{l}_{i-2} \hat{u}_{i-1}} f_{1}^{i-2}+\frac{f_{2}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}+\frac{g_{2}^{i-1}}{\hat{u}_{i-1}}\right) .
$$

Finally, from (6.22), (6.23) and (6.24), we obtain

$$
\hat{l}_{i}=u_{i+1}\left(1-d_{i+1}\right)+\left(l_{i+1}-g_{1}^{i+1}\right)-\frac{f_{1}^{i-1}}{\hat{l}_{i-1}} .
$$

For $i=n-1$,

$$
\begin{aligned}
\hat{l}_{n-1} & =\frac{u_{n}-\hat{u}_{n-1}}{\hat{l}_{n-2} \hat{u}_{n-1}^{n-2}}+\frac{u_{n}}{\hat{u}_{n-1}} g_{1}^{n-1} \\
& =\frac{u_{n}}{\hat{u}_{n-1}}\left(\frac{f_{1}^{n-2}}{\hat{l}_{n-2}}+g_{1}^{n-1}\right)-\frac{f_{1}^{n-2}}{\hat{l}_{n-2}} \\
& =\frac{u_{n}}{\hat{u}_{n-1}}\left(\hat{u}_{n-1}-d_{n-1} u_{n-1}\right)-\frac{f_{1}^{n-2}}{\hat{l}_{n-2}} \\
& =u_{n}\left(1-d_{n}\right)-\frac{f_{1}^{n-2}}{\hat{l}_{n-2}} .
\end{aligned}
$$

The case $i=1$ :

$$
\begin{aligned}
\hat{l}_{1} & =\frac{l_{2}+u_{2}-\hat{u}_{1}}{\hat{u}_{1}} f_{1}^{0}+\frac{u_{2}}{\hat{u}_{1}} g_{1}^{1}+\frac{g_{2}^{1}+f_{2}^{0}}{\hat{u}_{1}} \\
& =\frac{u_{2}}{\hat{u}_{1}}\left(f_{1}^{0}+g_{1}^{1}\right)+\frac{l_{2}-\hat{u}_{1}}{\hat{u}_{1}} f_{1}^{0}+\frac{g_{2}^{1}+f_{2}^{0}}{\hat{u}_{1}} \\
& =\frac{u_{2}}{\hat{u}_{1}}\left(\hat{u}_{1}-d_{1} u_{1}\right)+\left(\frac{l_{2}}{\hat{u}_{1}} f_{1}^{0}+\frac{g_{2}^{1}}{\hat{u}_{1}}+\frac{f_{2}^{0}}{\hat{u}_{1}}\right)-f_{1}^{0} \\
& =u_{2}\left(1-d_{2}\right)+\left(l_{2}-g_{1}^{2}\right)-f_{1}^{0}
\end{aligned}
$$

using lemma 6.3.7, theorem 6.3.1 and lemma 6.3.3.

Combining all these results we obtain the new version of tridqds algorithm that we will call 3dqds. We introduce a new temporary variable $a u x$.

The changes to the first minor step:

## Minor step 1 :

$$
\begin{aligned}
& d_{1}=1 ; g_{1}^{1}=l_{1} ; g_{2}^{1}=0 \\
& f_{1}^{0}=\left(u_{1}+l_{1}\right)^{2}+u_{2} l_{1}-2(\Re \sigma)\left(u_{1}+l_{1}\right)+|\sigma|^{2} \\
& f_{2}^{0}=u_{2} l_{1} u_{3} l_{2} / f_{1}^{0} \\
& f_{1}^{0}=u_{2} l_{1}\left(u_{1}+l_{1}+u_{2}+l_{2}-2(\Re \sigma)\right) / f_{1}^{0} \\
& \hat{u}_{1}=d_{1} u_{1}+g_{1}^{1}+f_{1}^{0} \\
& d_{2}=d_{1} u_{1} / \hat{u}_{1} \\
& a u x=\left(l_{2} f_{1}^{0}+f_{2}^{0}+g_{2}^{1}\right) / \hat{u}_{1} \\
& g_{1}^{2}=l_{2}-a u x \\
& \hat{l}_{1}=u_{2}\left(1-d_{2}\right)+a u x-f_{1}^{0} \\
& g_{2}^{2}=-l_{3} f_{2}^{0} / \hat{u}_{1} \\
& f_{1}^{1}=u_{3} \cdot a u x-g_{2}^{2}-f_{2}^{0} \\
& f_{2}^{1}=-u_{4} g_{2}^{1}
\end{aligned}
$$

The inner loop is changed to:

$$
\begin{aligned}
& \text { for } i=2, \ldots, n-3 \\
& \qquad \begin{array}{l}
\hat{u}_{i}=d_{i} u_{i}+g_{1}^{i}+f_{1}^{i-1} / \hat{l}_{i-1} \\
d_{i+1}=d_{i} u_{i} / \hat{u}_{i} \\
a u x=\left(l_{i+1} f_{1}^{i-1}+f_{2}^{i-1}+\hat{l}_{i-1} g_{2}^{i}\right) /\left(\hat{l}_{i-1} \hat{u}_{i}\right) \\
g_{1}^{i+1}=l_{i+1}-a u x \\
\hat{l}_{i}=u_{i+1}\left(1-d_{i+1}\right)+a u x-f_{1}^{i-1} / \hat{l}_{i-1} \\
g_{2}^{i+1}=-l_{i+2} f_{2}^{i-1} /\left(\hat{l}_{i-1} \hat{u}_{i}\right) \\
f_{1}^{i}=u_{i+2} \cdot a u x-g_{2}^{i+1}-f_{2}^{i-1} / \hat{l}_{i-1} \\
f_{2}^{i}=-u_{i+3} g_{2}^{i+1}
\end{array}
\end{aligned}
$$

## end for

The last three steps become:
Minor step n-2 :

$$
\begin{aligned}
& \hat{u}_{n-2}=d_{n-2} u_{n-2}+g_{1}^{n-2}+f_{1}^{n-3} / \hat{l}_{n-3} \\
& d_{n-1}=d_{n-2} u_{n-2} / \hat{u}_{n-2} \\
& a u x=\left(l_{n-1} f_{1}^{n-3}+f_{2}^{n-3}+\hat{l}_{n-3} g_{2}^{n-2}\right) /\left(\hat{l}_{n-3} \hat{u}_{n-2}\right) \\
& g_{1}^{n-1}=l_{n-1}-a u x \\
& \hat{l}_{n-2}=u_{n-1}\left(1-d_{n-1}\right)+a u x-f_{1}^{n-3} / \hat{l}_{n-3} \\
& f_{1}^{n-2}=u_{n} \cdot a u x-f_{2}^{n-3} / \hat{l}_{n-3}
\end{aligned}
$$

Minor step n-1 :

$$
\begin{aligned}
& \hat{u}_{n-1}=d_{n-1} u_{n-1}+g_{1}^{n-1}+f_{1}^{n-2} / \hat{l}_{n-2} \\
& d_{n}=1-\left(g_{1}^{n-1}+f_{1}^{n-2} / \hat{l}_{n-2}\right) / \hat{u}_{n-1} \\
& \hat{l}_{n-1}=u_{n}\left(1-d_{n}\right)-f_{1}^{n-2} / \hat{l}_{n-2}
\end{aligned}
$$

Minor step n:

$$
\hat{u}_{n}=d_{n} * u_{n}
$$

Now, turning into a computacional implementation we will go back again to variables $x_{l}$ and $y_{l}$ to play the role of $f_{1}^{i}$ and $f_{2}^{i}$, respectively, and to variables $y_{r}$ and $z_{r}$ to replace $g_{1}^{i}$ and $g_{2}^{i}$, respectively. Variable $d_{i}$ will substitute $x_{r}$. Recognizing common subexpressions, we will introduce some changes with the goal of reducing the number of divisions needed.

If we let

$$
\bar{l}=\frac{1}{\hat{l}_{i-1}} \quad \text { and } \quad \bar{u}=\frac{1}{\hat{u}_{i}}
$$

and initially

$$
\begin{aligned}
x_{l} & =f_{1}^{i-1} & & y_{r}=g_{1}^{i} \\
y_{l} & =f_{2}^{i-1} & & z_{r}=g_{2}^{i}
\end{aligned}
$$

then the inner loop will be

$$
\begin{aligned}
& \text { for } i=2, \ldots, n-3 \\
& \qquad \begin{array}{l}
x_{l}=x_{l} * \bar{l} \\
y_{l}=y_{l} * \bar{l} \\
\hat{u_{i}}=d_{i} * u_{i}+y_{r}+x_{l} \\
\bar{u}=1 / \hat{u}_{i} \\
d_{i+1}=d_{i} * u_{i} * \bar{u} \\
a u x=\left(l_{i+1} * x_{l}+y_{l}+z_{r}\right) * \bar{u} \\
y_{r}=l_{i+1}-a u x \\
\hat{l}_{i}=u_{i+1}-d_{i+1} * u_{i+1}+a u x-x_{l} \\
z_{r}=-l_{i+2} * y_{l} * \bar{u} \\
x_{l}=u_{i+2} * a u x-z_{r}-y_{l} \\
y_{l}=-u_{i+3} * z_{r} \\
\bar{l}=1 / \hat{l}_{i}
\end{array} \\
& \text { end for }
\end{aligned}
$$

So, we incorporated two additional auxiliary variables but we reduced the number of divisions to only 2 .

It is still possible to reduce the number or multiplications by noticing that the multiplication $d_{i} * u_{i}$ can be performed implicitly. If we define $t_{i}$ by

$$
t_{i} \equiv d_{i} * u_{i}
$$

then

$$
t_{i+1}=d_{i+1} * u_{i+1}=d_{i} * u_{i} * \bar{u} * u_{i+1}=t_{i} * \bar{u} * u_{i+1}
$$

Thus, $\hat{u}_{i}$ and $\hat{l}_{i}$ will be given by

$$
\begin{aligned}
\hat{u}_{i} & =t_{i}+y_{r}+x_{l} \\
t_{i+1} & =t_{i} * \bar{u} * u_{i+1} \\
\hat{l}_{i} & =u_{i+1}-t_{i+1}+a u x-x_{l}
\end{aligned}
$$

It can be verified that there is no need for both variables $\bar{l}$ and $\bar{u}$. The variable $a u x_{1}$ will play the role of both. The final version of 3dqds algorithm is presented below.

3dqds( $\sigma$ ):

$$
\begin{aligned}
& t=1 ; y_{r}=l_{1} ; z_{r}=0 \\
& x_{l}=\left(u_{1}+l_{1}\right)^{2}+u_{2} * l_{1}-2(\Re \sigma) *\left(u_{1}+l_{1}\right)+|\sigma|^{2} \\
& y_{l}=u_{2} * l_{1} * u_{3} * l_{2} / x_{l} \\
& x_{l}=u_{2} * l_{1} *\left(u_{1}+l_{1}+u_{2}+l_{2}-2(\Re \sigma)\right) / x_{l} \\
& t=t * u_{1} \\
& \hat{u}_{1}=t+y_{r}+x_{l} \\
& a u x_{1}=1 / \hat{u}_{1} \\
& a u x=\left(l_{2} * x_{l}+y_{l}+z_{r}\right) * a u x_{1} \\
& y_{r}=l_{2}-a u x \\
& t=t * u_{2} * a u x_{1} \\
& \hat{l}_{1}=u_{2}-t+a u x-x_{l} \\
& z_{r}=-l_{3} * y_{l} * a u x_{1} \\
& x_{l}=u_{3} * a u x-z_{r}-y_{l}
\end{aligned}
$$

$$
\begin{aligned}
& y_{l}=-u_{4} * z_{r} \\
& \text { aux } x_{1}=1 / \hat{l}_{1} \\
& \text { for } i=2, \ldots, n-3 \\
& \qquad \begin{array}{l}
x_{l}=x_{l} * a u x_{1} \\
y_{l}=y_{l} * a u x_{1} \\
\hat{u_{i}}=t+y_{r}+x_{l} \\
\quad a u x_{1}=1 / \hat{u}_{i} \\
\quad a u x=\left(l_{i+1} * x_{l}+y_{l}+z_{r}\right) * a u x_{1} \\
y_{r}=l_{i+1}-a u x \\
t=u_{i+1} * t * a u x_{1} \\
\hat{l}_{i}=u_{i+1}-t+a u x-x_{l} \\
z_{r}=-l_{i+2} * y_{l} * a u x_{1} \\
x_{l}=u_{i+2} * a u x-z_{r}-y_{l} \\
y_{l}=-u_{i+3} * z_{r} \\
a u x_{1}=1 / \hat{l}_{i}
\end{array}
\end{aligned}
$$

## end for

$x_{l}=x_{l} *$ aux $_{1}$
$y_{l}=y_{l} * a u x_{1}$
$\hat{u}_{n-2}=t+y_{r}+x_{l}$
$a u x_{1}=1 / \hat{u}_{n-2}$
$a u x=\left(l_{n-1} * x_{l}+y_{l}+z_{r}\right) * a u x_{1}$
$y_{r}=l_{n-1}-a u x$
$t=u_{n-1} * t * a u x_{1}$
$\hat{l}_{n-2}=u_{n-1}-t+a u x-x_{l}$
$x_{l}=u_{n} * a u x-y_{l}$
$a u x_{1}=1 / \hat{l}_{n-2}$

$$
\begin{aligned}
& x_{l}=x_{l} * a u x_{1} \\
& \hat{u}_{n-1}=t+y_{r}+x_{l} \\
& a u x_{1}=\left(y_{r}+x_{l}\right) / \hat{u}_{n-1} \\
& \hat{l}_{n-1}=u_{n} * a u x_{1}-x_{l} \\
& \hat{u}_{n}=\left(1-a u x_{1}\right) * u_{n}
\end{aligned}
$$

### 6.3.3 Operation count for 3dqds

Table 6.1 below shows that the more elegant version 3dqds of triple dqds algorithm has the advantage of performing less three divisions than the version tridqds. The price to pay for this reduction on the number of divisions is worthwhile considering that in certain machines the cost of a division is much more expensive than the cost of a multiplication.

|  | tridqds | 3dqds |
| :---: | :---: | :---: |
| Divisions | 5 | 2 |
| Multiplications | 6 | 10 |
| Additions | 5 | 5 |
| Subtractions | 6 | 5 |
| Assignments | 16 | 12 |
| Auxiliary variables | 5 | 7 |

Figure 6.1: Operation Count of tridqds and 3dqds

The simplicity of 3dqds is very attractive and preliminary numerical tests reveal its robustness.

### 6.3.4 Entries of $(U L)^{-1}$

In this section we will present results for the entries of $(U L)^{-1}$ analogous to the results shown for dqd in theorem 6.2.4 (page 176). So, the goal is to obtain expressions for

$$
\left[(U L)^{-1}\right]_{i: i+1, i: i+1}=\left[\begin{array}{cc}
{\left[(U L)^{-1}\right]_{i, i}} & {\left[(U L)^{-1}\right]_{i, i+1}} \\
{\left[(U L)^{-1}\right]_{i+1, i}} & {\left[(U L)^{-1}\right]_{i+1, i+1}}
\end{array}\right], \quad i=1, \ldots, n-1
$$

when we apply triple dqds.
The first lemma gives the expression for entries $(i-2, i),(i-1, i)$ and $(i, i)$ of $\left(F^{(i)} G^{(i)}\right)^{-1}$ that will be used to obtain the diagonal of $(U L)^{-1}$. See page 179 for recalling the properties of the matrices $F^{(i)}$ and $G^{(i)}$.

Lemma 6.3.10 Consider $F^{(i)}$ and $G^{(i)}$ at the beginning of minor step $i, i=1, \ldots, n$, when we apply tridqds algorithm to $L$ and $U$. If $U$ is invertible then

$$
\begin{align*}
& {\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i i}=\frac{1}{d_{i} u_{i}}, \quad i=1, \ldots, n}  \tag{6.25}\\
& {\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-1, i}=-\frac{1}{d_{i} u_{i} \hat{u}_{i-1}}, \quad i=2, \ldots, n}  \tag{6.26}\\
& {\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-2, i}=\frac{1}{d_{i} u_{i} \hat{u}_{i-1} \hat{u}_{i-2}}, \quad i=3, \ldots, n} \tag{6.27}
\end{align*}
$$

Proof. We start by using the interesting fact that row $i$ of $G^{(i)}$ and column $i$ of $F^{(i)}$ are singletons (see (6.17)). We have, for $i=1, \ldots, n$,

$$
F^{(i)} \boldsymbol{e}_{i}=u_{i} \boldsymbol{e}_{i} \quad \text { and } \quad \boldsymbol{e}_{i}^{T} G^{(i)}=d_{i} \boldsymbol{e}_{i}^{T}
$$

and, equivalently,

$$
u_{i}^{-1} \boldsymbol{e}_{i}=\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i} \quad \text { and } \quad d_{i}^{-1} \boldsymbol{e}_{i}^{T}=\boldsymbol{e}_{i}^{T}\left(G^{(i)}\right)^{-1}
$$

Therefore,

$$
\begin{aligned}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i i} } & =\boldsymbol{e}_{i}^{T}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right] \boldsymbol{e}_{i} \\
& =\left[\boldsymbol{e}_{i}^{T}\left(G^{(i)}\right)^{-1}\right]\left[\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right] \\
& =d_{i}^{-1} \boldsymbol{e}_{i}^{T} u_{i}^{-1} \boldsymbol{e}_{i} \\
& =\frac{1}{d_{i} u_{i}}
\end{aligned}
$$

Analogously, for $i=2, \ldots, n$, we can write

$$
\begin{align*}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-1, i} } & =\boldsymbol{e}_{i-1}^{T}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right] \boldsymbol{e}_{i} \\
& =\left[\boldsymbol{e}_{i-1}^{T}\left(G^{(i)}\right)^{-1}\right]\left[\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right] \\
& =\left[\boldsymbol{e}_{i-1}^{T}\left(G^{(i)}\right)^{-1}\right] u_{i}^{-1} \boldsymbol{e}_{i} . \tag{6.28}
\end{align*}
$$

But now row $(i-1)$ of $G^{(i)}$ is not a singleton. We have

$$
\boldsymbol{e}_{i-1}^{T}\left(G^{(i)}\right)=\hat{u}_{i-1} \boldsymbol{e}_{i-1}^{T}+\boldsymbol{e}_{i}^{T}
$$

Then

$$
\begin{aligned}
\boldsymbol{e}_{i-1}^{T} & =\left(\hat{u}_{i-1} \boldsymbol{e}_{i-1}^{T}+\boldsymbol{e}_{i}^{T}\right)\left(G^{(i)}\right)^{-1} \\
& =\hat{u}_{i-1} \boldsymbol{e}_{i-1}^{T}\left(G^{(i)}\right)^{-1}+\boldsymbol{e}_{i}^{T}\left(G^{(i)}\right)^{-1} \\
& =\hat{u}_{i-1} \boldsymbol{e}_{i-1}^{T}\left(G^{(i)}\right)^{-1}+d_{i}^{-1} \boldsymbol{e}_{i}^{T} .
\end{aligned}
$$

And

$$
\begin{equation*}
\boldsymbol{e}_{i-1}^{T}\left(G^{(i)}\right)^{-1}=\left(\boldsymbol{e}_{i-1}^{T}-d_{i}^{-1} \boldsymbol{e}_{i}^{T}\right) \hat{u}_{i-1}^{-1} \tag{6.29}
\end{equation*}
$$

Thus, from (6.28), we get

$$
\begin{aligned}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-1, i} } & =\left[\left(\boldsymbol{e}_{i-1}^{T}-d_{i}^{-1} \boldsymbol{e}_{i}^{T}\right) \hat{u}_{i-1}^{-1}\right] u_{i}^{-1} \boldsymbol{e}_{i} \\
& =-d_{i}^{-1} \hat{u}_{i-1}^{-1} u_{i}^{-1} \\
& =-\frac{1}{d_{i} u_{i} \hat{u}_{i-1}},
\end{aligned}
$$

as we wanted to show.
Finally, for $i=3, \ldots, n$,

$$
\begin{align*}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-2, i} } & =\boldsymbol{e}_{i-2}^{T}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right] \boldsymbol{e}_{i} \\
& =\left[\boldsymbol{e}_{i-2}^{T}\left(G^{(i)}\right)^{-1}\right]\left[\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right] \\
& =\left[\boldsymbol{e}_{i-2}^{T}\left(G^{(i)}\right)^{-1}\right] u_{i}^{-1} \boldsymbol{e}_{i} . \tag{6.30}
\end{align*}
$$

We have

$$
\begin{equation*}
\boldsymbol{e}_{i-2}^{T}\left(G^{(i)}\right)=\hat{u}_{i-2} \boldsymbol{e}_{i-2}^{T}+\boldsymbol{e}_{i-1}^{T} \tag{6.31}
\end{equation*}
$$

and then

$$
\begin{align*}
\boldsymbol{e}_{i-2}^{T}\left(G^{(i)}\right)^{-1} & =\left(\boldsymbol{e}_{i-2}^{T}-\boldsymbol{e}_{i-1}^{T}\left(G^{(i)}\right)^{-1}\right) \hat{u}_{i-2}^{-1} \\
& =\left[\boldsymbol{e}_{i-2}^{T}-\left(\boldsymbol{e}_{i-1}^{T}-d_{i}^{-1} \boldsymbol{e}_{i}^{T}\right) \hat{u}_{i-1}^{-1}\right] \hat{u}_{i-2}^{-1} \tag{6.32}
\end{align*}
$$

From (6.30) we get

$$
\begin{aligned}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-2, i} } & =\left[\boldsymbol{e}_{i-2}^{T}-\left(\boldsymbol{e}_{i-1}^{T}-d_{i}^{-1} \boldsymbol{e}_{i}^{T}\right) \hat{u}_{i-1}^{-1}\right] \hat{u}_{i-2}^{-1} u_{i}^{-1} \boldsymbol{e}_{i} \\
& =d_{i}^{-1} \hat{u}_{i-1}^{-1} \hat{u}_{i-2}^{-1} u_{i}^{-1} \\
& =\frac{1}{d_{i} u_{i} \hat{u}_{i-1} \hat{u}_{i-2}} .
\end{aligned}
$$

The next theorem gives the expression for the diagonal of $(U L)^{-1}$.
In (6.16) we already saw that

$$
\begin{aligned}
F^{(i)} G^{(i)} & =\left(\mathcal{L}_{i-1}^{-1} \ldots \mathcal{L}_{2}^{-1} \mathcal{L}_{1}^{-1}\right)(U L)\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right) \\
& =\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)^{-1}(U L)\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right) .
\end{aligned}
$$

Then

$$
U L=\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right) F^{(i)} G^{(i)}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)^{-1} .
$$

and

$$
\begin{equation*}
(U L)^{-1}=\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)\left(F^{(i)} G^{(i)}\right)^{-1}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)^{-1} \tag{6.33}
\end{equation*}
$$

Theorem 6.3.2 Consider $L$ and $U$ as described in (6.1). If $U$ is invertible and tridqds algorithm applied to $L$ and $U$ is successful, then the quantities $d_{i}, f_{1}^{i-1}$ and $f_{2}^{i-1}$ generated by the algorithm satisfy

$$
\begin{aligned}
& {\left[(U L)^{-1}\right]_{1,1}=\frac{1}{d_{1} u_{1}}} \\
& {\left[(U L)^{-1}\right]_{2,2}=\frac{1}{d_{2} u_{2}}\left(1-\frac{f_{1}^{0}}{\hat{u}_{1}}\right)} \\
& {\left[(U L)^{-1}\right]_{3,3}=\frac{1}{d_{3} u_{3}}\left(1-\frac{f_{1}^{1}}{\hat{l}_{1} \hat{u}_{2}}+\frac{f_{2}^{0}}{\hat{u}_{2} \hat{u}_{1}}\right)} \\
& {\left[(U L)^{-1}\right]_{i, i}=\frac{1}{d_{i} u_{i}}\left(1-\frac{f_{1}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}+\frac{f_{2}^{i-3}}{\hat{l}_{i-3} \hat{u}_{i-1} \hat{u}_{i-2}}\right), \quad i=4, \ldots, n .}
\end{aligned}
$$

Proof. Using (6.19), (6.20) and (6.33), for $i=4, \ldots, n$,

$$
\left.\begin{array}{rl}
{\left[(U L)^{-1}\right]_{i, i}} & =\boldsymbol{e}_{i}^{T}\left[(U L)^{-1}\right] \boldsymbol{e}_{i} \\
& =\boldsymbol{e}_{i}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)\left(F^{(i)} G^{(i)}\right)^{-1}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)^{-1} \boldsymbol{e}_{i} \\
& =\left[\begin{array}{lllll}
\left.\boldsymbol{e}_{i}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)\right]\left[\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right.
\end{array}\right] \\
& =\left[\begin{array}{llllll}
0 & \ldots & 0 & h_{2}^{i-2} & h_{1}^{i-1} & 1
\end{array} 0\right. \\
0 & \ldots
\end{array}\right]\left[\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right] .
$$

Using lemma 6.3.10 yields

$$
\begin{align*}
{\left[(U L)^{-1}\right]_{i, i} } & =\frac{h_{2}^{i-2}}{d_{i} u_{i} \hat{u}_{i-1} \hat{u}_{i-2}}-\frac{h_{1}^{i-1}}{d_{i} u_{i} \hat{u}_{i-1}}+\frac{1}{d_{i} u_{i}} \\
& =\frac{1}{d_{i} u_{i}}\left(1-\frac{h_{1}^{i-1}}{\hat{u}_{i-1}}+\frac{h_{2}^{i-2}}{\hat{u}_{i-1} \hat{u}_{i-2}}\right) . \tag{6.34}
\end{align*}
$$

Finally, by lemma 6.3.1, we have

$$
h_{1}^{i-1}=f_{1}^{i-2} / \hat{l}_{i-2} \quad \text { and } \quad h_{2}^{i-2}=f_{2}^{i-3} / \hat{l}_{i-3},
$$

and then

$$
\left[(U L)^{-1}\right]_{i, i}=\frac{1}{d_{i} u_{i}}\left(1-\frac{f_{1}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}+\frac{f_{2}^{i-3}}{\hat{l}_{i-3} \hat{u}_{i-1} \hat{u}_{i-2}}\right)
$$

Now let us see the cases $i=1,2,3$. By lemma 6.3.10, we have

$$
\left[(U L)^{-1}\right]_{1,1}=\left[\left(F^{(1)} G^{(1)}\right)^{-1}\right]_{11}=\frac{1}{d_{1} u_{1}}
$$

For $i=2$, using lemmas 6.3.10 and 6.3.1 again,

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{2,2} } & =\boldsymbol{e}_{2}^{T}\left[(U L)^{-1}\right] \boldsymbol{e}_{2} \\
& =\boldsymbol{e}_{2}^{T} \mathcal{L}_{1}\left(F^{(2)} G^{(2)}\right)^{-1} \mathcal{L}_{1}^{-1} \boldsymbol{e}_{2} \\
& =\boldsymbol{e}_{2}^{T} \mathcal{L}_{1}\left[\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{2}\right] \\
& =\left[\begin{array}{llll}
h_{1}^{1} & 1 & 0 & \ldots
\end{array}\right]\left[\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{2}\right] \\
& =\left(h_{1}^{1} \boldsymbol{e}_{1}^{T}+\boldsymbol{e}_{2}^{T}\right)\left[\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{2}\right] \\
& =h_{1}^{1} \boldsymbol{e}_{1}^{T}\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{2}+\boldsymbol{e}_{2}^{T}\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{2} \\
& =h_{1}^{1}\left[\left(F^{(2)} G^{(2)}\right)^{-1}\right]+\left[\left(F_{1,2}^{(2)} G^{(2)}\right)^{-1}\right]_{2,2} \\
& =-\frac{h_{1}^{1}}{d_{2} u_{2} \hat{u}_{1}}+\frac{1}{d_{2} u_{2}} \\
& =\frac{1}{d_{2} u_{2}}\left(1-\frac{f_{1}^{0}}{\hat{u}_{1}}\right) .
\end{aligned}
$$

Finally, for $i=3,(6.34)$ also yields,

$$
\left[(U L)^{-1}\right]_{3,3}=\frac{1}{d_{3} u_{3}}\left(1-\frac{h_{1}^{2}}{\hat{u}_{2}}+\frac{h_{2}^{1}}{\hat{u}_{2} \hat{u}_{1}}\right)
$$

and by lemma 6.3 .1 we have

$$
h_{1}^{2}=f_{1}^{1} / \hat{l}_{1} \quad \text { and } \quad h_{2}^{1}=f_{2}^{0}
$$

So,

$$
\left[(U L)^{-1}\right]_{3,3}=\frac{1}{d_{3} u_{3}}\left(1-\frac{f_{1}^{1}}{\hat{l}_{1} \hat{u}_{2}}+\frac{f_{2}^{0}}{\hat{u}_{2} \hat{u}_{1}}\right) .
$$

Next lemma will drive us to the off-diagonal elements of $(U L)^{-1}$.

Lemma 6.3.11 Consider $F^{(i)}$ and $G^{(i)}$ at the beginning of minor step $i, i=1, \ldots, n$, when we apply tridqds algorithm to $L$ and $U$. If $U$ is invertible then

$$
\begin{aligned}
& {\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i+1, i}=-\frac{g_{1}^{i}}{d_{i} u_{i}}, \quad i=1, \ldots, n-1} \\
& {\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i, i+1}=-\frac{1}{d_{i} u_{i} u_{i+1}}, \quad i=1, \ldots, n-1} \\
& {\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-1, i+1}=\frac{1}{d_{i} u_{i} u_{i+1} \hat{u}_{i-1}}, \quad i=2, \ldots, n-1} \\
& {\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-2, i+1}=-\frac{1}{d_{i} u_{i} u_{i+1} \hat{u}_{i-1} \hat{u}_{i-2}}, \quad i=3, \ldots, n-1}
\end{aligned}
$$

Proof. For $i=1, \ldots, n-1$, we can write

$$
\begin{align*}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i+1, i} } & =\boldsymbol{e}_{i+1}^{T}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right] \boldsymbol{e}_{i} \\
& =\left[\boldsymbol{e}_{i+1}^{T}\left(G^{(i)}\right)^{-1}\right]\left[\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right] \tag{6.35}
\end{align*}
$$

We already know that

$$
u_{i}^{-1} \boldsymbol{e}_{i}=\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i} \quad \text { and } \quad d_{i}^{-1} \boldsymbol{e}_{i}^{T}=\boldsymbol{e}_{i}^{T}\left(G^{(i)}\right)^{-1}
$$

and for row $(i+1)$ of $G^{(i)}$ we have

$$
\boldsymbol{e}_{i+1}^{T}\left(G^{(i)}\right)=g_{1}^{i} \boldsymbol{e}_{i}^{T}+\boldsymbol{e}_{i+1}^{T}
$$

Then

$$
\begin{aligned}
\boldsymbol{e}_{i+1}^{T} & =\left(g_{1}^{i} \boldsymbol{e}_{i}^{T}+\boldsymbol{e}_{i+1}^{T}\right)\left(G^{(i)}\right)^{-1} \\
& =g_{1}^{i} \boldsymbol{e}_{i}^{T}\left(G^{(i)}\right)^{-1}+\boldsymbol{e}_{i+1}^{T}\left(G^{(i)}\right)^{-1} \\
& =g_{1}^{i} d_{i}^{-1} \boldsymbol{e}_{i}^{T}+\boldsymbol{e}_{i+1}^{T}\left(G^{(i)}\right)^{-1}
\end{aligned}
$$

And

$$
\boldsymbol{e}_{i+1}^{T}\left(G^{(i)}\right)^{-1}=\boldsymbol{e}_{i+1}^{T}-g_{1}^{i} d_{i}^{-1} \boldsymbol{e}_{i}^{T}
$$

Thus, from (6.35) we get

$$
\begin{aligned}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i+1, i} } & =\left(\boldsymbol{e}_{i+1}^{T}-g_{1}^{i} d_{i}^{-1} \boldsymbol{e}_{i}^{T}\right) u_{i}^{-1} \boldsymbol{e}_{i} \\
& =-\frac{g_{1}^{i}}{d_{i} u_{i}}
\end{aligned}
$$

For $i=1, \ldots, n-1$, we have

$$
\begin{align*}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i, i+1} } & =\boldsymbol{e}_{i}^{T}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right] \boldsymbol{e}_{i+1} \\
& =\left[\boldsymbol{e}_{i}^{T}\left(G^{(i)}\right)^{-1}\right]\left[\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}\right] \\
& =d_{i}^{-1} \boldsymbol{e}_{i}^{T}\left[\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}\right] \tag{6.36}
\end{align*}
$$

Column $(i+1)$ of $F^{(i)}$ is

$$
\left(F^{(i)}\right) \boldsymbol{e}_{i+1}=\boldsymbol{e}_{i}+u_{i+1} \boldsymbol{e}_{i+1}
$$

Then

$$
\boldsymbol{e}_{i+1}=\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i}+u_{i+1}\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}=u_{i}^{-1} \boldsymbol{e}_{i}+u_{i+1}\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}
$$

which gives

$$
\begin{equation*}
\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}=\left(\boldsymbol{e}_{i+1}-u_{i}^{-1} \boldsymbol{e}_{i}\right) u_{i+1}^{-1} \tag{6.37}
\end{equation*}
$$

Thus, from (6.36), we get

$$
\begin{aligned}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i, i+1} } & =d_{i}^{-1} \boldsymbol{e}_{i}^{T}\left(\boldsymbol{e}_{i+1}-u_{i}^{-1} \boldsymbol{e}_{i}\right) u_{i+1}^{-1} \\
& =-\frac{1}{d_{i} u_{i} u_{i+1}}
\end{aligned}
$$

For $i=2, \ldots, n-1$, using (6.29) and (6.37), we have

$$
\begin{aligned}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-1, i+1} } & =\boldsymbol{e}_{i-1}^{T}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right] e_{i+1} \\
& =\left[\boldsymbol{e}_{i-1}^{T}\left(G^{(i)}\right)^{-1}\right]\left[\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}\right] \\
& =\left(\boldsymbol{e}_{i-1}^{T}-d_{i}^{-1} e_{i}^{T}\right) \hat{u}_{i-1}^{-1}\left(\boldsymbol{e}_{i+1}-u_{i}^{-1} e_{i}\right) u_{i+1}^{-1} \\
& =\frac{1}{d_{i} u_{i} u_{i+1} \hat{u}_{i-1}} .
\end{aligned}
$$

Finally, for $i=3, \ldots, n-1$, using (6.32) and (6.37) gives

$$
\begin{aligned}
{\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-2, i+1} } & =\boldsymbol{e}_{i-2}^{T}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right] \boldsymbol{e}_{i+1} \\
& =\left[\boldsymbol{e}_{i-2}^{T}\left(G^{(i)}\right)^{-1}\right]\left[\left(F^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}\right] \\
& =\left(\left[\boldsymbol{e}_{i-2}^{T}-\left(\boldsymbol{e}_{i-1}^{T}-d_{i}^{-1} \boldsymbol{e}_{i}^{T}\right) \hat{u}_{i-1}^{-1}\right] \hat{u}_{i-2}^{-1}\right)\left(\boldsymbol{e}_{i+1}-u_{i}^{-1} \boldsymbol{e}_{i}\right) u_{i+1}^{-1} \\
& =-\frac{1}{d_{i} u_{i} u_{i+1} \hat{u}_{i-1} \hat{u}_{i-2}} .
\end{aligned}
$$

Next theorem give the entries $(i, i+1)$ of $(U L)^{-1}, i=1, \ldots, n-1$.

Theorem 6.3.3 Consider $L$ and $U$ as described in (6.1). If $U$ is invertible and tridqds algorithm applied to $L$ and $U$ is successful, then the quantities $d_{i}, f_{1}^{i-1}$ and $f_{2}^{i-1}$ generated by the algorithm satisfy

$$
\begin{aligned}
& {\left[(U L)^{-1}\right]_{1,2}=-\frac{1}{d_{1} u_{1} u_{2}},} \\
& {\left[(U L)^{-1}\right]_{2,3}=-\frac{1}{d_{2} u_{2} u_{3}}\left(1-\frac{f_{1}^{0}}{\hat{u}_{1}}\right),} \\
& {\left[(U L)^{-1}\right]_{3,4}=-\frac{1}{d_{3} u_{3} u_{4}}\left(1-\frac{f_{1}^{1}}{\hat{l}_{1} \hat{u}_{2}}+\frac{f_{2}^{0}}{\hat{u}_{2} \hat{u}_{1}}\right),} \\
& {\left[(U L)^{-1}\right]_{i, i+1}=-\frac{1}{d_{i} u_{i} u_{i+1}}\left(1-\frac{f_{1}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}+\frac{f_{2}^{i-3}}{\hat{l}_{i-3} \hat{u}_{i-1} \hat{u}_{i-2}}\right), \quad i=4, \ldots, n-1 .}
\end{aligned}
$$

Proof. Using (6.19), (6.20) and (6.33), for $i=4, \ldots, n$,

$$
\begin{aligned}
& {\left[(U L)^{-1}\right]_{i, i+1}=\boldsymbol{e}_{i}^{T}\left[(U L)^{-1}\right] \boldsymbol{e}_{i+1}} \\
& =\boldsymbol{e}_{i}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)\left(F^{(i)} G^{(i)}\right)^{-1}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)^{-1} \boldsymbol{e}_{i+1} \\
& =\left[\boldsymbol{e}_{i}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)\right]\left[\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}\right] \\
& =\left[\begin{array}{lllllllll}
0 & \ldots & 0 & h_{2}^{i-2} & h_{1}^{i-1} & 1 & 0 & \ldots & 0
\end{array}\right]\left[\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}\right] \\
& =\left(h_{2}^{i-2} \boldsymbol{e}_{i-2}^{T}+h_{1}^{i-1} \boldsymbol{e}_{i-1}^{T}+\boldsymbol{e}_{i}^{T}\right)\left[\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}\right] \\
& =h_{2}^{i-2} \boldsymbol{e}_{i-2}^{T}\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}+h_{1}^{i-1} \boldsymbol{e}_{i-1}^{T}\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i+1}+\boldsymbol{e}_{i}^{T}\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i+1} \\
& =h_{2}^{i-2}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-2, i+1}+h_{1}^{i-1}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-1, i+1}+\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i, i+1} .
\end{aligned}
$$

Using lemma 6.3.11 yields

$$
\begin{align*}
{\left[(U L)^{-1}\right]_{i, i+1} } & =-\frac{h_{2}^{i-2}}{d_{i} u_{i} u_{i+1} \hat{u}_{i-1} \hat{u}_{i-2}}+\frac{h_{1}^{i-1}}{d_{i} u_{i} u_{i+1} \hat{u}_{i-1}}-\frac{1}{d_{i} u_{i} u_{i+1}} \\
& =-\frac{1}{d_{i} u_{i} u_{i+1}}\left(1-\frac{h_{1}^{i-1}}{\hat{u}_{i-1}}+\frac{h_{2}^{i-2}}{\hat{u}_{i-1} \hat{u}_{i-2}}\right) . \tag{6.38}
\end{align*}
$$

Finally, by lemma 6.3.1, we have

$$
h_{1}^{i-1}=f_{1}^{i-2} / \hat{l}_{i-2} \quad \text { and } \quad h_{2}^{i-2}=f_{2}^{i-3} / \hat{l}_{i-3},
$$

and then

$$
\left[(U L)^{-1}\right]_{i, i+1}=-\frac{1}{d_{i} u_{i} u_{i+1}}\left(1-\frac{f_{1}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}+\frac{f_{2}^{i-3}}{\hat{l}_{i-3} \hat{u}_{i-1} \hat{u}_{i-2}}\right) .
$$

Now let us see the cases $i=1,2,3$. Using lemma 6.3.11, we have

$$
\left[(U L)^{-1}\right]_{1,2}=\left[\left(F^{1} G^{(1)}\right)^{-1}\right]_{1,2}=-\frac{1}{d_{1} u_{1} u_{2}} .
$$

For $i=2$, lemmas 6.3.11 and 6.3.1 give

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{2,3} } & =\boldsymbol{e}_{2}^{T}\left[(U L)^{-1}\right] \boldsymbol{e}_{3} \\
& =\boldsymbol{e}_{2}^{T} \mathcal{L}_{1}\left(F^{(2)} G^{(2)}\right)^{-1} \mathcal{L}_{1}^{-1} \boldsymbol{e}_{3} \\
& =\boldsymbol{e}_{2}^{T} \mathcal{L}_{1}\left[\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{3}\right] \\
& =\left[\begin{array}{llll}
h_{1}^{1} & 1 & 0 & \ldots \\
0
\end{array}\right]\left[\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{3}\right] \\
& =\left(h_{1}^{1} \boldsymbol{e}_{1}^{T}+\boldsymbol{e}_{2}^{T}\right)\left[\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{3}\right] \\
& =h_{1}^{1} \boldsymbol{e}_{1}^{T}\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{3}+\boldsymbol{e}_{2}^{T}\left(F^{(2)} G^{(2)}\right)^{-1} \boldsymbol{e}_{3} \\
& =h_{1}^{1}\left[\left(F^{(2)} G^{(2)}\right)^{-1}\right]_{1,3}+\left[\left(F^{(2)} G^{(2)}\right)^{-1}\right]_{2,3} \\
& =\frac{h_{1}^{1}}{d_{2} u_{2} u_{3} \hat{u}_{1}}-\frac{1}{d_{2} u_{2} u_{3}} \\
& =-\frac{1}{d_{2} u_{2} u_{3}}\left(1-\frac{f_{1}^{0}}{\hat{u}_{1}}\right) .
\end{aligned}
$$

Finally, for $i=3$, (6.38) also yields,

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{3,4} } & =-\frac{1}{d_{3} u_{3} u_{4}}\left(1-\frac{h_{1}^{2}}{\hat{u}_{2}}+\frac{h_{2}^{1}}{\hat{u}_{2} \hat{u}_{1}}\right) \\
& =-\frac{1}{d_{3} u_{3} u_{4}}\left(1-\frac{f_{1}^{1}}{\hat{l}_{1} \hat{u}_{2}}+\frac{f_{2}^{0}}{\hat{u}_{2} \hat{u}_{1}}\right),
\end{aligned}
$$

since by lemma 6.3 .1 we have

$$
h_{1}^{2}=f_{1}^{1} / \hat{l}_{1} \quad \text { and } \quad h_{2}^{1}=f_{2}^{0} .
$$

Finally the entries $(i+1, i)$ of $(U L)^{-1}, i=1, \ldots, n-1$.

Theorem 6.3.4 Consider $L$ and $U$ as described in (6.1). If $U$ is invertible and tridqds algorithm applied to $L$ and $U$ is successful, then the quantities $d_{i}, f_{1}^{i-1}, f_{2}^{i-1}$ and $g_{1}^{i}$ generated
by the algorithm satisfy

$$
\begin{aligned}
& {\left[(U L)^{-1}\right]_{2,1}=-\frac{1}{d_{1} u_{1}} g_{1}^{1},} \\
& {\left[(U L)^{-1}\right]_{3,2}=-\frac{1}{d_{2} u_{2}}\left(g_{1}^{2}+\frac{f_{2}^{0}}{\hat{u}_{1}}\right),} \\
& {\left[(U L)^{-1}\right]_{i+1, i}=-\frac{1}{d_{i} u_{i}}\left(g_{1}^{i}+\frac{f_{2}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}\right), \quad i=3, \ldots, n-1 .}
\end{aligned}
$$

Proof. Using (6.19), (6.20) and (6.33), for $i=3, \ldots, n$,

$$
\begin{align*}
& {\left[(U L)^{-1}\right]_{i+1, i}=\boldsymbol{e}_{i+1}^{T}\left[(U L)^{-1}\right] \boldsymbol{e}_{i}} \\
& =\boldsymbol{e}_{i+1}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)\left(F^{(i)} G^{(i)}\right)^{-1}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)^{-1} \boldsymbol{e}_{i}  \tag{6.39}\\
& =\left[\boldsymbol{e}_{i+1}^{T}\left(\mathcal{L}_{1} \mathcal{L}_{2} \ldots \mathcal{L}_{i-1}\right)\right]\left[\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right] \\
& =\left[\begin{array}{lllllllll}
0 & \ldots & 0 & 0 & h_{2}^{i-1} & 0 & 1 & \ldots & 0
\end{array}\right]\left[\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right] \\
& =\left(h_{2}^{i-1} \boldsymbol{e}_{i-1}^{T}+\boldsymbol{e}_{i+1}^{T}\right)\left[\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i}\right] \\
& =h_{2}^{i-1} \boldsymbol{e}_{i-1}^{T}\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i}+\boldsymbol{e}_{i+1}^{T}\left(F^{(i)} G^{(i)}\right)^{-1} \boldsymbol{e}_{i} \\
& =h_{2}^{i-1}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i-1, i}+\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{i+1, i} . \tag{6.40}
\end{align*}
$$

Using lemmas 6.3.10, 6.3.11 and 6.3.1,

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{i+1, i} } & =-\frac{h_{2}^{i-1}}{d_{i} u_{i} \hat{u}_{i-1}}-\frac{g_{1}^{i}}{d_{i} u_{i}} \\
& =-\frac{1}{d_{i} u_{i}}\left(g_{1}^{i}+\frac{h_{2}^{i-1}}{\hat{u}_{i-1}}\right) \\
& =-\frac{1}{d_{i} u_{i}}\left(g_{1}^{i}+\frac{f_{2}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}\right)
\end{aligned}
$$

For $i=2$, (6.40) also holds,

$$
\left[(U L)^{-1}\right]_{3,2}=h_{2}^{1}\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{1,2}+\left[\left(F^{(i)} G^{(i)}\right)^{-1}\right]_{3,2}
$$

Again, using lemmas 6.3.10, 6.3.11 and 6.3.1,

$$
\begin{aligned}
{\left[(U L)^{-1}\right]_{3,2} } & =-\frac{h_{2}^{1}}{d_{2} u_{2} \hat{u}_{1}}-\frac{g_{1}^{2}}{d_{2} u_{2}} \\
& =-\frac{1}{d_{2} u_{2}}\left(g_{1}^{2}+\frac{h_{2}^{1}}{\hat{u}_{1}}\right) \\
& =-\frac{1}{d_{2} u_{2}}\left(g_{1}^{2}+\frac{f_{2}^{0}}{\hat{u}_{1}}\right) .
\end{aligned}
$$

Finally, for $i=1$, by lemma 6.3.11,

$$
\left[(U L)^{-1}\right]_{2,1}=\left[\left(F^{1} G^{(1)}\right)^{-1}\right]_{2,1}=-\frac{g_{1}^{1}}{d_{1} u_{1}}
$$

Finally, theorems 6.3.2, 6.3.3 and 6.3.4 together give the entries for the matrix

$$
\left[(U L)^{-1}\right]_{i: i+1, i: i+1}=\left[\begin{array}{cc}
\alpha^{i} & \gamma^{i} \\
\beta^{i} & \alpha^{i+1}
\end{array}\right], \quad i=1, \ldots, n-1,
$$

where

$$
\begin{aligned}
\alpha^{i} & :=\left[(U L)^{-1}\right]_{i, i}, & \gamma^{i} & :=\left[(U L)^{-1}\right]_{i, i+1} \\
\beta^{i} & :=\left[(U L)^{-1}\right]_{i+1, i}, & \alpha^{i+1} & :=\left[(U L)^{-1}\right]_{i+1, i+1} .
\end{aligned}
$$

Theorem 6.3.5 Given the notation above, for $i=4, \ldots, n-1$, we have

$$
\begin{aligned}
\alpha^{i} & =\frac{1-a_{i}}{d_{i} u_{i}}, & \gamma^{i} & =-\frac{1-a_{i}}{d_{i} u_{i} u_{i+1}} \\
\beta^{i} & =-\frac{1}{d_{i} u_{i}}\left(g_{1}^{i}+\frac{f_{2}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}\right), & \alpha^{i+1} & =\frac{1-a_{i+1}}{d_{i+1} u_{i+1}}
\end{aligned}
$$

where $a_{i}=\frac{f_{1}^{i-2}}{\hat{l}_{i-2} \hat{u}_{i-1}}+\frac{f_{2}^{i-3}}{\hat{l}_{i-3} \hat{u}_{i-1} \hat{u}_{i-2}}$.
Proof. Use theorems 6.3.4, 6.3.3 and 6.3.2.

## Chapter 7

## Implementation details and numerical examples

The implementation of 3dqds and all our numerical experiments were carried out in Matlab 7.2.0.232 (R2006a) on a Pentium M 1.6 GHz machine under Windows XP. All our computations were done in IEEE standard floating point arithmetic using double precision arithmetic with unit roundoff $\varepsilon=2^{-53} \approx 1.1 \times 10^{-16}$. We used the advantages of the exception handling feature incorporated in this arithmetic.

In what follows eigval refers to our function to obtain the eigenvalues of an unreduced real tridiagonal matrix using 3dqds and eig to MATLAB's function which uses LAPACK [2] routines that implement the QR algorithm on Hessenberg matrices. We report comparative results on carefully chosen tridiagonal matrices for test purposes: Bessel, Toeplitz, Clement, Liu's and symmetric matrices. We also present the first numerical results for a comparative study of some of the condition numbers presented in Chapter 2.

### 7.1 Implementation details

Next sections describe some of the details we had to consider in our implementation of eigval function. They are concerned with the shift strategy, the stopping criterion for deflation, the tolerance for element growth in 3dqds code, the need to choose a different shift after a failure and the possibility that the initial $L U$ factorization of the $J$ form of the input matrix $C$ does not exist.

### 7.1.1 Choosing a shift for 3 dqd s

Consider matrices $L$ and $U$ as described in (6.1), page 166, and remember that the shift information for 3dqds is coded in the first column of $(U L)^{2}-2(\Re \sigma) U L+|\sigma|^{2} I$ that, for a shift $\sigma$, is given by

$$
\left[\begin{array}{c}
\left(u_{1}+l_{1}\right)^{2}+u_{2} l_{1}-2(\Re \sigma)\left(u_{1}+l_{1}\right)+|\sigma|^{2} \\
u_{2} l_{1}\left(u_{1}+l_{1}+u_{2}+l_{2}-2(\Re \sigma)\right) \\
u_{2} l_{1} u_{3} l_{2} \\
0 \\
\vdots \\
0
\end{array}\right]
$$

So, to completely specify one iteration of 3dqds we need

$$
2(\Re \sigma)=\sigma+\bar{\sigma} \quad \text { and } \quad|\sigma|^{2}=\sigma \bar{\sigma}
$$

Since 3dqds is a restoring shift transformation, in the case of a real eigenvalue, a reasonable choice of a single shift is the $(n, n)$ element of $U_{k} L_{k}$, where $L_{k}$ and $U_{k}$ are the factors obtained after the $k^{\text {th }}$ iteration. The generalization for double shifting is to use the Francis shift, which means that $\sigma$ and $\bar{\sigma}$ are the eigenvalues of the bottom $2 \times 2$ corner of $U_{k} L_{k}$,

$$
\left[\begin{array}{cc}
l_{n-1}+u_{n-1} & 1  \tag{7.1}\\
u_{n} l_{n-1} & u_{n}
\end{array}\right]
$$

This will let us converge to either two real eigenvalues in the bottom $2 \times 2$ or a single $2 \times 2$ block with complex conjugate eigenvalues. When we are close to convergence, we expect
the $(n-1, n-2)$ entry, and possibly the $(n, n-1)$ entry, to be tiny so that the eigenvalues of this $2 \times 2$ submatrix are good approximations to eigenvalues of $J$.

If $\sigma$ and $\bar{\sigma}$ are the eigenvalues of (7.1) then

$$
\begin{array}{r}
\sigma+\bar{\sigma}=l_{n-1}+u_{n-1}+u_{n}=: \text { Sum } \\
\sigma \bar{\sigma}=u_{n} u_{n-1}=: \operatorname{Prod}
\end{array}
$$

and Sum and Prod will be the shift information given to 3dqds. When the ( $n, n-1$ ) entry is negligible, then $u_{n}$ will be an approximation for a real eigenvalue of $J$. If not, and if the ( $n-1, n-2$ ) entry is negligible then

$$
\begin{aligned}
& \sigma_{1}=\frac{1}{2}\left(\operatorname{Sum}+\sqrt{\text { Sum }^{2}-4 \operatorname{Prod}}\right) \\
& \sigma_{2}=\operatorname{Prod} / \sigma_{1}
\end{aligned}
$$

will be the approximations to the eigenvalues of the $2 \times 2$ corner block which may be both real or complex conjugate. Then deflation takes place and the algorithm proceeds with the remaining submatrix until approximations to all eigenvalues are found.

In our present shift strategy, each time deflation occurs, we may choose to use first zero shifts (Sum $=0$ and Prod $=0$ ) until $l_{n-1}$ or $l_{n-2}$ is small, and only then switch to the Francis shift. In general, this procedure improves convergence.

### 7.1.2 Criterion for deflation

Since we expect $U_{k} L_{k}$ to converge to a quasi-bidiagonal form, we use only tests to decide when $l_{n-1}$ or $l_{n-2}$ is negligible. These tests are performed before a 3dqds transformation takes place (not after) and consist in verifying if

$$
\left|l_{n-1}\right|<\operatorname{Tol}_{\text {Def }} \quad \text { or } \quad\left|l_{n-2}\right|<\operatorname{Tol}_{\text {Def }}
$$

for a given tolerance $T o l_{D e f}$. If we choose $T_{o l} l_{\text {Def }}=\varepsilon$, it may be too severe and it will take more time to converge; if we relax too much, we will converge faster but we may loose accuracy seriously. So, we must find a compromise between the speed of convergence and
the accuracy we want to attain. By default, in our experiments we used $T_{\text {ol }}^{\text {Def }}=\varepsilon$ since we were mainly concerned with accuracy. This is the most crucial parameter and needs more study.

### 7.1.3 Tolerance for element growth

The LU transformation fails when zero pivots do occur. Furthermore, for reasons of numerical stability, we do not accept the result of a particular transformation if, in consequence of the occurrence of pivots of very small size, elements of excessive size do appear in $L$ or $U$ factors. For this reason, we monitor the element growth and reject $L$ and $U$ when

$$
\max _{i}\left|l_{i}\right|>\text { Tol }_{\text {Growth }}\|J\| \quad \text { or } \quad \max _{i}\left|u_{i}\right|>\operatorname{Tol}_{\text {Growth }}\|J\|
$$

for a given tolerance Tol $_{\text {Growth }}$. When this happens, we signal a breakdown but it must be emphasized that is not a dramatic situation since we may recover. It is just a transformation whose result is neglected.

If we allow less element growth, we will have more breakdowns and it will take more time to converge. Allowing huge element growth will imply less accuracy. Again, there must be a trade-off between element growth and accuracy. Usually we let $\operatorname{Tol}_{\text {Growth }}=1 / \sqrt{\varepsilon}$.

The possibility of dealing with divisions by zero in IEEE arithmetic allows 3dqds transformation to be carried out till the very end, even when a pivot is zero or too small. Comparison of the size of the new values of $l_{i}$ and $u_{i}$ with $T o l_{\text {Growth }}\|J\|$ is only performed at the end of a transformation.

### 7.1.4 The shift after a failure

After a failure, there is the question of what shift to use next. We don't want to move away from the previous shift too much, just the necessary amount so that the 3dqds transformation does not breakdown. We also have to admit the possibility of a succession of failures.

So, while we do not have a successful transformation we will repeat the following

```
if \(\operatorname{Prod} \neq 0\)
    Sum \(=(1+\delta) *\) Sum
    Prod \(=(1+\delta)^{2} * \operatorname{Prod}\)
    \(\delta=2 * \delta\)
else
    small \(=\sqrt{\varepsilon} *\) local
    Sum \(=2 *\) small
    Prod \(=\) small \(^{2}\)
end if
```

where "local" is a local norm and, initially, we let $\delta=\sqrt{\varepsilon}$. This procedure forces to define a limited number of failures allowed.

### 7.1.5 Initial LU factorization

The $J$ form of $C$ may not have an LU factorization. In this case we must choose an initial shift $\tau$ so that the factorization

$$
J-\tau I=L U
$$

is possible and the 3dqds transform may start. Initially, we let $\tau=\theta$, for a small $\theta$, and then, if we are still not successful, we increase $\tau$ by $\theta$ until we have success. In the end $\tau$ must be added to the approximations given by eigval. But this choice needs more thought because it seems to affect the accuracy more than we expected.

### 7.2 Numerical examples

In the figures containing the plots of the computed eigenvalues that are presented in this section, the results obtained with eig are represented by a plus sign and the ones provided by eigval are represented with a big dot.

### 7.2.1 Bessel matrices

Bessel matrices, associated with generalized Bessel polynomials (see section 1.3.2, page 22), are nonsymmetric tridiagonals matrices defined by $B_{n}^{(a, b)}=\operatorname{tridiag}\left(\boldsymbol{\beta}^{(a, b)}, \boldsymbol{\alpha}^{(a, b)}, \boldsymbol{\gamma}^{(a, b)}\right)$ with

$$
\alpha_{1}^{(a, b)}=-\frac{a}{b}, \quad \gamma_{1}^{(a, b)}=-\alpha_{1}^{(a, b)}, \quad \beta_{1}^{(a, b)}=\frac{\alpha_{1}^{(a, b)}}{a+1},
$$

and

$$
\begin{aligned}
\alpha_{j}^{(a, b)}:=-b \frac{a-2}{(2 j+a-2)(2 j+a-4)}, \quad j=2, \ldots, n, \\
\gamma_{j}^{(a, b)}:=b \frac{j+a-2}{(2 j+a-2)(2 j+a-3)}, \\
\beta_{j}^{(a, b)}:=-b \frac{j}{(2 j+a-1)(2 j+a-2)}, \quad j=2, \ldots, n-1 .
\end{aligned}
$$

All the eigenvalues are simple and below we report three theorems from [44] that contain information about the localization of these eigenvalues. These theorems apply to the cases

$$
\begin{equation*}
a \in \mathbb{R}, \quad n>1-a, \quad b=2 . \tag{7.2}
\end{equation*}
$$

Theorem 7.2.1 Let the condition (7.2) be satisfied. Also let $n \geq 1$. Then all the eigenvalues of $B_{n}^{(a, b)}$ lie in the cardioil region

$$
C(n, a):=\left\{z=\rho \mathrm{e}^{\mathrm{i} \theta} \in \mathbb{C}: 0<\rho<\frac{1-\cos \theta}{n+a-1}\right\} \cup\left\{\frac{-2}{n+a-1}\right\} .
$$

Theorem 7.2.2 Let the condition (7.2) be satisfied. Also let $n \geq 2$. Then all the eigenvalues of $B_{n}^{(a, b)}$ belong to the sector

$$
S(n, a):=\left\{z=\rho \mathrm{e}^{\mathrm{i} \theta} \in \mathbb{C}:|\theta|>\cos ^{-1}\left(\frac{-a}{2 n+a-2}\right),-\pi<\theta \leq \pi\right\} .
$$

Theorem 7.2.3 Let the condition (7.2) be satisfied. Also let $n \geq 1$. Then all the eigenvalues of $B_{n}^{(a, b)}$ belong to the infinite region

$$
I(n, a):=\left\{z \in \mathbb{C}:|z|>\frac{2}{2 n+a-\frac{2}{3}}\right\} .
$$

Thus, the eigenvalues of the matrix $B_{n}^{(a, b)}$ lie in the intersection of the inclusion regions defined by the sets $C(n, a), S(n, a)$ and $I(n, a)$. But these eigenvalues suffer from ill-conditioning that increases with $n$. Matrices $B_{n}^{(a, b)}$ are close to defective matrices.

Figure 7.1 shows, along with the inclusion regions defined above, the results obtained with Matlab function eig (plus sign) and with eigval (big dot) for the classical case $a=b=2$, for $n=30$ and $n=40$. As $n$ increases, computed eigenvalues even fail to belong to the inclusion regions. But the answers of eigval are usually better.


Figure 7.1: Bessel matrix with $a=2$ and $b=2$

In [44] it is mentioned that the ill-conditining seems to reach its maximum when $a$ ranges from -8.5 to -4.5 . We also report the cases $a=-8.5, b=2$ (for $n=18$ and $n=25$ ), $a=-4.5, b=2$ (for $n=20$ and $n=25$ ) and $a=12, b=2$ (for $n=40$ and $n=50$ ). See Figure 7.2. The number of iterations needed for eigval to converge is approximately equal to $2 n$.

Bessel matrices are notoriously difficult. There is an interesting phenomenon of bifurcation in the eigenvalues computed by both algorithms and the regularity, in spite of the presence of ill-conditioning, is quite interesting. The approximations given by eigval follow the same pattern as the ones given by Matlab but seem to be better. As $n$ increases the huge deviation from the expected curve is not due to element growth - it rarely occurred.

(a) $a=-8.5, b=2 ; n=18$

(c) $a=-4.5, b=2 ; n=20$

(e) $a=12, b=2 ; n=40$

(b) $a=-8.5, b=2 ; n=25$

(d) $a=-4.5, b=2 ; n=25$

(f) $a=12, b=2 ; n=50$

Figure 7.2: Bessel matrices with $a=-8.5, b=2 ; a=-4.5, b=2$ and $a=12, b=2$

## Condition numbers

For a simple eigenvalue $\lambda \neq 0$ of a given matrix $A$, redefine now $\kappa_{\lambda}$ as the Wilkinson's relative condition number for $\lambda$ (see section 1.2.3, page 14), that is,

$$
\kappa_{\lambda}(A) \equiv \frac{\left\|y^{*}\right\|\|x\|}{\left|\lambda \| y^{*} x\right|}
$$

For $J=L U$ the relative condition numbers (see section 2.4.2, page 44)

$$
\operatorname{relcond}_{1}(\lambda ; L U)=\frac{|y|^{T} M_{1}|x|}{\left|y^{*} x\right||\lambda|}
$$

and

$$
\operatorname{relcond}_{2}(\lambda ; L U)=\frac{|y|^{T}(v+w)}{\left|y^{*} x\right|}
$$

give the sensitivity of eigenvalues to perturbations in the entries of the factors $L$ and $U$.
In Table 7.1 we display the condition number $\kappa_{\lambda}$ for $B_{n}^{(a, b)}$ with $a=-4.5, b=2$ and $n=20$, and for the corresponding $J$ form. We also display the condition numbers relcond ${ }_{1}$ and relcond 2 for the LU representation of $J$. It may be seen that, although the eigenvalues are much more sensitive to perturbations in $J$ than to perturbations in $B_{n}^{(a, b)}$, the situation improves significantly in the $L, U$ representation, that is, these factors define the eigenvalues not only much better than $J$ does but, also, they define better the eigenvalues than $B_{n}^{(a, b)}$ itself.

The eigenvectors we used to produce these numbers were the eigenvectors delivered by eig function of Matlab.

### 7.2.2 Clement matrices

The so-called Clement matrix (see [4])

$$
C_{n}=\operatorname{tridiag}(\boldsymbol{\beta}, \mathbf{0}, \boldsymbol{\gamma})
$$

with $\gamma=\left(\gamma_{j}\right), \gamma_{j}=j$ and $\boldsymbol{\beta}=\left(\beta_{j}\right), \beta_{j}=\gamma_{n-j}, j=1, \ldots, n-1$, has exact eigenvalues

$$
\begin{aligned}
& \pm n-1, n-3, \ldots, 1, \quad \text { for } n \text { even, } \\
& \pm n-1, n-3, \ldots, 0, \quad \text { for } n \text { odd. }
\end{aligned}
$$

| $\lambda$ | $\kappa_{\lambda}(B)$ | $\kappa_{\lambda}(J)$ | $\operatorname{relcond}_{1}(\lambda ; L U)$ | $\operatorname{relcond}_{2}(\lambda ; L U)$ |
| :---: | :---: | :---: | :---: | :---: |
| $-3.810^{-3}-6.710^{-2} \mathbf{i}$ | $610^{8}$ | $310^{22}$ | $310^{5}$ | $610^{5}$ |
| $-3.810^{-3}+6.710^{-2} \mathbf{i}$ | $610^{8}$ | $410^{22}$ | $510^{5}$ | $910^{5}$ |
| $-7.110^{-2}-1.610^{-2} \mathbf{i}$ | $510^{10}$ | $410^{25}$ | $510^{5}$ | $1010^{5}$ |
| $-7.110^{-2}+1.610^{-2} \mathbf{i}$ | $510^{10}$ | $410^{25}$ | $310^{6}$ | $1010^{6}$ |
| $-1.910^{-2}-7.210^{-2} \mathbf{i}$ | $110^{14}$ | $310^{23}$ | $310^{6}$ | $210^{6}$ |
| $-1.910^{-2}+7.210^{-2} \mathbf{i}$ | $110^{14}$ | $310^{23}$ | $910^{8}$ | $110^{8}$ |
| $-3.410^{-2}-7.210^{-2} \mathbf{i}$ | $110^{14}$ | $210^{24}$ | $810^{8}$ | $110^{8}$ |
| $-3.410^{-2}+7.210^{-2} \mathbf{i}$ | $110^{12}$ | $210^{24}$ | $310^{10}$ | $310^{9}$ |
| $-6.510^{-2}-4.610^{-2} \mathbf{i}$ | $110^{12}$ | $310^{25}$ | $310^{10}$ | $210^{9}$ |
| $-6.510^{-2}+4.610^{-2} \mathbf{i}$ | $110^{14}$ | $310^{25}$ | $1010^{10}$ | $610^{9}$ |
| $-4.710^{-2}-6.910^{-2} \mathbf{i}$ | $110^{14}$ | $610^{24}$ | $1010^{10}$ | $710^{9}$ |
| $-4.710^{-2}+6.910^{-2} \mathbf{i}$ | $110^{13}$ | $610^{24}$ | $510^{11}$ | $310^{10}$ |
| $-6.010^{-2}-6.610^{-2} \mathbf{i}$ | $110^{13}$ | $910^{24}$ | $410^{11}$ | $310^{10}$ |
| $-6.010^{-2}+6.610^{-2} \mathbf{i}$ | $110^{14}$ | $910^{24}$ | $110^{12}$ | $910^{10}$ |
| $-8.010^{-2}-5.910^{-2} \mathbf{i}$ | $110^{14}$ | $410^{24}$ | $110^{12}$ | $810^{10}$ |
| $-8.010^{-2}+5.910^{-2} \mathbf{i}$ | $110^{14}$ | $410^{24}$ | $310^{12}$ | $210^{11}$ |
| $-1.010^{-1}-4.310^{-2} \mathbf{i}$ | $110^{14}$ | $110^{24}$ | $210^{12}$ | $110^{11}$ |
| $-1.010^{-1}+4.310^{-2} \mathbf{i}$ | $210^{14}$ | $110^{24}$ | $310^{12}$ | $210^{11}$ |
| $-1.210^{-1}-1.610^{-2} \mathbf{i}$ | $210^{14}$ | $510^{23}$ | $210^{12}$ | $210^{11}$ |
| $-1.210^{-1}+1.610^{-2} \mathbf{i}$ | $210^{14}$ | $510^{23}$ | $210^{12}$ | $110^{11}$ |

Table 7.1: Relative condition numbers for the eigenvalues of $B_{20}^{(-4.5,2)}$

Figure 7.3 illustrates the results of both algorithms for the Clement matrices with $n=150$ and $n=200$. All the approximations given by eigval are real while the approximations obtained with Matlab have large imaginary parts.


Figure 7.3: Clement matrix

Table 7.2 shows the largest and the smallest Wilkinson's relative condition number, $\kappa_{\text {max }} \equiv \max _{\lambda} \kappa_{\lambda}\left(C_{n}\right)$ and $\kappa_{\text {min }} \equiv \min _{\lambda} \kappa_{\lambda}\left(C_{n}\right)$, and also the largest and smallest of the relative condition number relcond ${ }_{1}(\lambda ; L U), \max _{\text {relcond }_{1}}$ and $\min _{\text {relcond }_{1}}$, respectively.

| $n$ | $\kappa_{\max }$ | $\kappa_{\min }$ | $\max _{\text {relcond }_{1}}$ | $\min _{\text {relcond }_{1}}$ |
| :---: | :---: | :---: | :---: | :---: |
| 150 | $3.510^{20}$ | $1.010^{4}$ | $6.410^{4}$ | $2.810^{1}$ |
| 200 | $1.210^{28}$ | $2.310^{11}$ | $2.310^{4}$ | $2.310^{1}$ |

Table 7.2: Relative condition numbers for the Clement matrix

The minimum and maximum relative errors, $r e l_{\text {min }}$ and $r e l_{\text {max }}$, respectively, for $n=150$ and $n=200$, as well as for $n=300$ and $n=450$, are presented in Table 7.3. When $n$ increases almost all the approximations given by Matlab exhibit significantly high relative errors, while with eigval we can consider that all the approximations produced have satisfactory relative accuracy.

| eigval |  |  |  | eig |  |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $n$ | $r e l_{\text {min }}$ | $r e l_{\text {max }}$ | rel $_{\text {min }}$ | rel $_{\text {max }}$ |  |
| 150 | $7.810^{-12}$ | $8.110^{-9}$ | $3.210^{-15}$ | $7.810^{1}$ |  |
| 200 | $1.210^{-11}$ | $6.410^{-9}$ | $2.910^{-16}$ | $1.110^{2}$ |  |
| 300 | $3.610^{-11}$ | $1.110^{-8}$ | $6.610^{-3}$ | $1.710^{2}$ |  |
| 450 | $4.510^{-12}$ | $1.810^{-8}$ | $4.510^{-3}$ | $2.610^{2}$ |  |

Table 7.3: Relative errors for the Clement matrix
Matlab function eig does not detect that Clement matrices are symmetrizable matrices. Our function eigval gives much better results than eig, almost as good as the ones Matlab delivers if we give as input a symmetric matrix similar to $C_{n}$.

### 7.2.3 Liu's matrices

Z. S. Liu [31] devised an algorithm to obtain one-point spectrum unreduced tridiagonal matrices of arbitrary dimension $n \times n$. These matrices, that we already introduced in Chapter 4 and called Liu's matrices, have only one eigenvalue, zero with multiplicity $n$, and the Jordan form consists of one Jordan block. We represent Liu's matrices as

$$
\operatorname{Liu}_{n}=\operatorname{tridiag}\left(\mathbf{1}^{n}, \boldsymbol{\alpha}^{n}, \boldsymbol{\gamma}^{n}\right)
$$

where $\mathbf{1}^{n}$ always stands for a vector of 1 's of length $n-1$. Notice that the transpose of $L i u_{n}$ is already in $J$ form.

We considered $L i u_{n}$ for $n=6$,

$$
\operatorname{Liu}_{6}=\left[\begin{array}{cccccc}
0 & -1 & & & & \\
1 & 0 & 1 & & & \\
& 1 & -1 & -1 & & \\
& & 1 & 1 & 1 & \\
& & & 1 & 0 & -1 \\
& & & & 1 & 0
\end{array}\right]
$$

for $n=14$,

$$
\begin{aligned}
& \boldsymbol{\alpha}^{14}=\left[\begin{array}{lllllllllllll}
0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0
\end{array}\right]^{T} \\
& \gamma^{14}=\left[\begin{array}{lllllllllllll}
-1 & 1 & 1 & -1 & 1 & -1 & -1 & -1 & 1 & -1 & 1 & 1 & -1
\end{array}\right]^{T},
\end{aligned}
$$

and for $n=28$,

$$
\begin{aligned}
\boldsymbol{\alpha}^{28} & =[0,0,0,0,0,0,-1,1,0,0,0,0,0,-1,1,0,0,0,0,0,1,-1,0,0,0,0,0,0]^{T} \\
\boldsymbol{\gamma}^{28} & =[-1,1,1,-1,1,-1,-1,-1,1,-1,1,1,-1,-1,-1,1,1,-1,1,-1,-1,-1,1,-1,1,1,-1]^{T} .
\end{aligned}
$$

Now a note on perturbation theory for multiple eigenvalues. Consider the example of perturbing by $\epsilon$ the $(n, 1)$ entry of an $n \times n$ Jordan block. Let

$$
\widetilde{C}=\left[\begin{array}{cccc}
0 & 1 & & \\
& \ddots & \ddots & \\
& & \ddots & 1 \\
\epsilon & & & 0
\end{array}\right]
$$

Then the characteristic equation changes from $\lambda^{n}=0$ to $\lambda^{n}-\epsilon=0$. So the eigenvalues of the perturbed matrix $\widetilde{C}$ are the $n$ possible complex roots of $\epsilon, \lambda_{k}=\sqrt[n]{\epsilon} \mathrm{e}^{\frac{\mathbf{i} \frac{2 k \pi}{n}}{n}}, k=0, \ldots, n-1$. The $n^{\text {th }}$ root of $\epsilon$ grows much faster than any multiple of $\epsilon$ for small $\epsilon$. More formally, the condition number of a multiple eigenvalue is infinite because at $\epsilon=0$, for $n \geq 2$,

$$
\frac{d \lambda}{d \epsilon}=\frac{1}{n \epsilon^{1-\frac{1}{n}}}=\infty .
$$

For example, if we take $n=16$ and $\epsilon=10^{-16}$, then for each eigenvalue $\lambda_{k}$, we have $\left|\lambda_{k}\right|=0.1$, a change $10^{15}$ times greater than $\epsilon$. However, having an infinite condition number does not mean that the eigenvalues cannot be computed with any correct digits.

The more apart from the diagonal entries the perturbations occur, the worse the effect on the eigenvalues seems to be. So, we decided to show our numerical results together with the circles $|z|=\sqrt[n]{\varepsilon}$ where $\varepsilon$ is the unit roundoff. If the approximations are inside this circle we may consider the results good enough, since they are not worse than the ones that we would get in exact arithmetic for the Jordan block similar to Liu's matrix perturbing by $\varepsilon$ the $(1, n)$ entry .

The results we obtained with eigval and eig for $n=6, n=14$ and $n=28$ are shown in the following figures.


Figure 7.4: Liu's matrices


Figure 7.5: Liu's matrix for $n=28$

The accuracy of our approximations is slightly better than the accuracy of those provided by Matlab and the convergence is surprisingly quick. The number of iterations needed for eigval to converge is less than $2 n$.

### 7.2.4 Toeplitz matrices

A tridiagonal Toeplitz matrix has the form

$$
T_{n}=\left[\begin{array}{ccccc}
a & c & & & \\
b & a & c & & \\
& \ddots & \ddots & \ddots & \\
& & b & a & c \\
& & & b & a
\end{array}\right]
$$

Such matrices arise, for example, when discretizing partial differential equations or boundary value problems for ordinary differential equations [Higham, p.522] . The eigenvalues are known explicitly:

$$
a+2(b c)^{1 / 2} \cos \left(\frac{k \pi}{n+1}\right), \quad k=1, \ldots, n .
$$

If $b c<0$ then the exact eigenvalues are complex with real part equal to $a$.
We first show the results for $T_{50}$ and $T_{80}$ with $a=1, b=2$ and $c=-1$. In Figure 7.6 we also represent the exact solutions with a times sign $(\times)$. For an exact eigenvalue $\lambda$ and an approximation $\hat{\lambda}$ we do not represent $\lambda$ or $\hat{\lambda}$ but $\lambda-a$ and $\hat{\lambda}-a$.


Figure 7.6: Toeplitz matrix with $a=1, b=2$ and $c=-1$

We also show the results for $n=150$ and $n=200$.


Figure 7.7: Toeplitz matrix with $a=1, b=2$ and $c=-1$

The accuracy of the approximations given by eigval is better than the accuracy of the approximations delivered by Matlab. See Table 7.4 for a report on the relative errors for the different values of $n$.

| eigval |  |  |  | eig |  |
| ---: | :---: | :---: | :---: | :---: | :---: |
| $n$ | $r e l_{\text {min }}$ | $r e l_{\text {max }}$ | $r e l_{\text {min }}$ | rel $_{\text {max }}$ |  |
| 50 | $3.810^{-14}$ | $2.610^{-11}$ | $3.810^{-12}$ | $3.610^{-10}$ |  |
| 80 | $5.410^{-14}$ | $3.510^{-10}$ | $5.010^{-9}$ | $1.210^{-6}$ |  |
| 150 | $1.210^{-13}$ | $4.310^{-5}$ | $3.210^{-3}$ | $1.810^{-1}$ |  |
| 200 | $2.710^{-13}$ | $2.110^{-1}$ | $2.110^{-3}$ | $2.310^{-1}$ |  |

Table 7.4: Relative errors for the Toeplitz matrix with $a=1, b=2$ and $c=-1$

### 7.2.5 Symmetric matrices

Finally, as an example of a symmetric matrix we will just consider the tridiagonal Toeplitz $T_{n}$ as described in the previous section with $b=c$. For this case the exact eigenvalues are

$$
a+2 b \cos \left(\frac{k \pi}{n+1}\right), \quad k=1, \ldots, n
$$

The behavior of eig function of Matlab is better than eigval. In table 7.5 we show the bounds for the relative errors of the approximations obtained with eigval for the case of $a=5, b=1$ and $c=1$. The approximations provided by Matlab have full accuracy.

| eigval |  |  |
| ---: | :---: | :---: |
| $n$ | rel $_{\text {min }}$ | rel $l_{\text {max }}$ |
| 50 | $4.810^{-15}$ | $5.410^{-12}$ |
| 100 | $1.810^{-15}$ | $1.410^{-11}$ |
| 200 | $3.010^{-14}$ | $2.110^{-8}$ |

Table 7.5: Relative errors for the Toeplitz matrix with $a=5$ and $b=c=1$

## Chapter 8

## Summary

The unsymmetric eigenvalue problem is, in general, much harder than the real symmetric eigenvalue problem for the basic reason that it is not always well posed. More precisely, the derivative of an eigenvalue with respect to certain matrix entries may be infinite. When executed in computer arithmetic any method finds itself aiming at a target (an eigenvalue) that may change at every step in the process. In contrast the derivatives in the symmetric case are all bounded by 1 .

The ideal algorithm should compute each eigenvalue to nearly the accuracy to which it is determined by the data (matrix entries). Ideally the algorithm should report what that attainable accuracy is. The traditional way to do this is to deliver well chosen condition numbers along with the computed eigenvalues. It turns out that computing the condition numbers is just as difficult as computing the eigenvalues themselves.

We focus on the tridiagonal eigenproblem which occurs in its own right (for instance with Bessel polynomials) and also as a condensed form of a real square matrix. The first contribution of this thesis is to consider the sensitivity of the eigenvalues to different representations of an unreduced tridiagonal matrix $C$ :
a) the non-trivial entries of $L$ and $D$ where $T \Delta=D_{1}^{-1} C D_{1}=L D U=L D L^{T} \Delta$ with $\Delta$ a signature matrix and $T$ symmetric;
b) the non-trivial entries of $L$ and $U$ where $J=D_{2}^{-1} C D_{2}=L U$, and the $(i, i+1)$ entries of $J$ are all 1 .

These new measures are, in general, smaller than Wilkinson's condition number and are sometimes quite realistic despite being upper bounds. Our measures take advantage of the tridiagonal form and play a fundamental role in assessing attainable accuracy. See Chapter 7 for examples.

A second contribution is purely mathematical. We follow a hint given in Wilkinson's monumental book "The Algebraic Eigenvalue Problem" and produce a rigorous proof that the LR algorithm, and our dqd algorithm, converge without breakdown to the nonzero eigenvalue of an unreduced tridiagonal matrix with a one-point spectrum. The context is exact arithmetic and the rate of convergence is very slow $(\mathcal{O}(1 / k)$ as $k \rightarrow \infty)$ but the surprise is that the algorithm actually does converge in the presence of a (large) Jordan block.

The main contribution of this thesis is practical: the presentation of a robust and efficient algorithm 3dqds that computes both real and complex eigenvalues of a real tridiagonal matrix while employing real arithmetic throughout the computation. In general, in our numerical tests, the output is more accurate than Matlab's procedure eig and requires $\mathcal{O}\left(n^{2}\right)$ arithmetic effort instead of Matlab's $\mathcal{O}\left(n^{3}\right)$.

To the best of our knowledge this is the first robust algorithm that takes advantage of tridiagonal form. However more work needs to be done in tuning some of the parameters. Deep study of a variety of cases is needed. The basic difficulty is easy to state: demand too much accuracy and any procedure will fail (never converge), demand too little accuracy and the unnecessary errors in accepting one eigenvalue can spoil the accuracy in another one.

### 8.1 Future work

Our test bed is far too small to draw conclusions. We are not yet in a position to claim that eigval function (that uses 3dqds code) always delivers more accurate approximations than eig function of Matlab. The first results we obtained are very good results but we need a a more comprehensive study of test matrices with a variety of distributions.

Finding a general shift strategy that is efficient over a wide range of matrices is very difficult. It is a research topic in itself. In future work we plan to dedicate effort in trying to improve the shift strategy implemented in the present version of our code.

We also plan to develop a more sophisticated deflation strategy that monitors convergence to each eigenvalue (or conjugate pair) to detect when the process has stagnated, i.e., when further steps will not reduce the corresponding off-diagonal entry.

A different aspect of 3dqds is that it converts the given tridiagonal matrix to $L U$ form via the factorization $J-\sigma I=L U$. This initial shift affects the attainable accuracy and needs more study. In this context, the issues of a numerical error analysis need to be considered.

We have now a tool to tackle the tridiagonal eigenvalue problem. It will be of interest to compare the performance of eigval with the possible rival algorithms: Françoise Tisseur's Ehrlich-Aberth method [3], the Extended HR algorithm of A. Liu [31] (unpublished finite precision implementation) and David Day's complex dqds algorithm [7] (unpublished preliminary version).

The computation of row and column eigenvectors is beyond the scope of this thesis but it is a necessary component to assess the attainable accuracy and to refine initial approximations. Our ultimate goal is to produce a software package that is both more efficient, more accurate and more informative than Matlab's eig procedure.

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[^0]:    ${ }^{1}$ Matlab, shortcut for "Matrix Laboratory", is a commercial program sold by The Mathworks, Inc.

[^1]:    ${ }^{1}$ In German the word eigen means characteristic or special.

[^2]:    ${ }^{1}$ LaPACK is a library of Fortran 77 subroutines for solving the most common problems in numerical linear algebra.

[^3]:    ${ }^{2}$ One ulp, "unit in last place", of the normalized floating point number $y= \pm \beta^{e} \times . d_{1} d_{2} \ldots d_{t}$ is $\operatorname{ulp}(y)=\beta^{e} \times .00 \ldots 1=\beta^{e-t}$. It is the right unit for discussing relative errors since it avoids reference to the magnitude of the numbers involved.

